

# VECTOR BOSON SIGNALS OF ELECTROWEAK SYMMETRY BREAKING

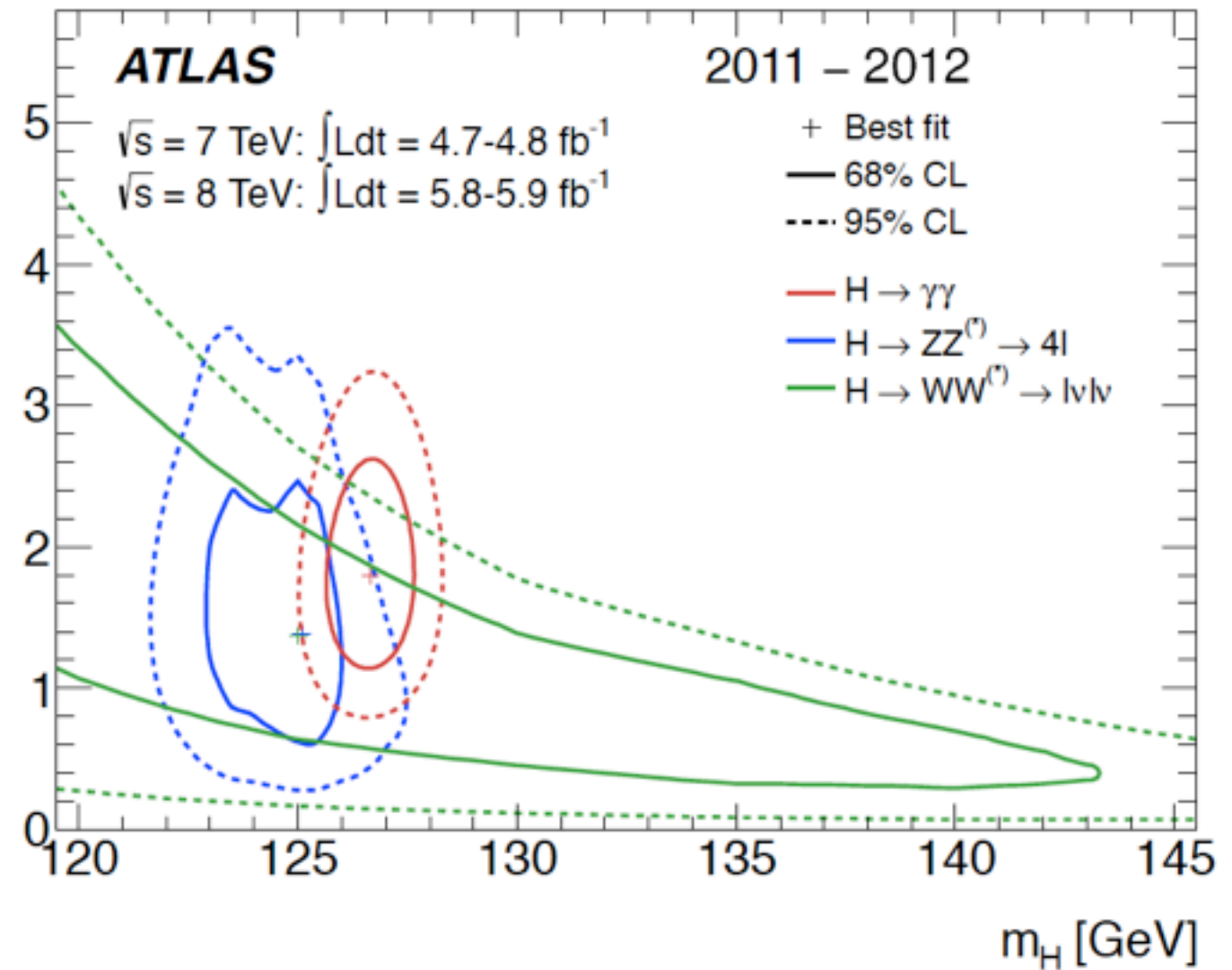
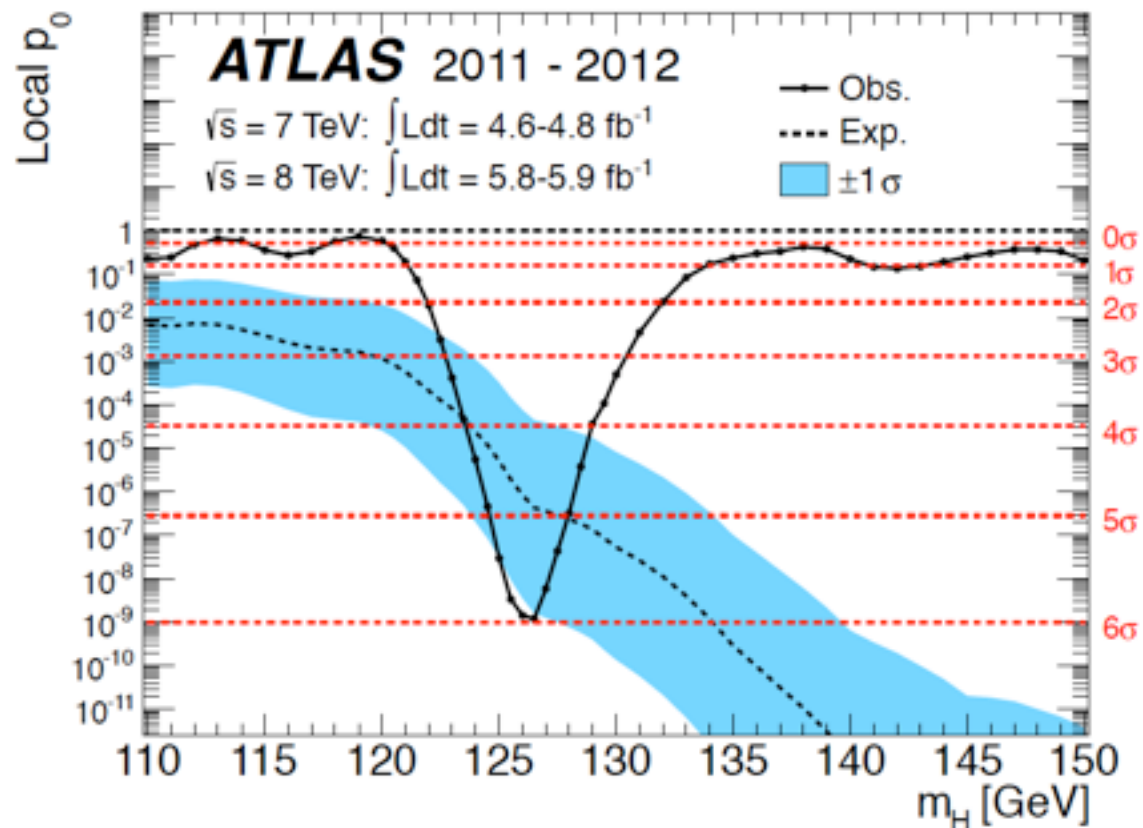
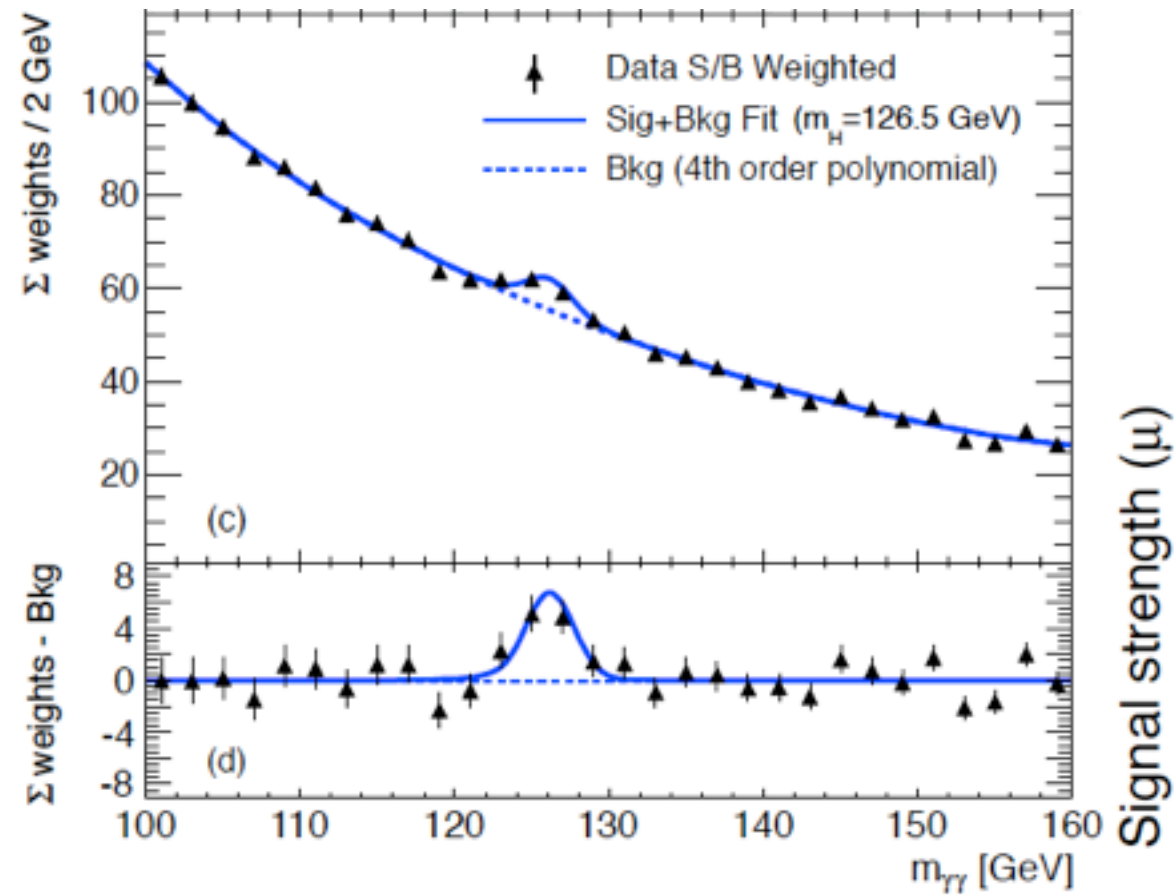
**R. SEKHAR CHIVUKULA**  
**MICHIGAN STATE UNIVERSITY**  
**NOVEMBER 13, 2012**

清华大学高能物理研究中心

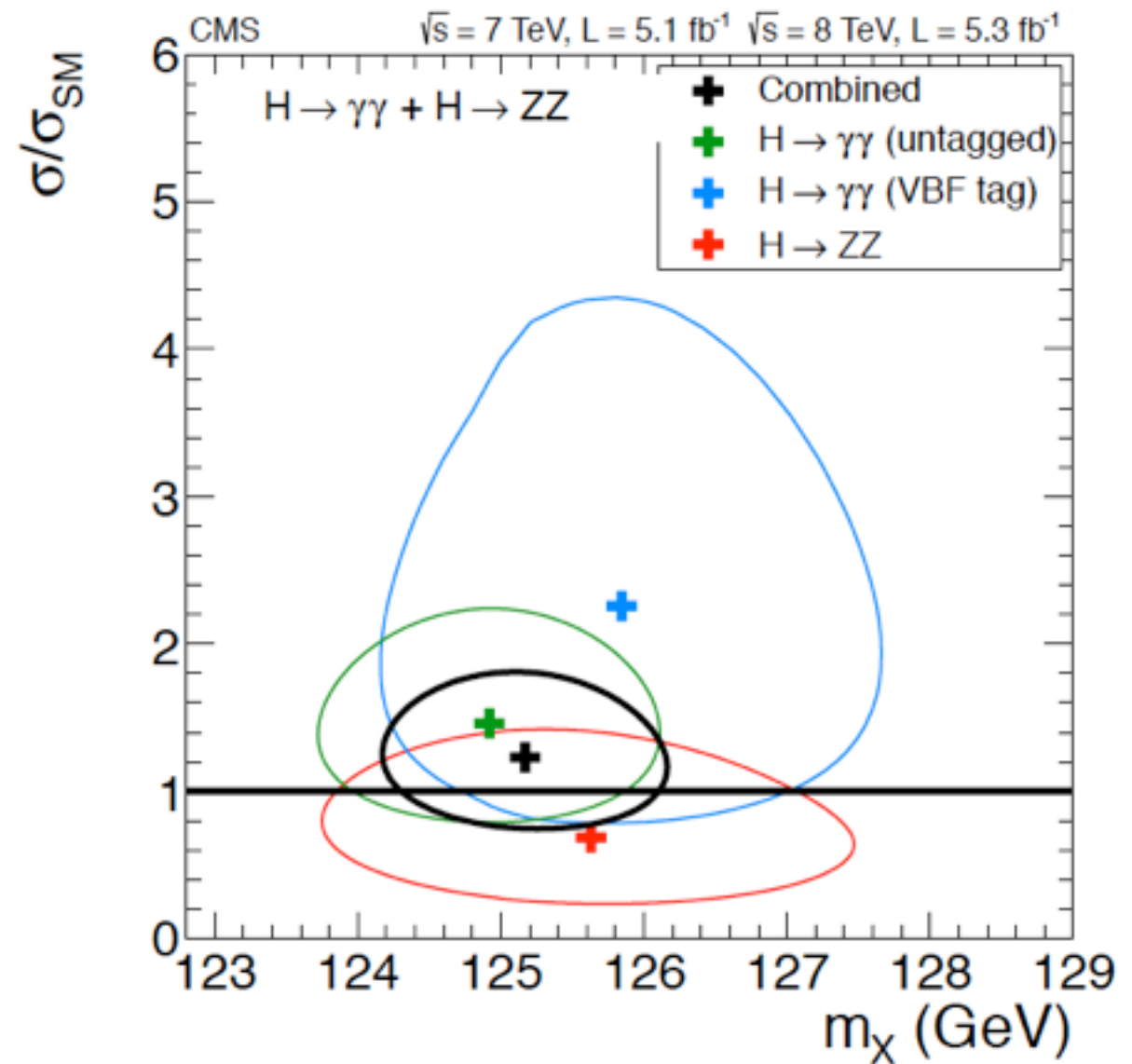
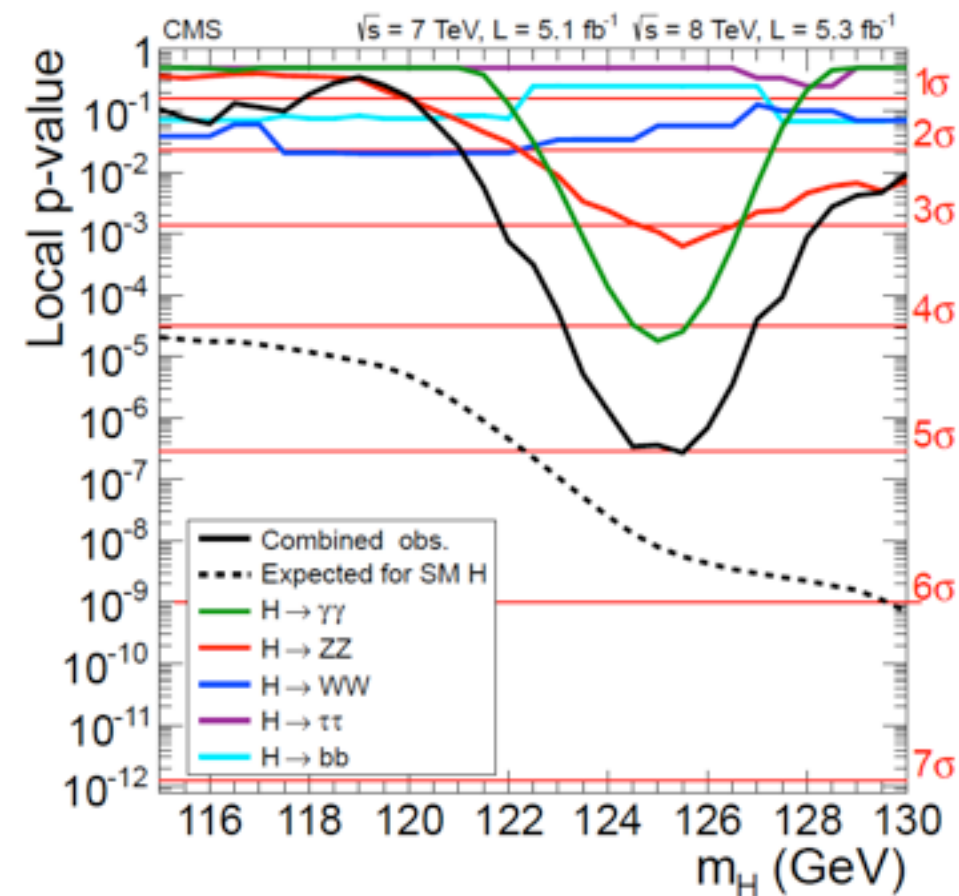
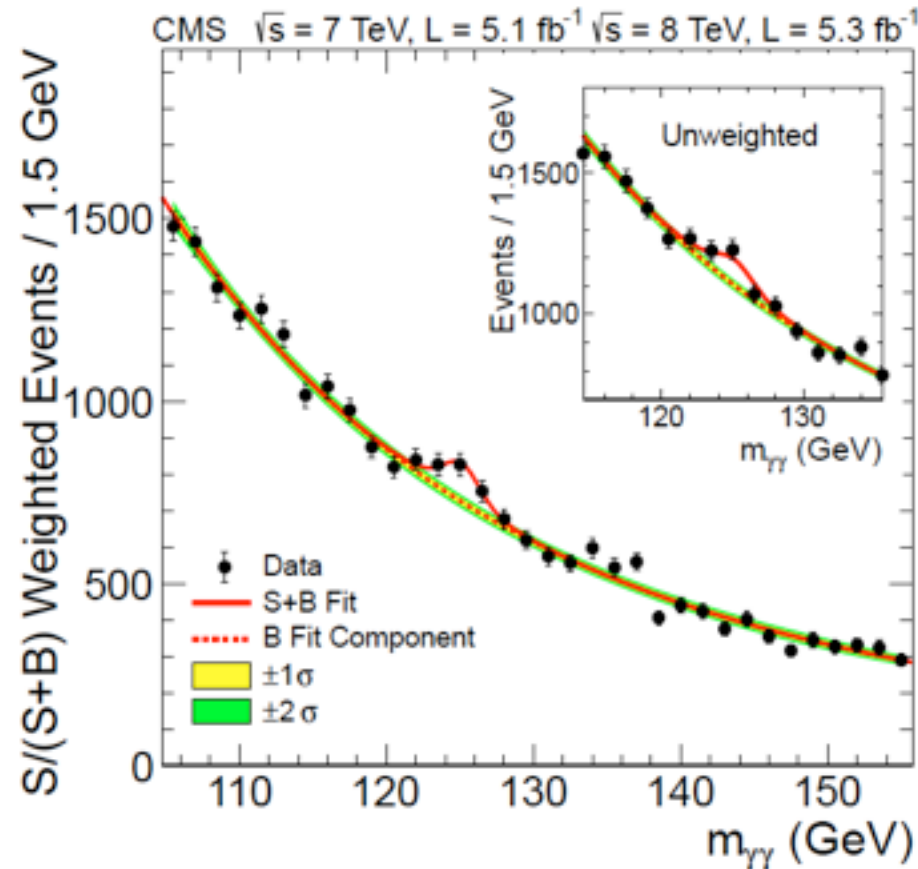
*Center for High Energy Physics, Tsinghua University*

