## Inflation in Light of Planck & BICEP2

**Big Bang** 

What happened here?



# Some Big Questions

- Why is the Universe so big and old?
- Why is it (almost) homogeneous on large scales?
- Why is its geometry (almost) Euclidean?
- What is the origin of structures in the Universe? A possible answer:
   Cosmological inflation
   What do the Planck and PICED2 data tell us?
- What do the Planck and BICEP2 data tell us?

Models of inflation

## General Relativity & Cosmology

• Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R + g_{\mu\nu} \Lambda = \frac{8\pi G}{c^4} T_{\mu\nu}$$

- G = Newton constant, energy scale ~  $10^{19}$  GeV, Distance ~  $10^{-32}$  cm, time ~  $10^{-43}$  s
- Our Universe is a special solution of equations Size >>  $10^{-32}$  cm, age ~  $10^{-43}$  s
- Due to a "cosmological constant" Λ?
  Due to a scalar "inflaton" field φ: Λ ~ V(φ) ≠ 0?

# **Cosmological Inflation**

Expansion driven by cosmological constant:

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{\Lambda}{3}$$

- Exponential expansion
- All the visible Universe was once very small
- In close contact: (almost) homogeneous
- Geometry ~ flat

$$(a = scale size)$$



## **Primordial Perturbations**

- "Cosmological constant" due to vacuum energy in
   "inflaton" field φ: Λ ~ V(φ) ≠ 0
- Quantum fluctuations in φ cause perturbations in energy density (scalar) and metric (tensor)
- (Almost) independent of scale size



• Visible in cosmic microwave background (CMB)

# Perturbations Generate Structures

Planck

Background



Dark matter distribution today (simulated)

You! Millennium Run 77.696.000 particles

### Inflationary Perturbations in a Nutshell

- Universe expanding like de Sitter
- Horizon  $\rightarrow$  information loss  $\rightarrow$  mixed state
- Effective 'Hawking temperature'  $H/2\pi$
- Expect quantum fluctuations:

 $\langle \phi \phi \rangle$ ,  $\langle hh \rangle \sim H^2$ (inflaton, gravitational field:  $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ )

Scalar (density), tensor perturbations

- Should have thermal spectrum
- Scalars enhanced by evolution of  $\phi \rightarrow r \sim \epsilon$

## **Slow-Roll Inflation**

• Expansion driven by cosmological constant:

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{\Lambda}{3}$$

- Getting small density perturbations requires a "small" potential:  $\left(\frac{V}{L}\right)^{\frac{1}{4}} = 0.0275 \times M_{Pl}$
- That is almost flat:  $\epsilon = \frac{1}{2}M_{Pl}^2 \left(\frac{V'}{V}\right)^2$ ,  $\eta = M_{Pl}^2 \left(\frac{V''}{V}\right)$  small so as to get sufficient e-folds of expansion:

$$N = \frac{v^2}{M_{Pl}^2} \int_{x_i}^{x_e} \left(\frac{V}{V'}\right) dx$$

- Main observables of scalar and tensor pert'ns:
- Scalar tilt:  $n_s = 1 6\epsilon + 2\eta$ , tensor/scalar ratio:  $r = 16\epsilon$
- Perturbations ~ Gaussian: look for deviation  $(f_{NL})$



## The Spectrum of Fluctuations in the Cosmic Microwave Background



### Cosmological Inflation in Light of BICEP2



### **Quantum gravity radiation in the sky?**

## Cosmological Inflation in Light of BICEP2

- QUANTUM GRAVITY RADIATION!
- BICEP2 observations support for inflation

• "Large" tensor-to-scalar ratio:  $r = 0.20^{+0}_{-0}$ 





# The BICEP2 Result is a Big Deal



- Gravitational waves! Probably quantum origin
- Need to maintain proper scientific scepticism





 Radio-emitting loops from supernova bubbles



#### morint on CMR temperature mans

- Leave imprint on CMB temperature maps
- Associated with magnetized dust?
- May polarize CMBBICEP2 towards a loop





# Reconciling Planck and BICEP2?

- Possible if scalar spectral index runs (fast)
  - $\frac{dn_s}{d\ln k} = \frac{2}{3} \left[ (n_s 1)^2 4\eta^2 \right] + 2\xi \text{ where } \xi = M_{Pl}^4 \left( \frac{V'V'''}{V^2} \right)$
- Bracketed terms O(10<sup>-4</sup>), need  $\xi = O(10^{-2})$
- In slow-roll approximation, expect  $\xi \ll O(10^{-2})$
- Such a large value of ξ is problematic for consistency of slow-roll approximation:
- $\varepsilon = O(10^{-2})$  implies V'/V = O(10^{-1})
- $\xi = O(10^{-2})$  then requires V'''=  $O(10^{-1})$
- Change  $\Delta V'' = O(1)$  for  $\Delta \phi = O(10)$

### Inflationary Models in Light of Planck

- Planck CMB observations consistent with inflation
- Tilted scalar perturbation spin  $n_s = 0.9585$
- **BUT** strengthen upper limi perturbations: r < 0.10
- Challenge for simple inflationary models
- Starobinsky R<sup>2</sup> to rescue?
  - Higgs/Supersymmetry/su



# Starobinsky Model

- Non-minimal general relativity (singularity-free cosmology):
  No scalar!?
- Inflationary interpretation, calculation of perturbations: Mukhanov & Chibisov, 1981

$$\delta S_b = \frac{1}{2} \int d^4 x \left[ \phi'^2 - \nabla_a \phi \nabla^a \phi + \left( \frac{a}{a} + M^2 a^2 \right) \phi^2 \right]$$

Conformally equivalent to scalar field model:

$$S = \frac{1}{2} \int d^4x \sqrt{-\tilde{g}} \left[ \tilde{R} + (\partial_\mu \varphi')^2 - \frac{3}{2} M^2 (1 - e^{-\sqrt{2/3}\varphi'})^2 \right]$$

# Higgs Inflation: a Single Scalar?

Bezrukov & Shaposhnikov, arXiv:0710.3755

• Standard Model with non-minimal coupling to gravity:  $S_J = \int d^4x \sqrt{-g} \left\{ -\frac{M^2 + \xi h^2}{2} R \right\}$ 

Consider case 
$$1 \ll \sqrt{\xi} \ll 10^{17}$$
 : in Einstein frame

 $+\frac{\partial_{\mu}h\partial^{\mu}h}{2}-\frac{\lambda}{4}\left(h^{2}-v^{2}\right)^{2}$ 

$$S_E = \int d^4x \sqrt{-\hat{g}} \left\{ -\frac{M_P^2}{2}\hat{R} + \frac{\partial_\mu \chi \partial^\mu \chi}{2} - U(\chi) \right\}$$

- With potential:  $U(\chi) = \frac{\lambda M_P^4}{4\xi^2} \left(1 + \exp\left(-\frac{2\chi}{\sqrt{6}M_P}\right)\right)^{-2}$ Similar to Starobinsky, but not identical
- Successful inflationary potential at  $\chi \gg M_P$

# Higgs Inflation: a Single Scalar?



• BUT:

Bezrukov & Shaposhnikov, arXiv:0710.3755

- Requires  $\lambda > 0$  beyond M<sub>P</sub>: need M<sub>H</sub> > 127 GeV?
- How natural is such a large value of  $\xi$ ?