# A NEW BOSON

# ATLAS



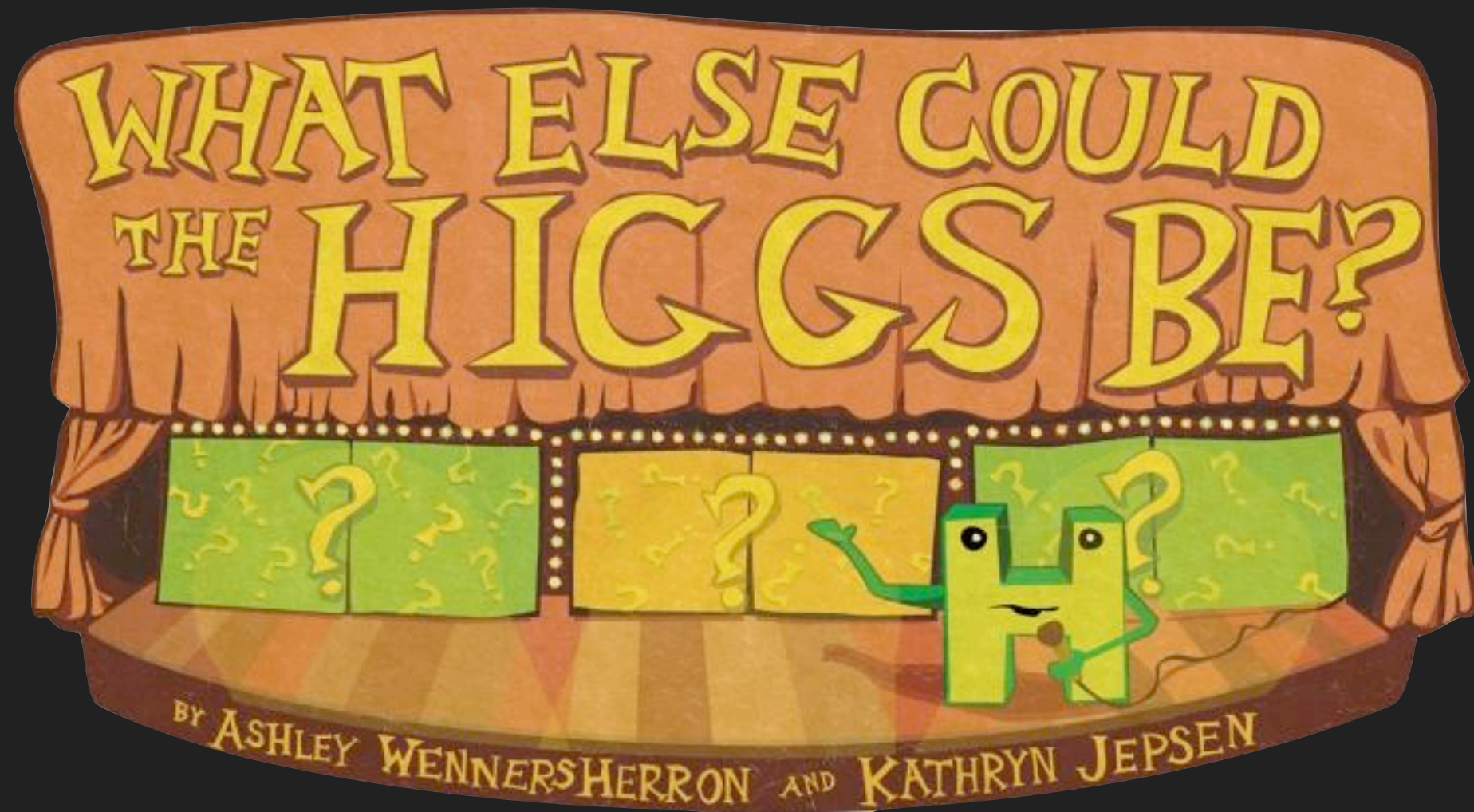
CERN-PH-EP-2012-218  
arxiv:1207.7214



CERN-PH-EP-2012-220  
arxiv:1207.7235

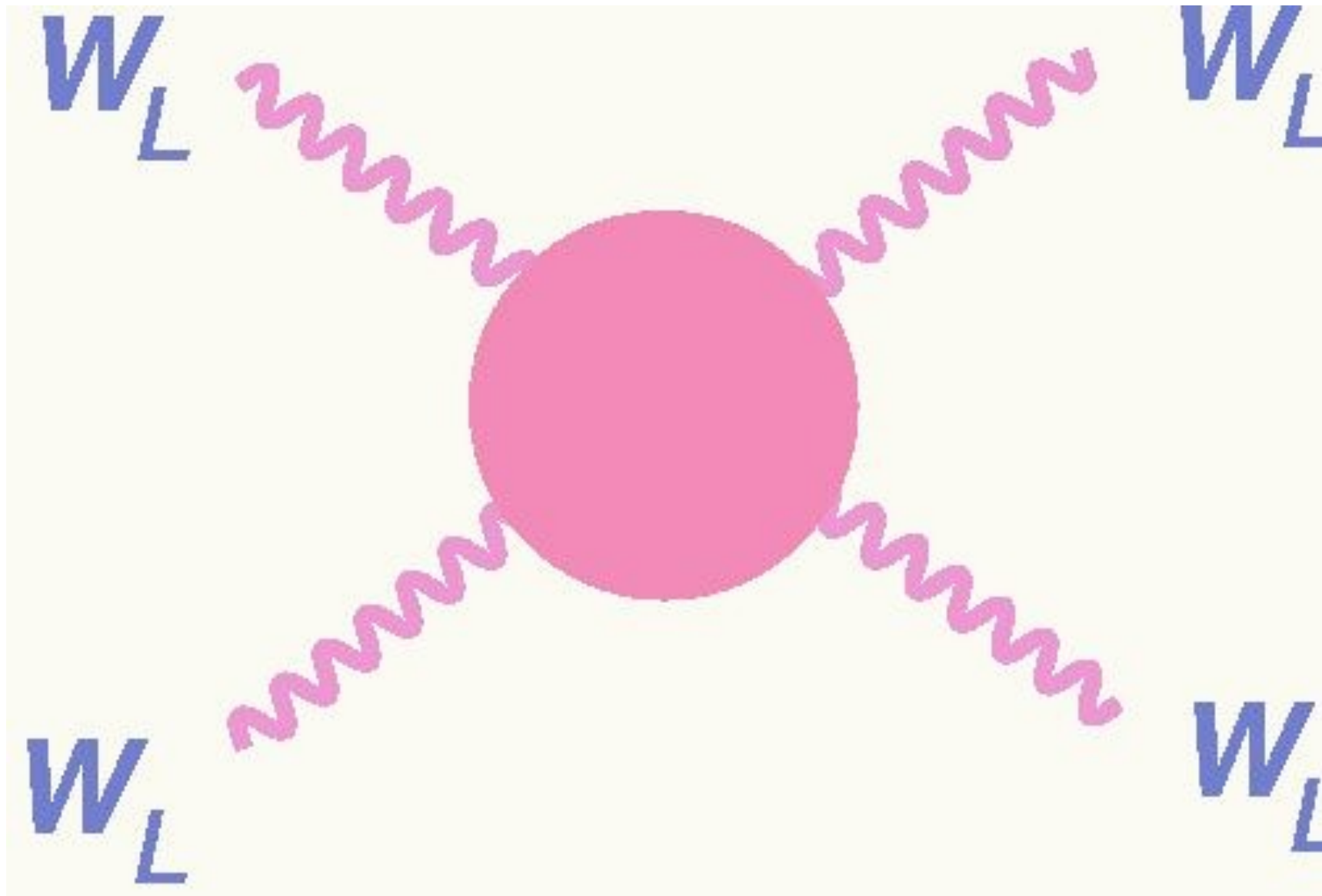


# IS THE NEW BOSON THE HIGGS BOSON?

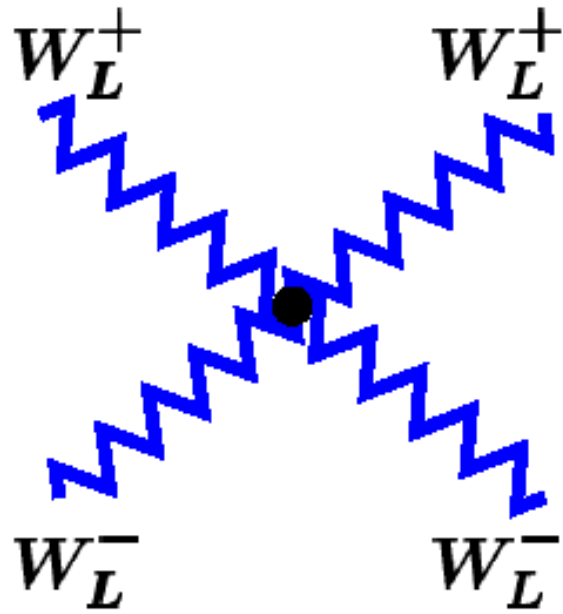


SYMMETRY MAGAZINE,  
OCT 30, 2012

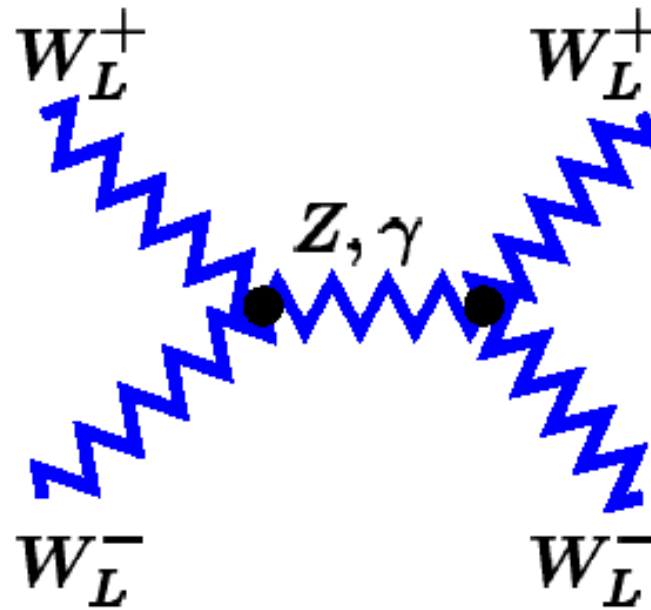
# Loss of Unitarity in



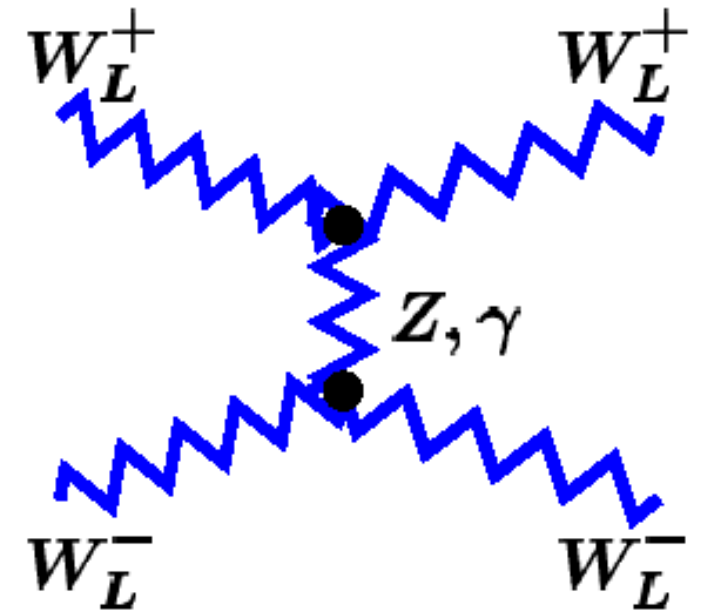
# SU(2) x U(1) @ E<sup>4</sup>



(a)



(b)



(c)

Graphs

$$g^2 \frac{E^4}{m_w^4}$$

(a)  $-3 + 6 \cos \theta + \cos^2 \theta$

(b)  $-4 \cos \theta$

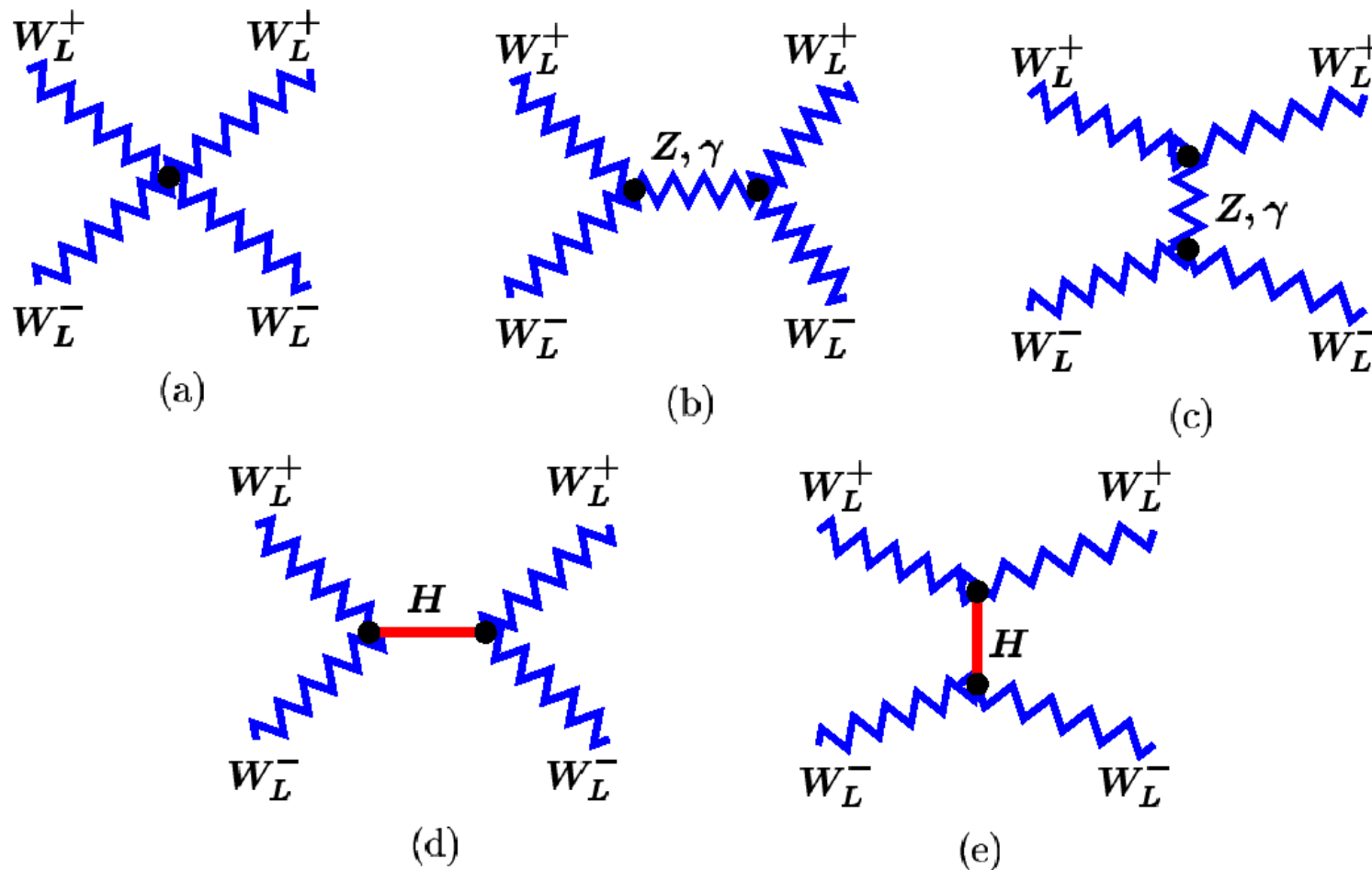
(c)  $+3 - 2 \cos \theta - \cos^2 \theta$

Sum

$$0$$

$$\epsilon_L^\mu(k) = \frac{k^\mu}{m_w} + \mathcal{O}\left(\frac{m_w}{E}\right)$$

# SU(2) x U(1) @ $E^2$ & THE HIGGS



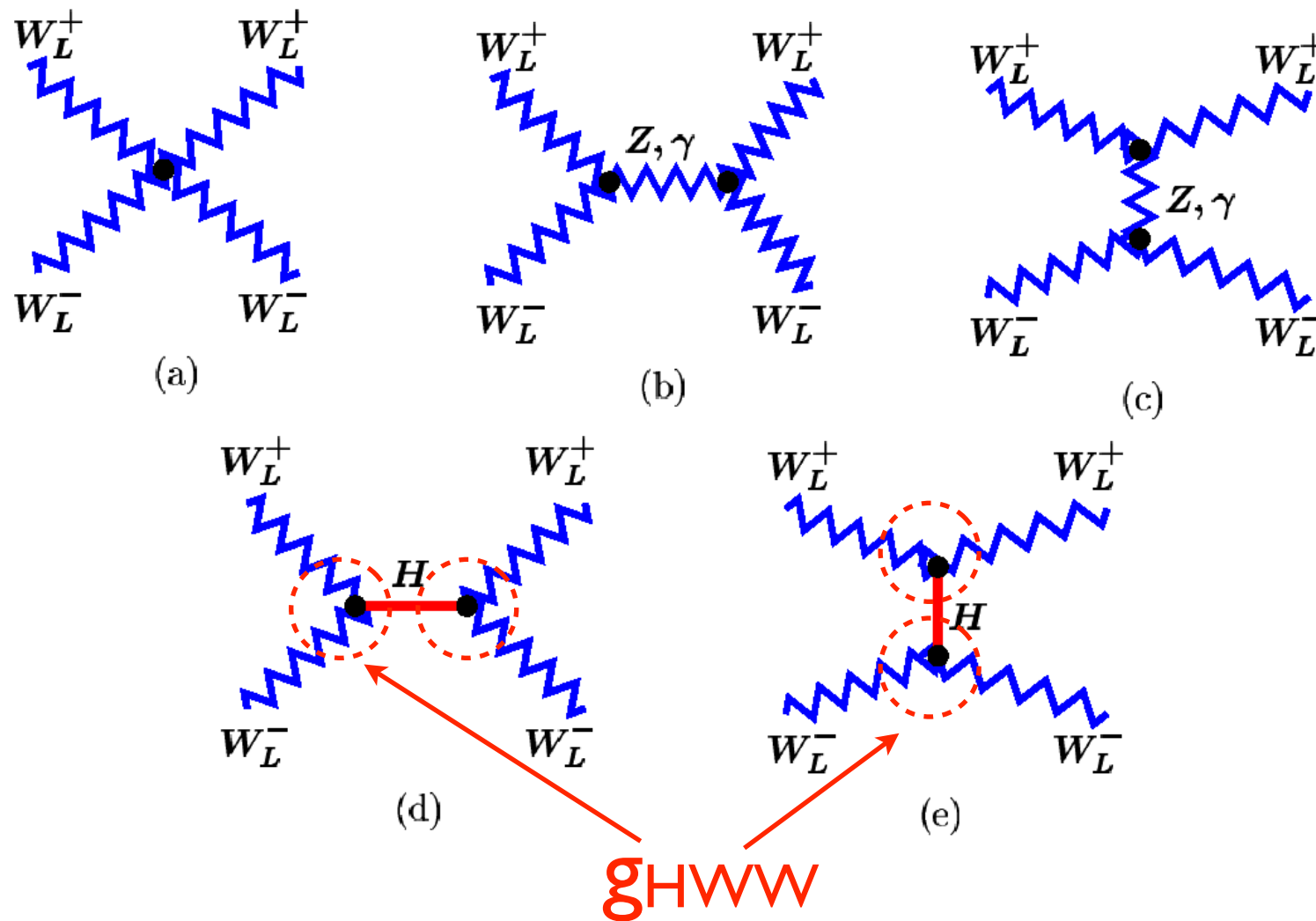
Graphs	$g^2 \frac{E^2}{m_w^2}$
(a)	$+2 - 6 \cos\theta$
(b)	$-\cos\theta$
(c)	$-\frac{3}{2} + \frac{15}{2} \cos\theta$
(d + e)	$-\frac{1}{2} - \frac{1}{2} \cos\theta$
<b>Sum</b> including (d+e)	<b>0</b>

►  $\mathcal{O}(E^0) \Rightarrow$  4d  $m_H$  bound:  $m_H < \sqrt{16\pi/3} v \simeq 1.0 \text{ TeV}$

► If no Higgs  $\Rightarrow \mathcal{O}(E^2) \Rightarrow E < \sqrt{4\pi} v \simeq 0.9 \text{ TeV}$



# SU(2) x U(1) @ $E^2$ & THE HIGGS



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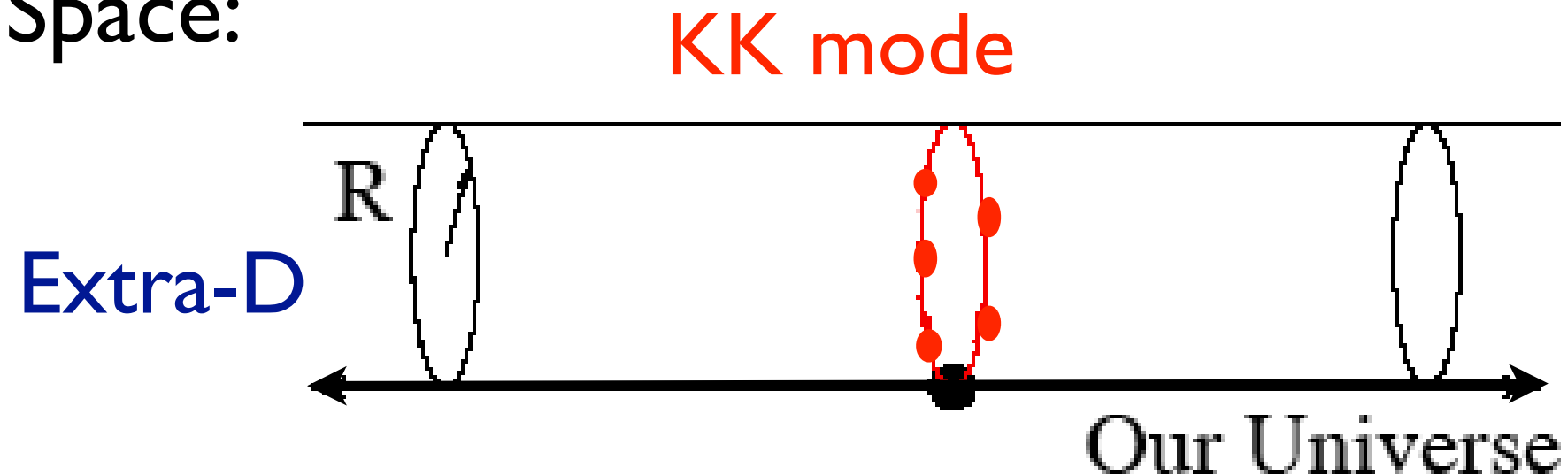
► If no Higgs  $\Rightarrow \mathcal{O}(E^2) \Rightarrow E < \sqrt{4\pi} v \simeq 0.9 \text{ TeV}$

# **ALTERNATIVES?**

## **“HIGGSLESS” MODELS**

# MASSIVE VECTOR BOSONS FROM 5-D

Flat Space:



Expand 5-D gauge bosons in eigenmodes: e.g.  
for  $S^1/Z_2$ :

$$\hat{A}_\mu^a = \frac{1}{\sqrt{\pi R}} \left[ A_\mu^{a0}(x_\nu) + \sqrt{2} \sum_{n=1}^{\infty} A_\mu^{an}(x_\nu) \cos\left(\frac{nx_5}{R}\right) \right]$$

$$\hat{A}_5^a = \sqrt{\frac{2}{\pi R}} \sum_{n=1}^{\infty} A_5^{an}(x_\nu) \sin\left(\frac{nx_5}{R}\right)$$