### Theoretical Constraints on Higgs Mass

- Large  $M_h \rightarrow$  large self-coupling  $\rightarrow$  blow up at  $\lambda(Q) = \lambda(v) - \frac{3m_t^4}{2\pi^2 v^4} \log \frac{Q}{v} \quad \int_{0.08}^{0.10} \left[ \int_{0.08} \text{Instability @} \right]$
- Small: renormalization due to t quark drives quartic coupling < 0 at some scale Λ
   → vacuum unstable



• Vacuum could be stabilized by **Supersymmetry** 

Degrassi, Di Vita, Elias-Miro, Giudice, Isodori & Strumia, arXiv:1205.6497

### Vacuum Instability in the Standard Model

 $V[\mathcal{H}]$ 





Instability scale: Buttazzo, Degrassi, Giardino, Giudice, Sala, Salvio & Strumia, arXiv:1307.3536

$$\log_{10} \frac{\Lambda_I}{\text{GeV}} = 11.3 + 1.0 \left( \frac{M_h}{\text{GeV}} - 125.66 \right) - 1.2 \left( \frac{M_t}{\text{GeV}} - 173.10 \right) + 0.4 \frac{\alpha_3(M_Z) - 0.1184}{0.0007}$$
New measurement of m<sub>4</sub> = 173.34 ± 0.76 GeV

### Inflationary Models in Light of BICEP2

- Back to basics: monomial power-law inflationary potential?  $V = \mu^{4-n}\phi^n$
- Slow-roll parameters:  $\epsilon = \frac{n^2}{2} \frac{M_{Pl}^2}{\phi^2}; \eta = n(n-1) \frac{M_{Pl}^2}{\phi^2}$
- Inflationary observables:  $r = 8n^2 \frac{M_{Pl}^2}{\phi^2}; n_s = 1 n(n+2) \frac{M_{Pl}^2}{\phi^2}$
- Number of e-folds:  $N = \frac{1}{2n} \frac{\phi^2}{M_{\text{Pl}}^2}$
- **Consistency condition:**  $r = 8\left(1 n_s \frac{1}{N}\right)$
- Comfortably satisfied by the measurements  $r=0.16^{+0.06}_{-0.05}, n_s=0.960\pm0.008, N=50\pm10$

JE, Garc n, Nanopoulos & Olive, arXiv:1403.7518

### Inflationary Models in Light of BICEP2

- What is the power in the monomial inflationary potential?  $V = \mu^{4-n}\phi^n$
- Two expressions for n:  $n = \frac{rN}{4}$ ;  $n = 2[N(1-n_s)-1]$
- Numerical values from measurements:

$$n = 2.0^{+0.9}_{-0.8}; n = 2.0 \pm 1.1$$

- n = 2 perfectly OK, n = 1, 3 potentials unbounded below, n = 4 has  $\Delta \chi^2 > 8$
- Looks like quadratic inflationary potential  $V = \mu^2 \phi^2 \simeq (2 \times 10^{16} \text{ GeV})^4 \ m = \sqrt{2}\mu = 1.8 \times 10^{13} \text{ GeV} \ \phi = \sqrt{200} M_{Pl}$

# Destabilization of Electroweak Vacuum during Inflation

Cobakhidze & Spencer-Smith: arXiv:1301.2846

- Transition out of unstable electroweak vacuum if  $H > 10^{10}$  to  $10^{12}$  GeV
- BICEP2 measurements corresponds to

 $V \sim (2 \times 10^{16} \,\text{GeV})^4$ 

• Corresponding to

 $H \sim 10^{14} \text{ GeV}$ 

 Strengthens case for new physics beyond the Standard Model, scale < 10<sup>11</sup> GeV

### Inflation Cries out for Supersymmetry

- Want "elementary" scalar field
  - (at least looks elementary at energies  $\langle M_P \rangle$ )
- To get right magnitude of perturbations
- Prefer mass << M<sub>P</sub>
  - (~  $10^{13}$  GeV in simple  $\varphi^2$  models)
- And/or prefer small self-coupling  $\lambda \ll 1$
- Both technically natural with supersymmetry

E, Nanopoulos, Olive, & Tamvakis: 1983

## **Sneutrino Inflation?**

- See-saw model of neutrino masses requires massive singlet (right-handed) neutrinos
- In supersymmetric models would be accompanied by massive sneutrinos, with  $m^2\phi^2$  potential
- Decay of sneutrino inflaton gives efficient leptogenesis
  BUT: disfavoured by

Planck data



## Multiple Sneutrino Inflation?

E, Fairbairn & Sueiro: arXiv:1312.1353

- Expect 3 massive singlet (right-handed) neutrinos
- Consistency with Planck data?
- Marginal improvement with 2 sneutrinos



Significant improvement with 3 sneutrinos

# Multiple Sneutrino Inflation?

JE. Fairbairn & Sueiro: arXiv:1312.1353

• Sneutrino inflation model compatible with all the neutrino and Planck data

 $\sin^2(2\theta_{13})$ 

0.2

0.4

 $\sin^2(2\theta_{\rm s})$ 

0.6

 $\sin^2(2\theta_{12})$ 

0.8

• Neutrino masses and mixing

• Full set of model parameters





# More General Polynomial Potential

- Consider single real field with double-well potential:  $V = A\phi^2(v \phi)^2$
- Shallower than  $\phi^2$  for

 $0 < \phi < v$ 

- Better tensor-to-scalar ratio r for  $0 < \phi < v$
- Steeper than  $\phi^2$  for  $\phi < 0$  or > v: worse r



Croon, JE & Mavromatos: arXiv:1303.6253

### Effective Potential in Wess-Zumino Model

• First (and minimal) supersymmetric field theory:

 $W = \frac{\mu}{2}\Phi^2 - \frac{\lambda}{3}\Phi^3$ 

- Effective potential:  $V = \left|\frac{\partial W}{\partial \phi}\right|^2 = Av^4(x^4 2\cos\theta x^3 + x^2)$
- Equivalent to single-field model for  $\theta = 0$  (good)
- Combination of  $\varphi^2 + \varphi^4$ for  $\theta = \pi/2$  (no good)
- Good inflation possible
   Inflaton could be sneutrino



Croon, JE & Mavromatos: arXiv:1303.6253

### Supersymmetric Inflation in Light of Planck



## **Exploring Wess-Zumino Inflation**

- Generic SUSY models have 2 scalars:  $\Phi = \psi + i \sigma$
- Explore 2-field parameter space of WZ model



• Different trajectories for different v

JE, Mavromatos & Mulryne: arXiv:1401.6078

# **Exploring Wess-Zumino Inflation**

### • Explore 2-field parameter space of WZ model



# **Exploring Wess-Zumino Inflation**

 Statistical distribution of possible predictions assuming rotationally-invariant distribution of initial conditions
 JE, Mayromatos & Mulryne: arXiv:1401.6078





• Planck-compatibility not implausible

# From Supersymmetry to Supergravity

- The only good symmetry is a local symmetry (cf, gauge symmetry in Standard Model)
- Local supersymmetry = supergravity
- Early Universe cosmology needs gravity
- **Supersymmetry** + **gravity** = **supergravity**
- Superpartner of graviton is gravitino fermion
- Gravitino condensation? (cf, quarks in QCD)
- Mechanism for inflation? JE & Mavromatos, arXiv:1308.1906

- Supersymmetry + gravity = Supergravity
- Include conventional matter?
- Potentials in generic supergravity models have 'holes' with depths  $\sim -\,M_P{}^4$
- Exception: no-scale supergravity
- Appears in compactifications of string
- Flat directions, scalar potential ~ global model + controlled corrections JE, Nanopoulos & Olive, arXiv:1305.1247, 1307.3537

- Simplest SU(2,1)/U(1) example:
- K ähler potential:  $K = -3\ln(T + T^* |\phi|^2/3)$
- Wess-Zumino superpotential:  $W = \frac{\mu}{2}\Phi^2 \frac{\lambda}{3}\Phi^3$
- Assume modulus T = c/2 fixed by 'string dynamics'

• Ef 
$$\mathcal{L}_{eff} = \frac{c}{(c-|\phi|^2/3)^2} |\partial_\mu \phi|^2 - \frac{\hat{V}}{(c-|\phi|^2/3)^2}$$
  $\hat{V} \equiv \left|\frac{\partial W}{\partial \phi}\right|^2$ 

 Modifications to globally supersymmetric case JE, Nanopoulos & Olive, arXiv:1305.1247
 Good inflation possible ...