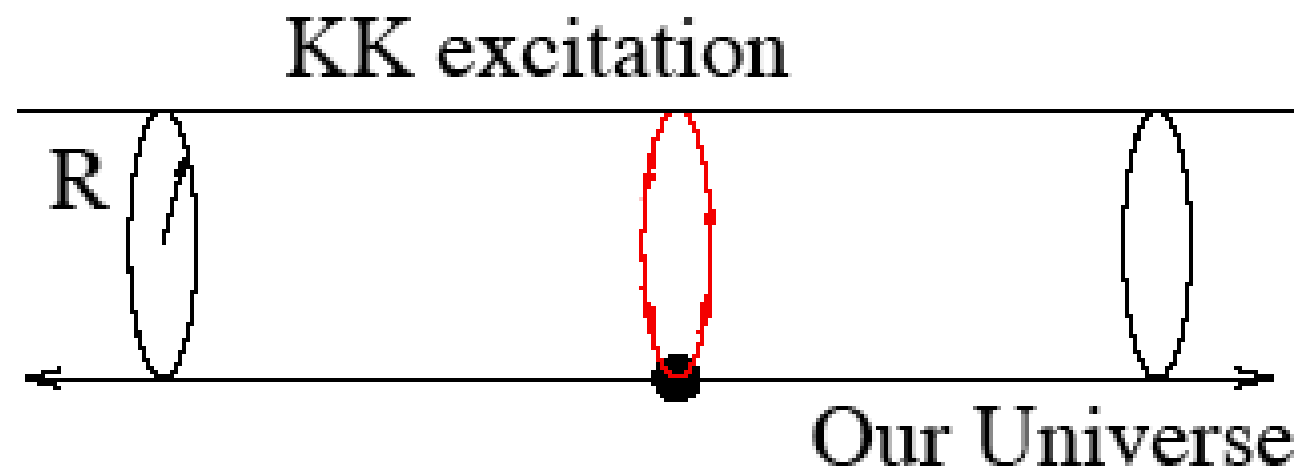
4-D gauge kinetic term contains

$$M_n = \frac{n}{R}$$

$$\frac{1}{2} \sum_{n=1}^{\infty} \left[ M_n^2 (A_\mu^{an})^2 - 2M_n A_\mu^{an} \partial^\mu A_5^{an} + (\partial_\mu A_5^{an})^2 \right]$$

i.e.,  $A_L^{an} \leftrightarrow A_5^{an}$

# ENERGY SCALES AND COUPLINGS



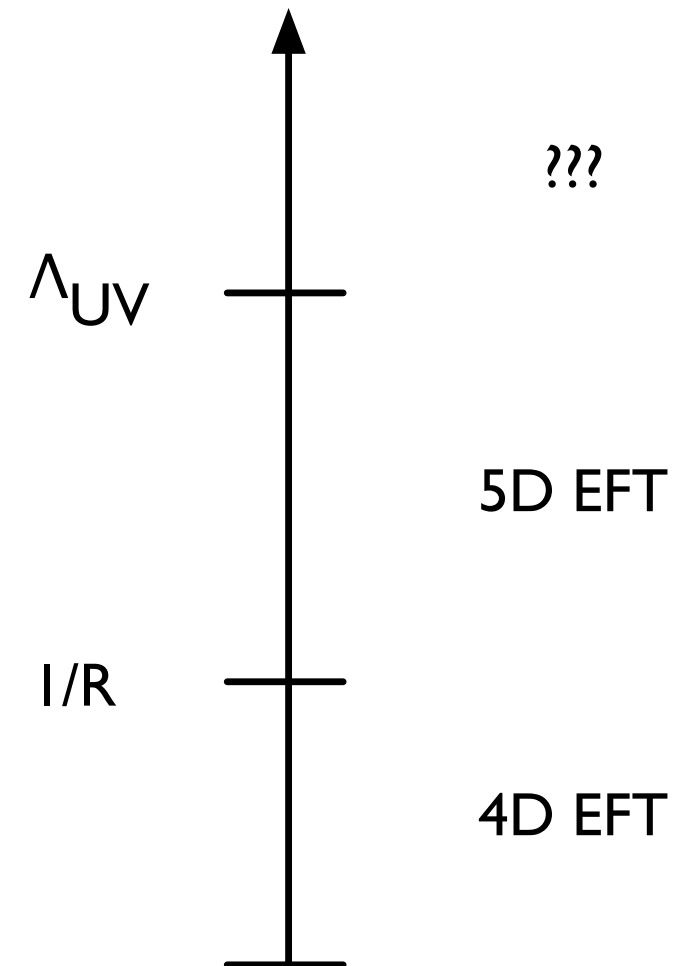
$$\hat{A}_\mu^a = \frac{1}{\sqrt{\pi R}} \left[ A_\mu^{a0}(x_\nu) + \sqrt{2} \sum_{n=1}^{\infty} A_\mu^{an}(x_\nu) \cos\left(\frac{nx_5}{R}\right) \right]$$

$$\hat{A}_5^a = \sqrt{\frac{2}{\pi R}} \sum_{n=1}^{\infty} A_5^{an}(x_\nu) \sin\left(\frac{nx_5}{R}\right)$$

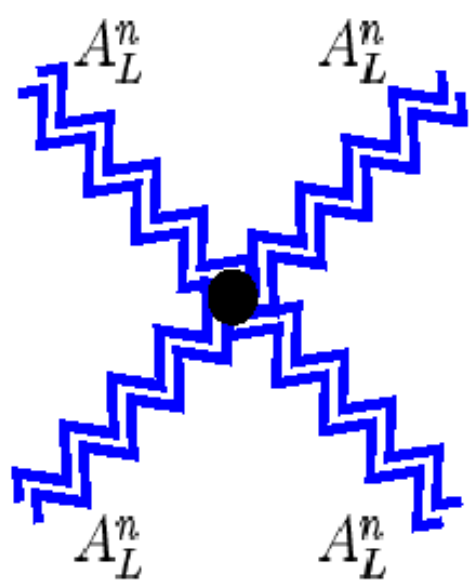
$$g_4 = \frac{g_5}{\sqrt{\pi R}}$$

$$M_n = \frac{n}{R}$$

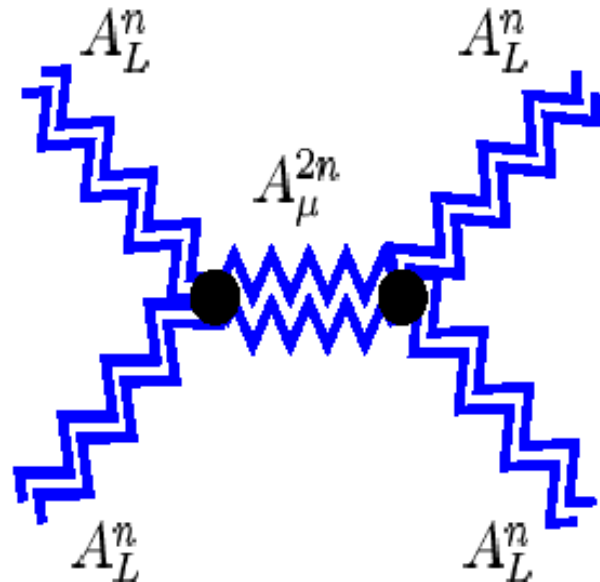
$$\Lambda_{UV} \propto \frac{1}{g_5^2}$$



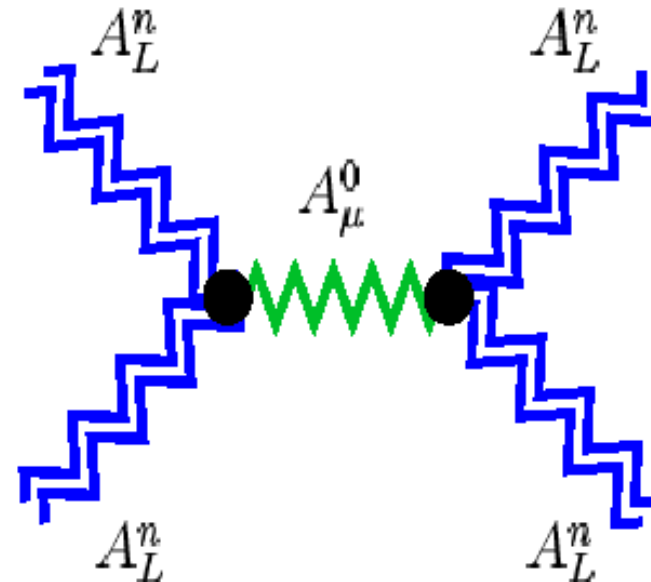
# 4-D KK MODE SCATTERING



(a)



(b1)



(c1)

+ Crossing Channels

(b2, b3) + (c2, c3)

**Cancellation of bad high-energy behavior through exchange of massive vector particles**

graph	$g^2 C^{eab} C^{ecd}$	$g^2 C^{eac} C^{edb}$	$g^2 C^{ead} C^{ebc}$
(a)	$6c(x^4 - x^2)$	$\frac{3}{2}(3 - 2c - c^2)x^4$ $-3(1 - c)x^2$	$\frac{-3}{2}(3 + 2c - c^2)x^4$ $+3(1 + c)x^2$
(b1)	$-2c(x^4 + x^2)$		
(c1)	$-4cx^4$		
(b2, 3)		$\frac{-1}{2}(3 - 2c + c^2)x^4$ $+3(1 - c)x^2$	$\frac{1}{2}(3 + 2c - c^2)x^4$ $-3(1 + c)x^2$
(c2, 3)		$(-3 + 2c + c^2)x^4$ $-8cx^2$	$(3 + 2c - c^2)x^4$ $-8cx^2$
<b>Sum</b>	$-8cx^2$	$-8cx^2$	$-8cx^2 \Rightarrow 0$

RSC, H.J. He, D. Dicus (2002)

Csaki, Grojean, Murayama, Pilo, Terning (2004)



# GENERAL PRINCIPLES

Higgsless models are low-energy effective theories of Dynamical Electroweak Symmetry Breaking with. They include:

- massive 4-d gauge bosons arise in the context of a 5-d gauge theory with appropriate boundary conditions
- $WW$  scattering is unitarized through exchange of KK modes (instead of scalar bosons)
- the language of Deconstruction allows a 4-d “Moose” representation of the model

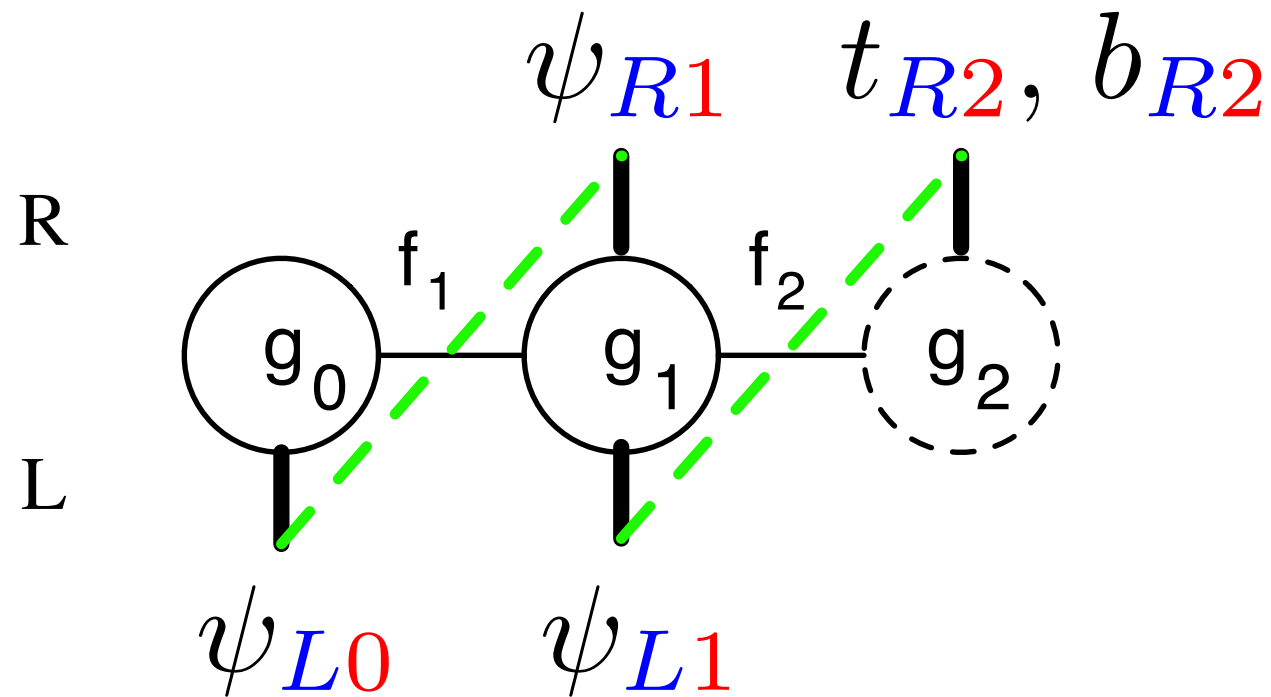
Chivukula & He hep-ph/0201164  
Csaki, Grojean, Murayama, Pilo, Terning hep-ph/0305237

# **A SIMPLE REALIZATION: THE THREE-SITE MODEL**

# 3-SITE MODEL: BASIC STRUCTURE

$$SU(2) \times SU(2) \times U(1)$$

$$g_0, g_2 \ll g_1$$



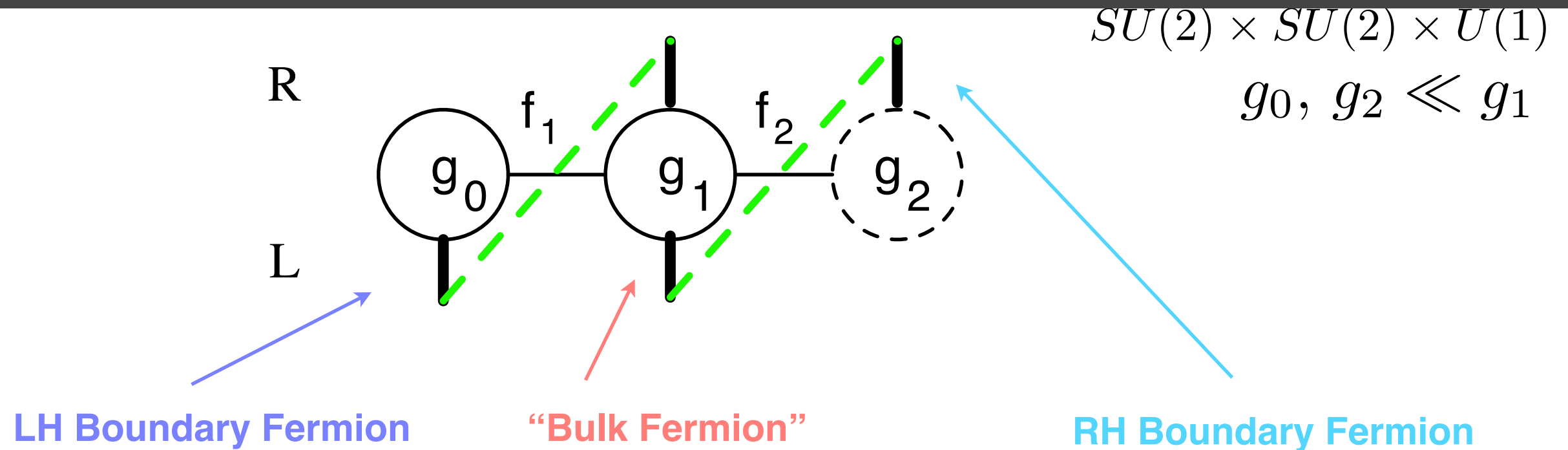
**Gauge boson spectrum:** photon, Z, Z', W, W' (as in BESS)