• Inflationary potential for  $\lambda \simeq \mu/3$ 



JE, Nanopoulos & Olive, arXiv:1305.1247

# Is there more profound connection?

• Starobinsky model:

$$S = \frac{1}{2} \int d^4x \sqrt{-g} (R + R^2/6M^2)$$

• After conformal transformation:

$$S = \frac{1}{2} \int d^4x \sqrt{-\tilde{g}} \left[ \tilde{R} + (\partial_\mu \varphi')^2 - \frac{3}{2} M^2 (1 - e^{-\sqrt{2/3}\varphi'})^2 \right]$$

- Effective potential:  $V = \frac{3}{4}M^2(1 e^{-\sqrt{2/3}\varphi'})^2$
- Identical with the no-scale Wess-Zumino model for the case  $\lambda = \mu/3$

... it actually IS Starobinsky

### See also ...

- Nakayama, Takahashi & Yanagida arXiv:1305.5099
- Kallosh & Linde arXiv:1306:3214
- Buchmuller, Domcke & Kamada arXiv:1306.3471
- Kallosh & Linde arXiv:1306.5220
- Farakos, Kehagias and Riotto arXiv:1307.1137
- Roest, Scalisi & Zavala arXiv:1307.4343
- Kiritsis arXiv:1307.5873
- Ferrara, Kallosh, Linde & Porrati arXiv:1307.7696

# No-Scale Supergravity Inflation in Light of BICEP2

- Possible to accommodate quadratic potential  $K = -3\ln(T+T^*) + |\phi|^2 W = e^{\frac{-\phi^2}{2}} \left(\mu - \frac{m}{2}\phi^2\right)$
- Can stabilize T, Im  $\phi$  while Re  $\phi$  inflates



JE, Garc á, Nanopoulos & Olive, arXiv:1403.7518

# No-Scale Supergravity Inflation in Light of BICEP2



- Re T and Im T also well-behaved
- Sufficient inflation
- BICEP2-compatible

JE, Garc n, Nanopoulos & Olive, arXiv:1403.7518



# Supersymmetric Starobinsky Inflation in Light of BICEP2

- Im  $\phi$  has m<sup>2</sup> $\phi$ <sup>2</sup> potential when Re  $\phi = 0$
- But Im \u03c6 does not inflate
- Inflation driven by Re  $\phi$
- Looks like Starobinsky
- BICEP2-incompatible

JE, Garc á, Nanopoulos & Olive, arXiv:1403.7518



# Supersymmetric Starobinsky Inflation in Light of BICEP2

• Modify K ähler potential? 0.5-0.5 $K = -3 \ln \left( T + T^* - rac{\phi \phi^*}{3} - rac{(T+T*)^n}{\Lambda^2} \right)$ 0 V/m<sup>2</sup> -800 • Unstable potential! -105 • Can stabilize Im T 10  $K = -3 \ln \left( T + T^* - rac{\phi \phi^*}{3} - rac{(T+T*)^n}{\Lambda^2} + rac{(\phi \phi^*)^2}{\Lambda_{\phi}^2} 
ight)$ • Alternative realization of 600  $V/m^2$ 300 quadratic inflation in no--10Re d scale supergravity -5 -0.1Im T

-0.3

10

JE, Garc á, Nanopoulos & Olive, arXiv:1403.7518

# Summary

- Inflation may solve some of the biggest problems in cosmology: age, size, homogeneity, flatness, ...
- Structures originated from quantum effects in the early Universe
- Another scalar field (inflaton) in the sky?
- Quantum gravitational radiation detected?
- Tension between Planck and BICEP2
- Inflation also cries out for supersymmetry
- Open season for building inflationary models
- Inflaton = superpartner of neutrino?