**Fermion spectrum:** t, T, b, B (  $\psi$  is an SU(2) doublet)

and also c,C, s,S, u,U, d,D plus the leptons

RSC, Coleppa, DiChiara, He, Kurachi, EHS, Tanabashi hep-ph/0607124

# 3-SITE FERMION MASSES



$$M \left[ \epsilon_L \bar{\psi}_{L0} \Sigma_{01} \psi_{R1} + \bar{\psi}_{R1} \psi_{L1} + \bar{\psi}_{L1} \Sigma_{12} \begin{pmatrix} \epsilon_{uR} & 0 \\ 0 & \epsilon_{dR} \end{pmatrix} \begin{pmatrix} u_{R2} \\ d_{R2} \end{pmatrix} \right]$$

degree of delocalization

ordinary fermion masses are of the form  $m_f \approx M \epsilon_L \epsilon_{fR}$   
 each ordinary fermion mass value is tied to  $\epsilon_{fR}$   
 flavor structure same as in standard model

heavy “KK” fermion masses are  $\sim M$

# 3-SITE IDEAL DELOCALIZATION

General **ideal delocalization** condition  $g_i(\psi_i^f)^2 = g_W v_i^w$

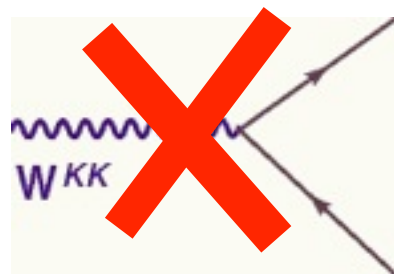
is realized as  $\frac{g_0(\psi_{L0}^f)^2}{g_1(\psi_{L1}^f)^2} = \frac{v_W^0}{v_W^1}$  in 3-site model

From the W, fermion eigenvectors, one solves for

$$\epsilon_L^2 \rightarrow (1 + \epsilon_{fR}^2)^2 \left[ \frac{x^2}{2} + \left( \frac{1}{8} - \frac{\epsilon_{fR}^2}{2} \right) x^4 + \dots \right] \quad x^2 \equiv \left( \frac{g_0}{g_1} \right)^2 \approx 4 \left( \frac{M_W}{M_{W'}} \right)^2$$

For all but top quark,  $\epsilon_{fR} \ll 1$  so the choice  $\epsilon_L^2 \approx 2 \left( \frac{M_W^2}{M_{W'}^2} \right)$

makes W' **fermiophobic** and Z' nearly so



$$\hat{S} = \hat{T} = W = 0$$

$$Y = M_W^2 (\Sigma_W - \Sigma_Z)$$

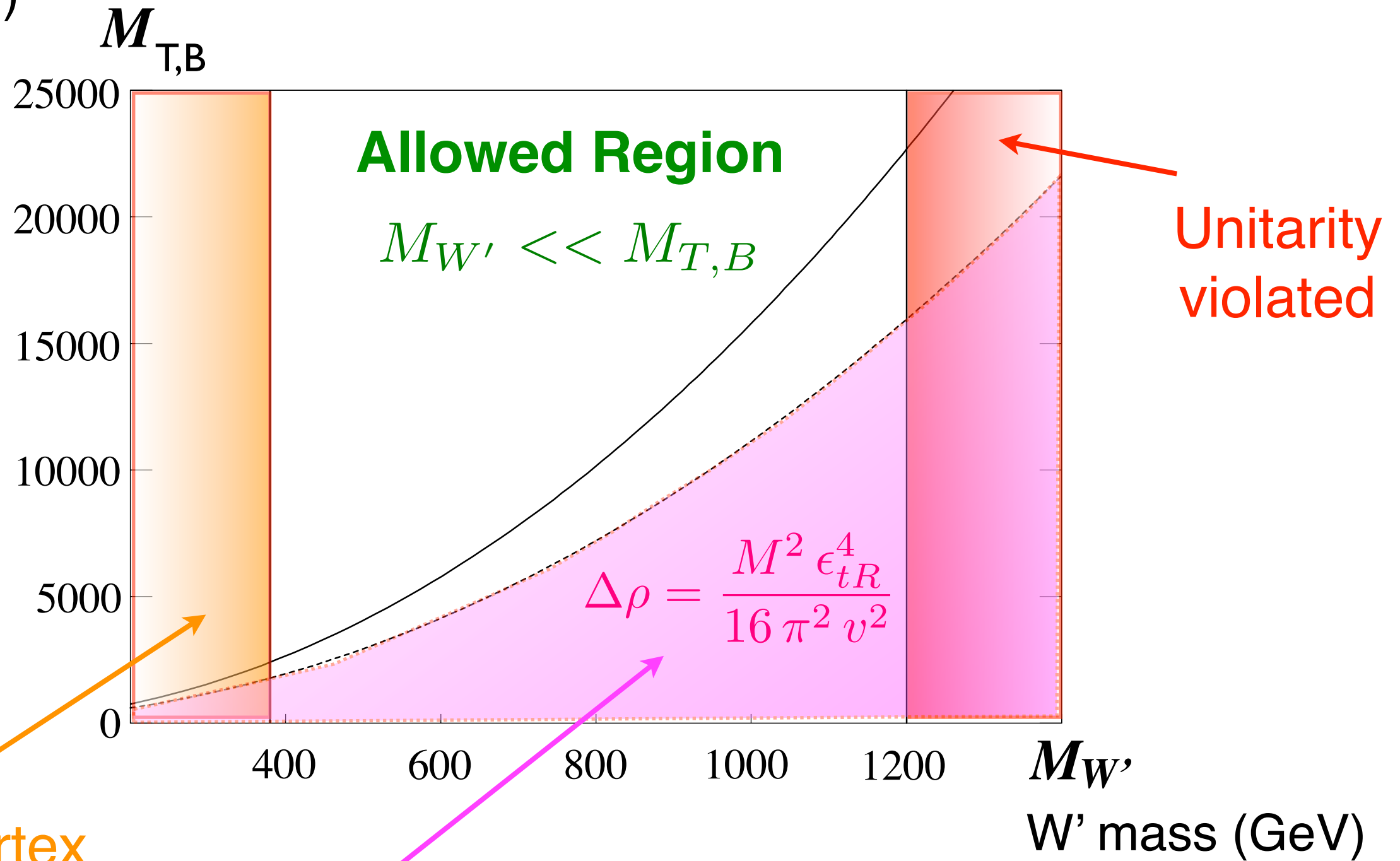
Use WW scattering to see W': Birkedal, Matchev, Perelstein hep-ph/0412278



# 3-SITE PARAMETER SPACE

Chivukula et al. hep-ph/0607124

KK fermion  
mass (GeV)



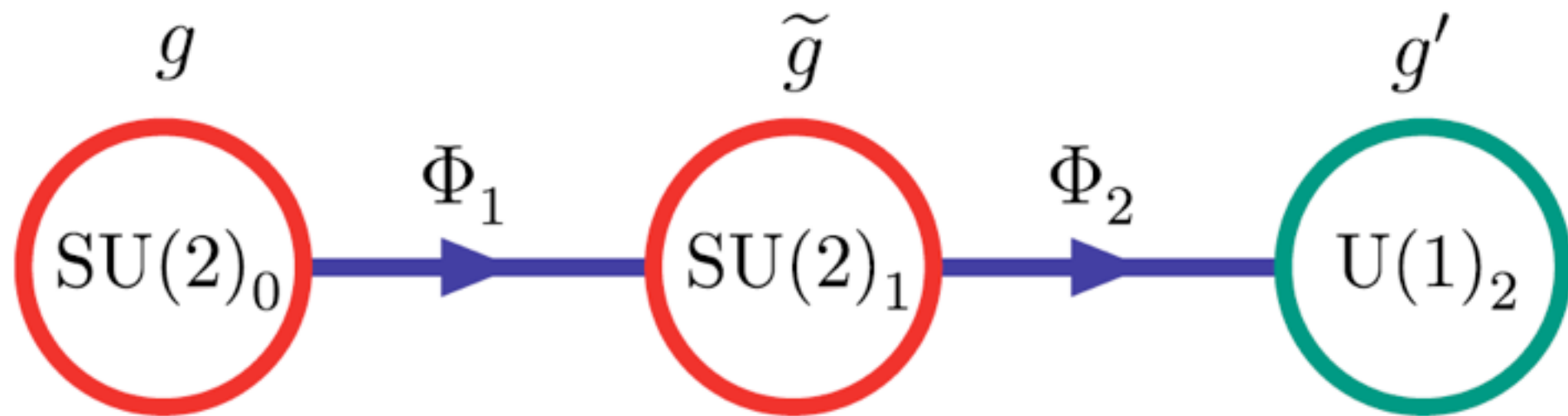
WWZ vertex  
visibly altered

1-loop fermionic EW  
precision corrections too large

# **BUT WHAT ABOUT THE NEW BOSON?**

**HONG-JIAN HE, NING CHEN, TOMOHIRO ABE: ARXIV 1207.4103**

# LINEAR 3-SITE MODEL

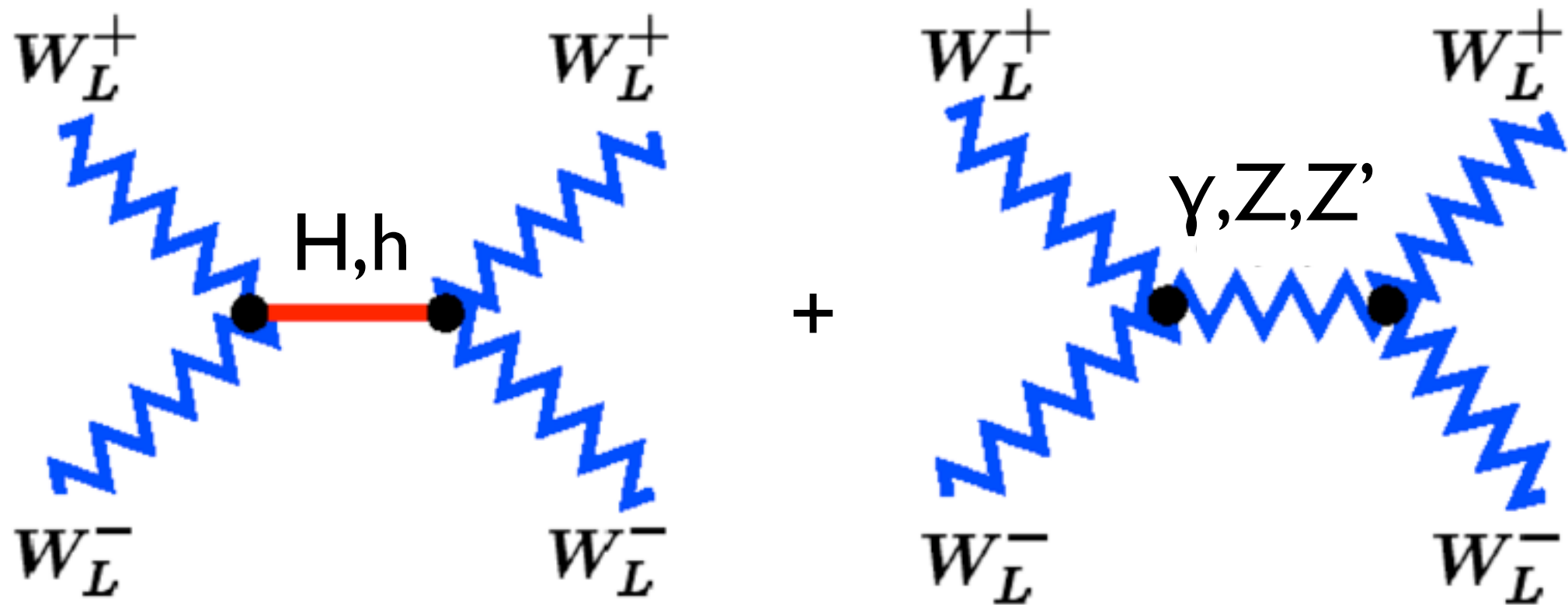


**Linear Scalar Link Fields:**  $\phi_1$  &  $\phi_2$

Leads to two-Higgs particles:  $h$ ,  $H$

He, Chen, Abe: arxiv 1207.4103

# WW UNITARIZATION



Unitarize *jointly* by scalar and vector exchange!  
Leads to sum rule:

$$G_{4W_0} - \frac{3M_{Z_0}^2}{4M_{W_0}^2} G_{W_0 W_0 Z_0}^2 = \sum_k \frac{3M_{Z_k}^2}{4M_{W_0}^2} G_{W_0 W_0 Z_k}^2 + \sum_k \frac{G_{W_0 W_0 h_k}^2}{4M_{W_0}^2}$$

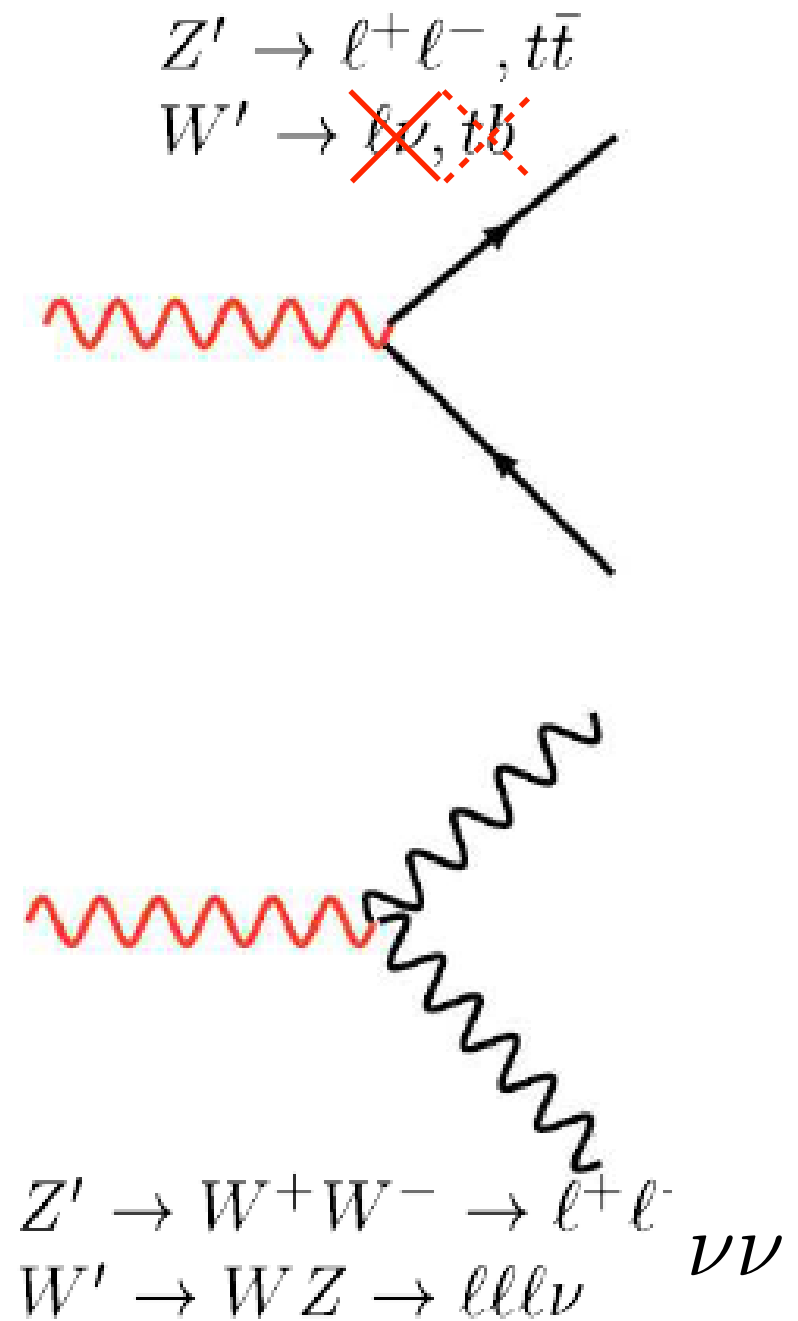
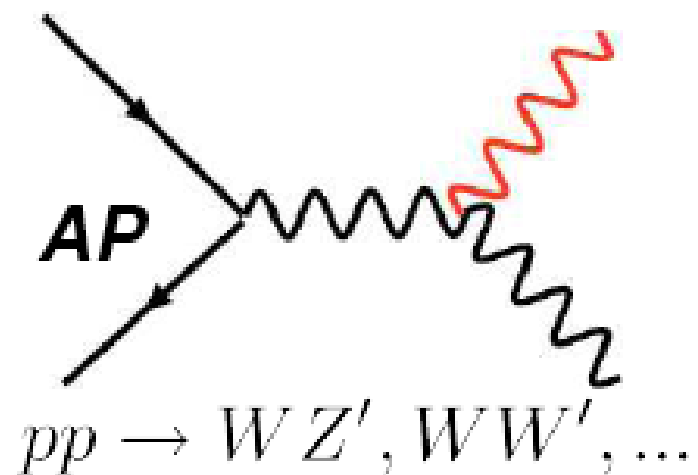
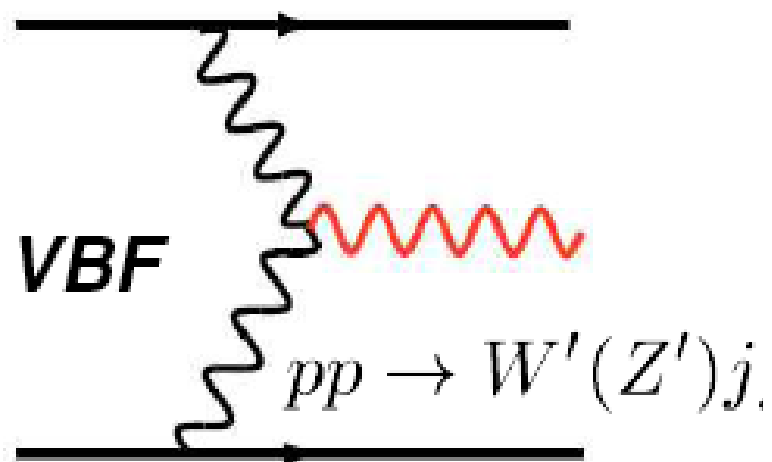
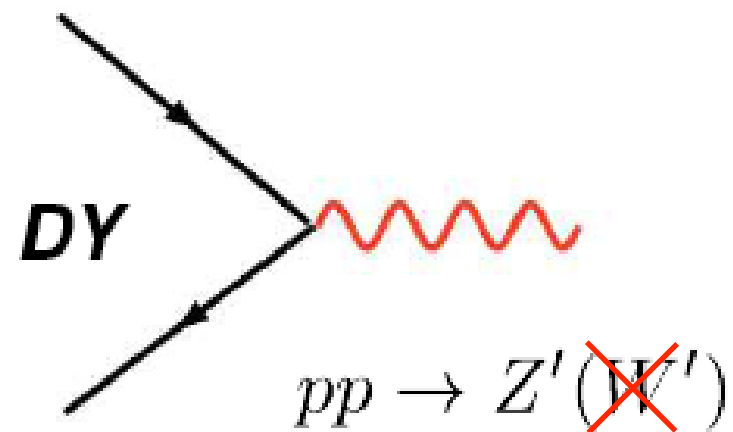
See next talk...

# LHC PHENOMENOLOGY OF VECTOR BOSONS

**RSC, EHS, H.-J. HE, Y.-P. KUANG, ET. AL., PHYS. REV. D78 (2008) 031701  
& ARXIV:1206.6022 AND PRD IN PRESS**



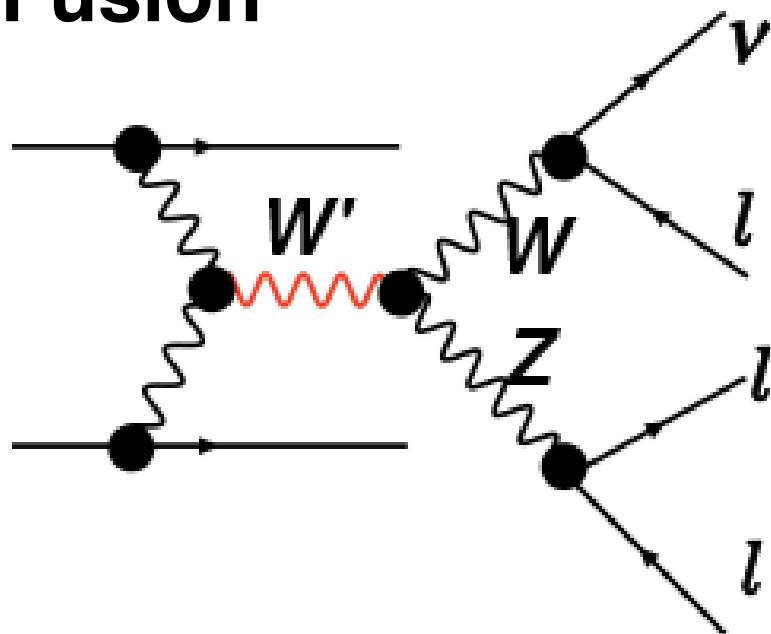
# W',Z' PRODUCTION AND DECAY AT LHC



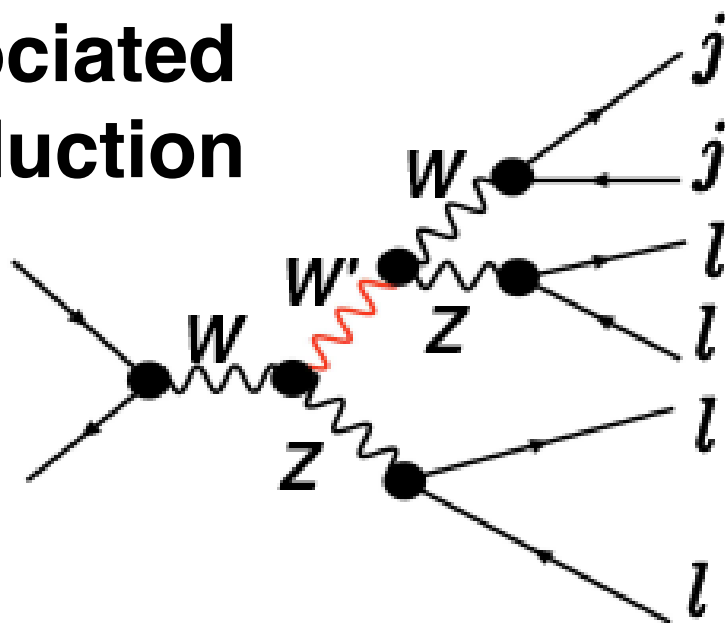
# W' PRODUCTION AT LHC

Two processes with large rates and clear signatures!

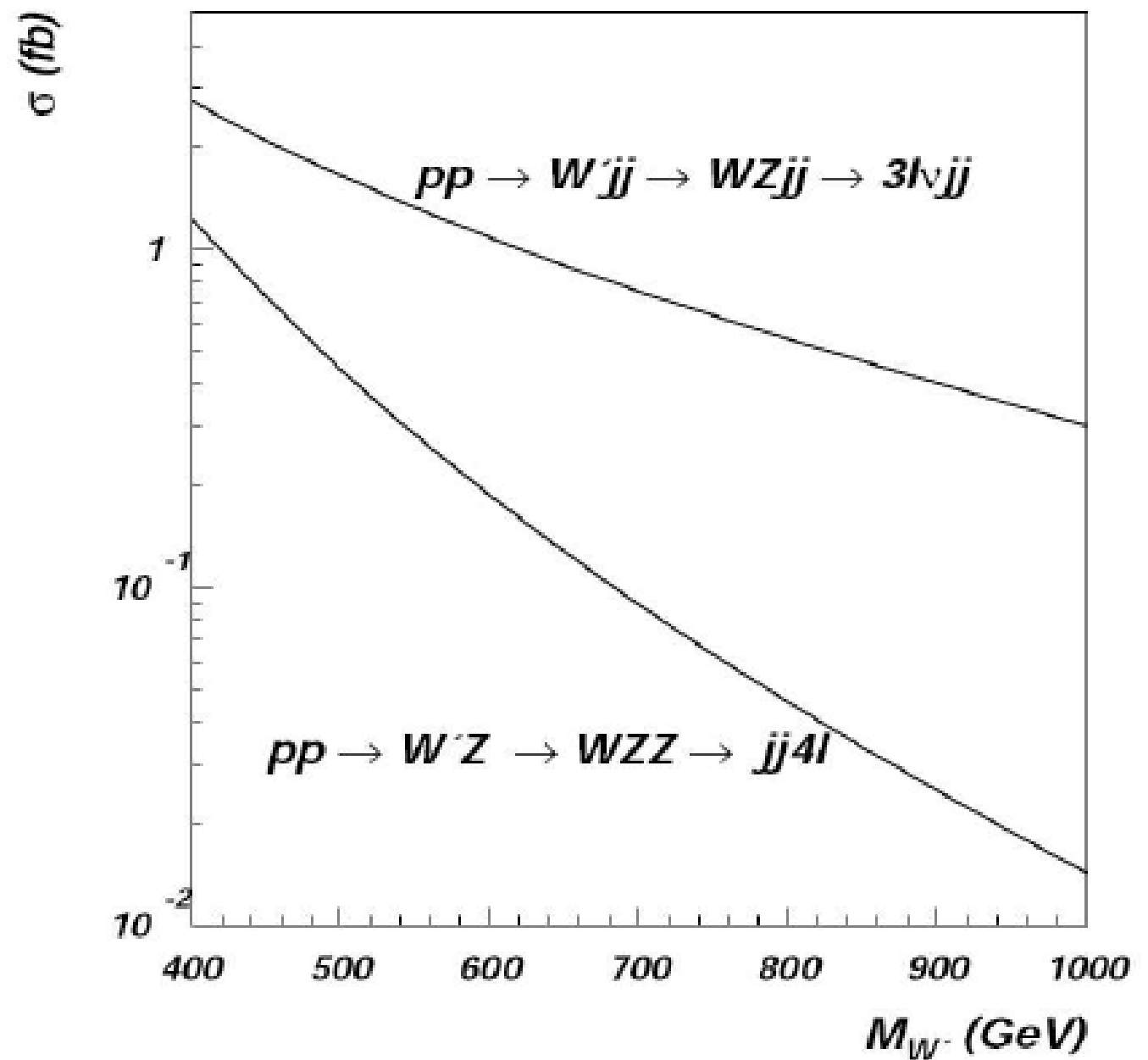
## Vector Boson Fusion



## Associated Production

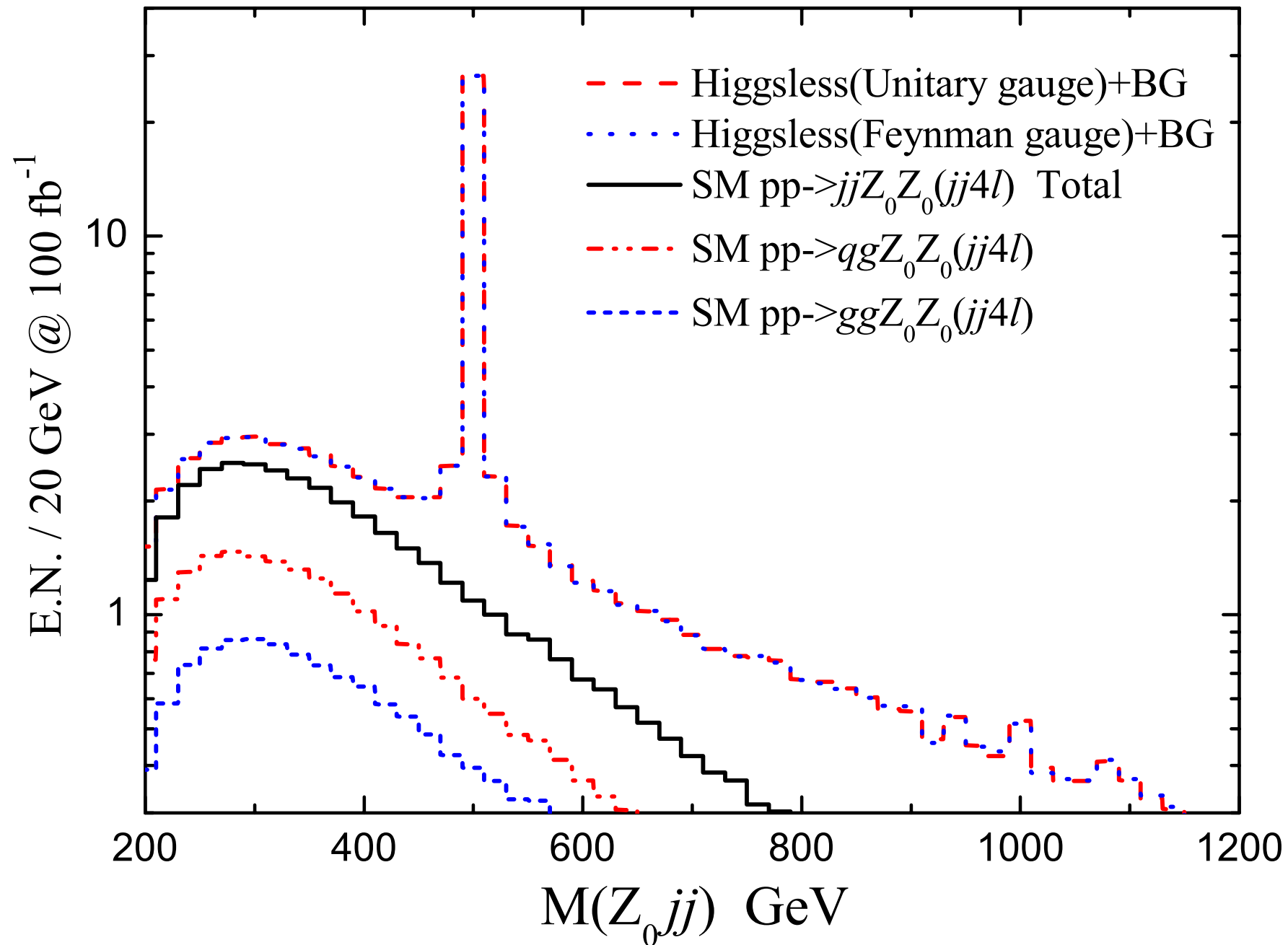


## LHC @14 TeV



# ASSOCIATED PRODUCTION (WZZ CHANNEL)

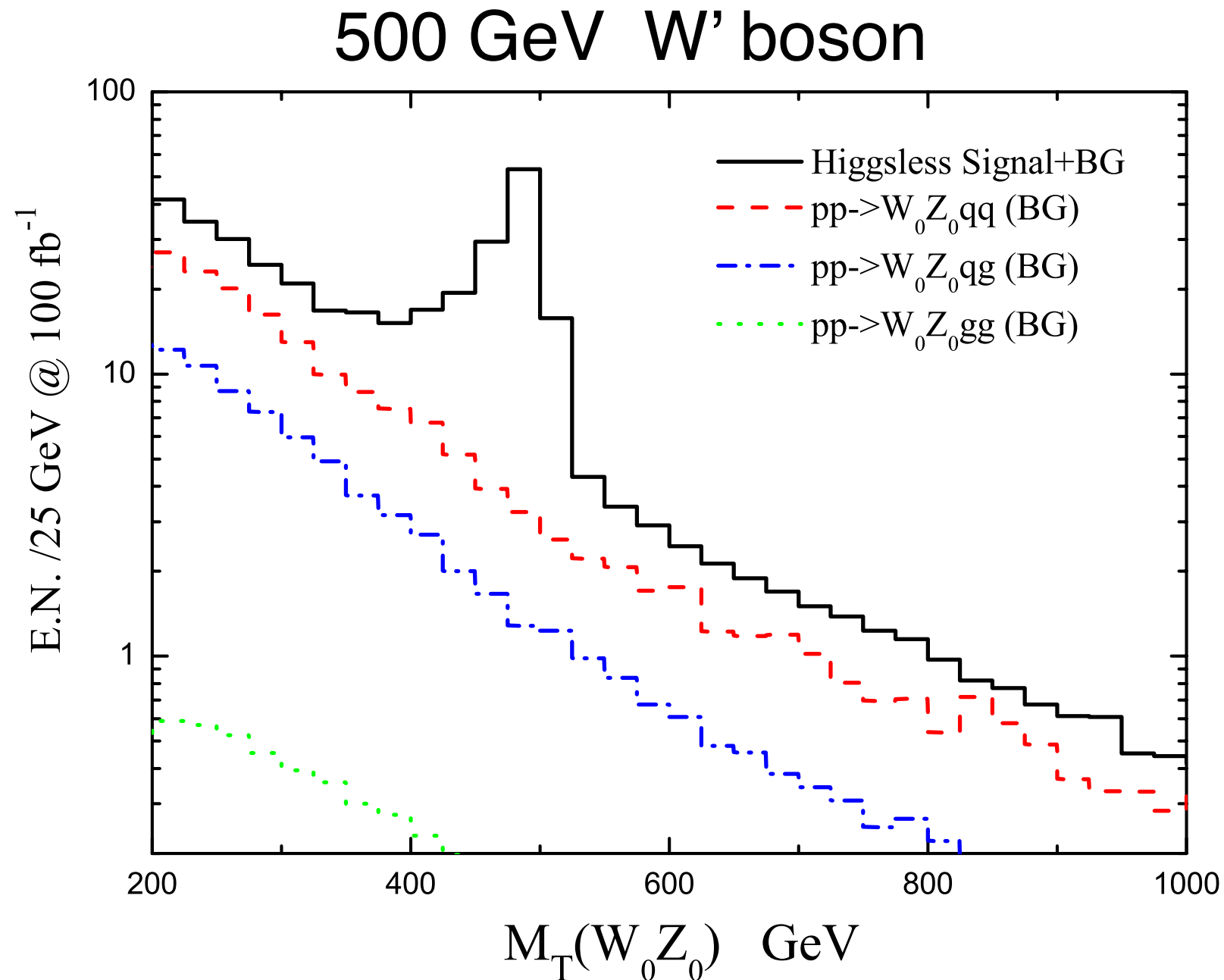
500 GeV  $W'$  boson



$$M_{jj} = 80 \pm 15 \text{ GeV}, \quad \Delta R(jj) < 1.5, \quad \sum_Z p_T(Z) + \sum_j p_T(j) = \pm 15 \text{ GeV}.$$

$$p_{T\ell} > 10 \text{ GeV}, \quad |\eta_\ell| < 2.5, \quad p_{Tj} > 15 \text{ GeV}, \quad |\eta_j| < 4.5.$$

# VECTOR BOSON FUSION (WZJJ CHANNEL)



Background is  
10x larger than  
estimated in  
Birkedal, Matchev  
& Perelstein (2005)

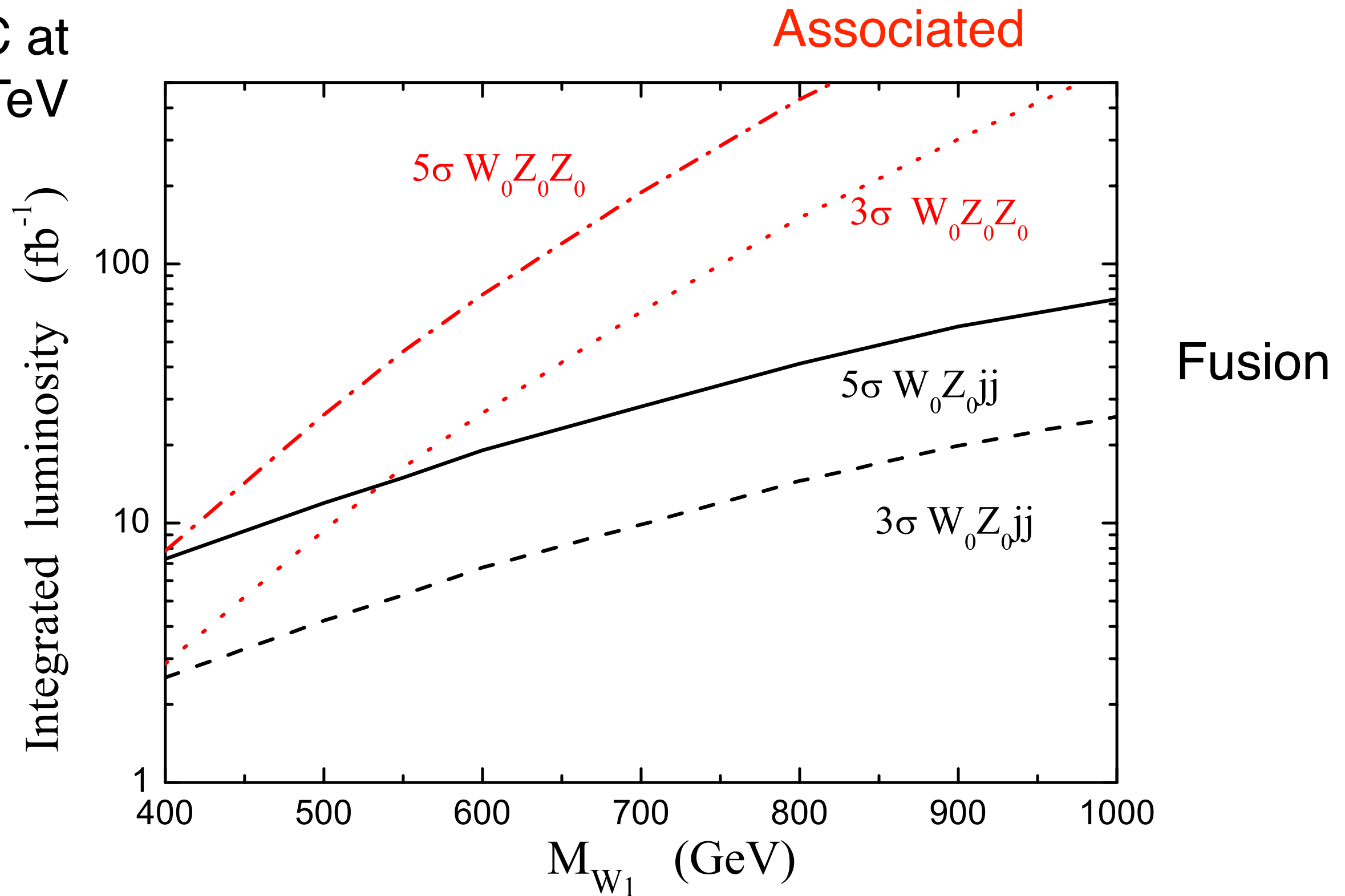
forward jet tag removes WZ background

$$E_j > 300 \text{ GeV}, \quad p_{Tj} > 30 \text{ GeV}, \quad |\eta_j| < 4.5, \quad |\Delta\eta_{jj}| > 4,$$

$$p_{T\ell} > 10 \text{ GeV}, \quad |\eta_\ell| < 2.5.$$

# INTEGRATED LUMINOSITY FOR $W'$ DISCOVERY

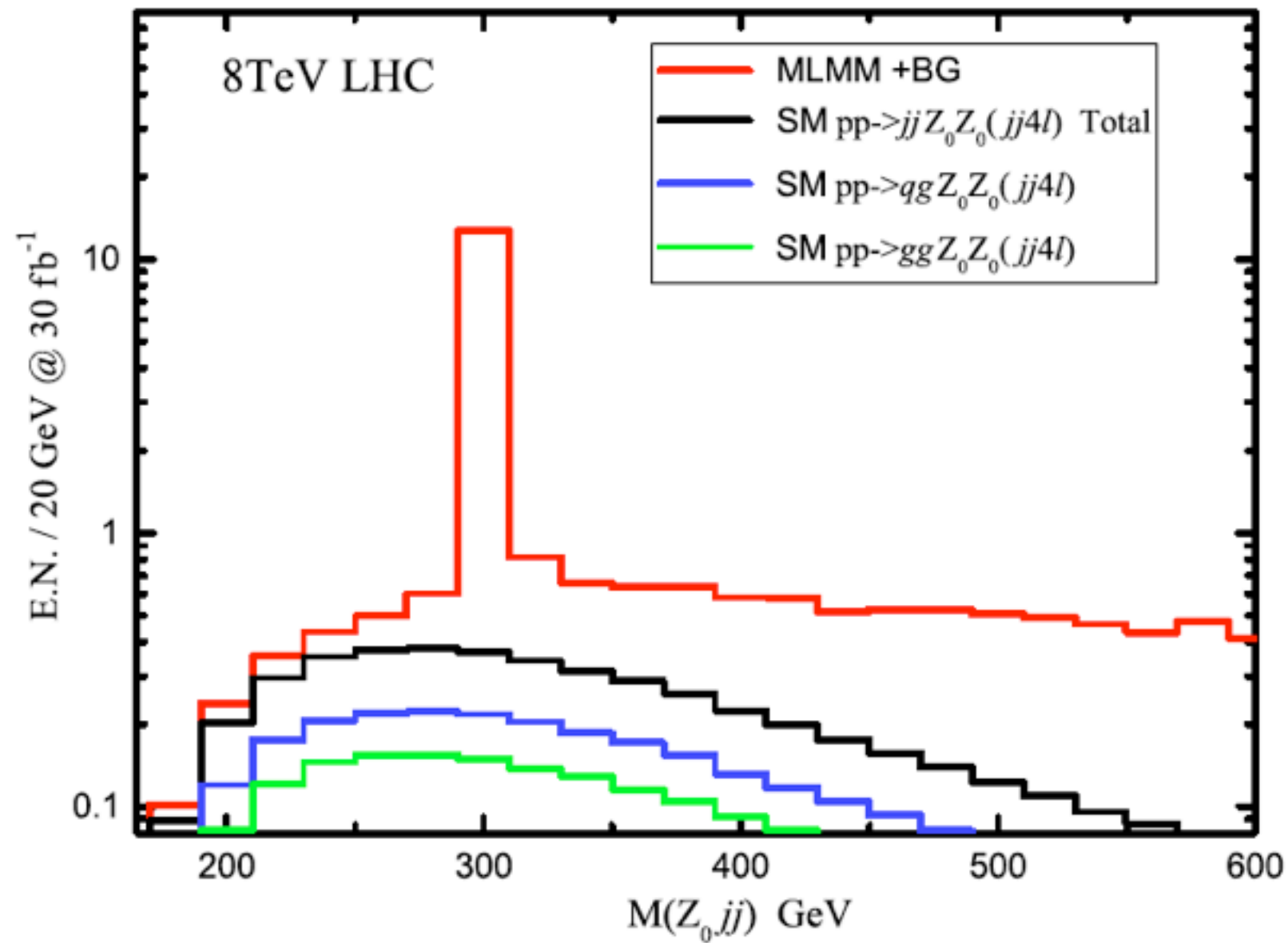
LHC at  
14 TeV





# W' DISCOVERY AT 8 TeV

## Associated Production



$$\begin{aligned} p_{T\ell} &> 10 \text{ GeV}, & |\eta_\ell| < 2.5, & M_{jj} = 80 \pm 15 \text{ GeV}. \\ p_{Tj} &> 15 \text{ GeV}, & |\eta_j| < 4.5, & \Delta R(jj) < 1.6. \end{aligned}$$

# W' PRODUCTION AT 8 TeV

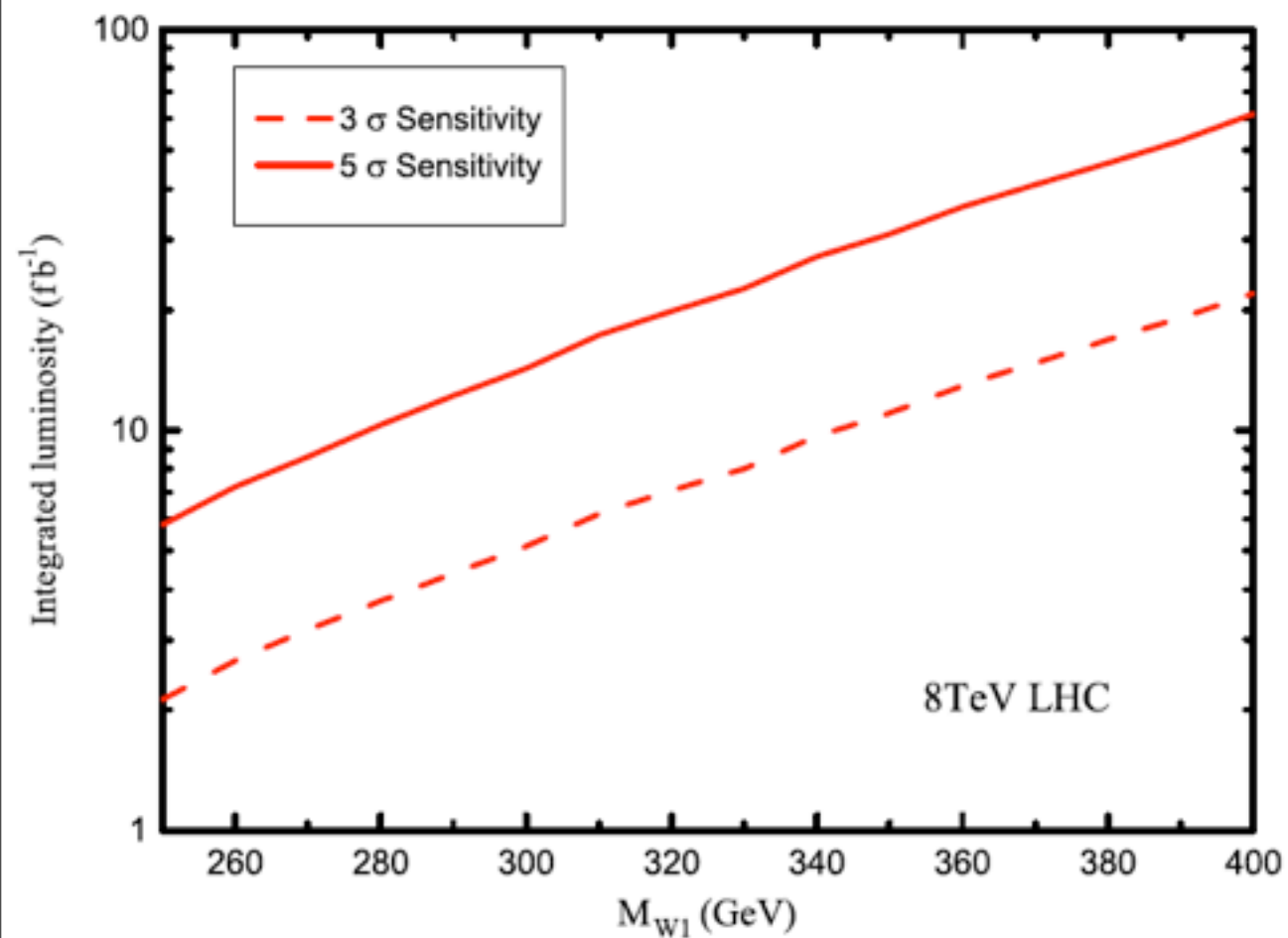


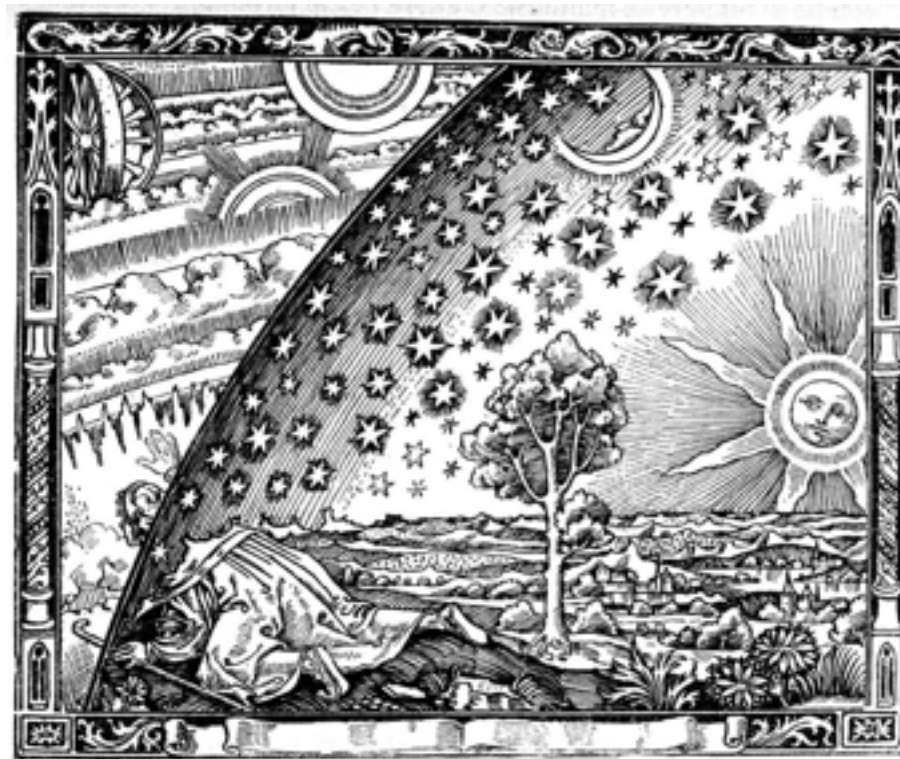
TABLE II. The  $5\sigma$  discovery reaches of the  $W_1^\pm$  bosons at the LHC-8, with the integrated luminosities  $\int \mathcal{L} = 10, 15, 20, 25, 30, 35, 40, 50, 60 \text{ fb}^{-1}$ , respectively.

$\int \mathcal{L} \text{ (fb}^{-1}\text{)}$	$M_{W_1} \text{ (GeV)}$
10	277
15	302
20	320
25	335
30	346
35	357
40	367
50	385
60	397

arxiv:1206.6022, PRD in press

# CONCLUSION

- ATLAS/CMS has discovered a **new** boson.
- Measure properties: is it **the** Higgs?
- If it isn't, there are potential new signatures in multi-gauge boson signals.



# MSU, TSINGHUA, AND PROF. KUANG

- C.-P. Yuan met Prof. Kuang at CCAST Workshop, Beijing in 1993.
- E. Simmons invited by Prof. Kuang to speak at ITP in Beijing in 1995.
- NSF USA-China International Program, Qing Wang and Yi Liao visited MSU, Carl Schmidt visited Tsinghua U, 1997-99.
- Hong-Jian He was a postdoc at MSU, 1997-2000.
- Chivukula, Dicus, & He on unitarity, 2001.
- Kuang et. al., LHC signature paper 2007...

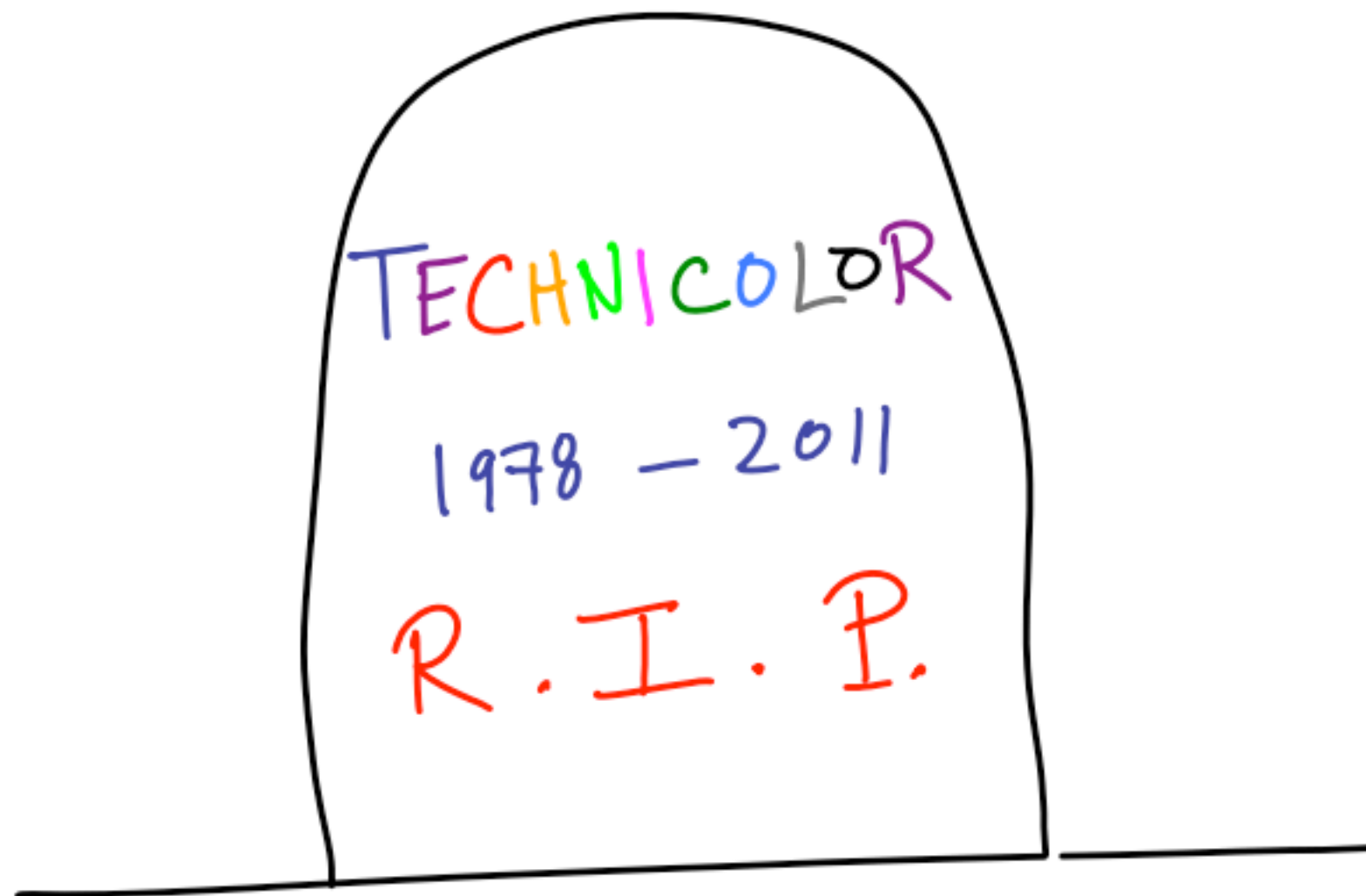
# 生日快乐!



# BACKUP SLIDES

# DISCUSSION QUESTION

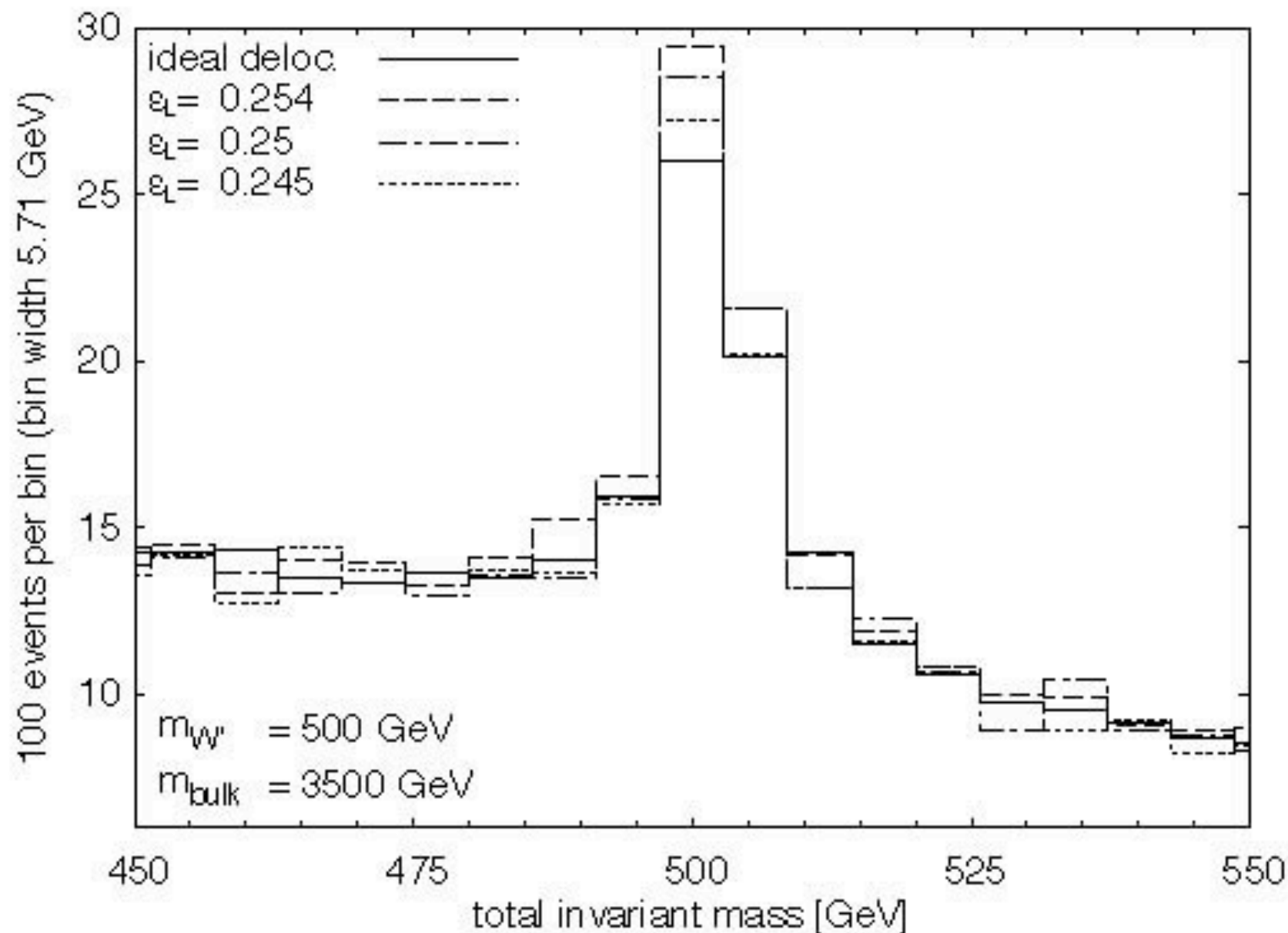
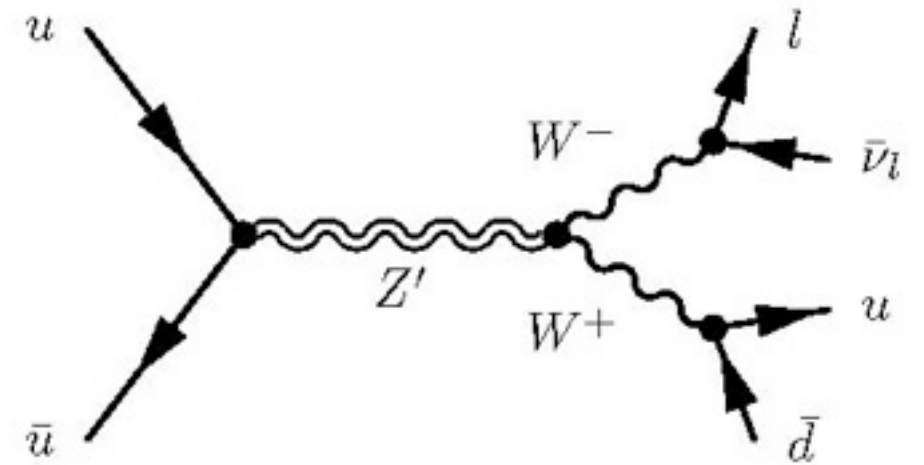
Obvious Implications of  $m_H \sim 125 \text{ GeV}$  ?





# Z' SEARCH AT LHC

Ohl & Speckner predict that the 3-site Z' boson (at or near ideal delocalization) should be visible in 100 fb<sup>-1</sup> of LHC data



$$p_T \geq 50 \text{ GeV}$$

$$|\cos \theta| \leq 0.95$$

$$75 \text{ GeV} \leq m_{jj} \leq 85 \text{ GeV}$$

$$M_{W'} = 500 \text{ GeV}$$

Ohl & Speckner arXiv:0809.0023