# Time for China to get involved in CMB Experiments?



# Crossing the Scalar Rubicon: Once or Twice?





# No-Scale Framework for Particle Physics & Dark Matter

- Incorporating Starobinsky-like og (M<sub>in</sub> (GeV)) inflation, leptogenesis, neutrino masses, LHC constraints, supersymmetric dark matter,
- Stringy origin...?

JE, Nanopoulos & Olive, arXiv:1310.4770



### Inflationary Models in Light of BICEP2

• Consistency condition for simple power-laws:

• They require  $r = 8\left(1 - n_s - \frac{1}{N}\right)$  OK with BICEP2





## Stabilization of Modulus Field

- With modified geometrical (K ähler) structure:  $K = -3\ln\left(T + T^* - \frac{|\phi|^2}{3} + \frac{(T + T^* - 1)^4 + d(T - T^*)^4}{\Lambda^2}\right)$
- Effective inflaton potential ~ Starobinsky
- Other field directions stabilized:



JE, Nanopoulos & Olive, arXiv:1307.3537

## Generalizations

• Common no-scale geometrical (Kähler) structure many superpotentials:

$$K = -3\ln\left(T + T^* - \frac{|\phi|^2}{3}\right), W = M\left[\frac{\phi^2}{2} - \frac{\phi^3}{3\sqrt{3}}\right]$$
$$W = \frac{M}{4}(T - 1/2)^2(1 + 10T + 2\sqrt{3}\phi)W = \frac{M}{4}(T - 1/2)^2(5 + 2T + 2\sqrt{3}\phi)$$
$$W = \sqrt{3}M\phi(T - 1/2)W = 2\sqrt{3}M\phi T(T - 1/2)W = M\left[\sqrt{3}(T^2 - 1/4)\phi + (T - 1/2)\phi^2\right]$$
$$\Delta W = \left[\frac{(T - 1/2)^n 2^{(n-3)}}{(2T + 1)^{(n-3)}}\right] \quad \Delta W = \left[\frac{(T - 1/2)^n 2^{(n-2)}\phi}{(2T + 1)^{(n-2)}}\right]$$

also yield Starobinsky-like inflationary models

JE, Nanopoulos & Olive, arXiv:1307.3537

# Beyond Starobinsky

• Exponential approach to constant potential:

$$V = A \left(1 - \delta e^{-Bx} + \mathcal{O}(e^{-2Bx})\right)$$

• Relations between observables:

$$n_s = 1 - 2B^2 \delta e^{-Bx}$$
  

$$r = 8B^2 \delta^2 e^{-2Bx},$$
  

$$N_* = \frac{1}{B^2 \delta} e^{+Bx}.$$

$$n_s \;=\; 1 - \frac{2}{N_*} \;, r \;=\; \frac{8}{B^2 N_*^2}$$

- E.g., multiple no-scale moduli:
- $K \ni -\Sigma_i N_i \ln(T_i + T_i^*): N_i > 0, \ \Sigma_i N_i = 3$
- Characteristic of generic string compactifications
- Tensor/scalar ratio may be < prediction of Starobinsky model:  $B = \sqrt{\left(\frac{2}{N_i}\right)}$  $r = \frac{4N_i}{N_*^2}.$
- String phenomenology via the CMB? JE, Nanopoulos & Olive, arXiv:1307.353

# Completing the Holy Trinity

- Scale hierarchy possible only in theory that can be calculated over many magnitudes of energy
   Renormalizable
- Theorem: (1) vectors (2) fermions (3) scalars
- Need to specify:

Cornwall, Levin & Tiktopoulos; Bell; Llewellyn-Smith

(1) group (2) representations (3) symmetry breaking
(1) = SU(3) × SU(2) × U(1) [so far]
(2) = Singlets + doublets + triplets

• Finally:

(3) A scalar, mechanism of symmetry breaking



# Measurements of CMB Fluctuations with Planck Satellite



### Wess-Zumino Inflation in Light of Planck

• Consistent with Planck for  $x_i = 0.3, 0.4$ 

Value of $x_i$	0.1	0.2	8.3	0.4
Derived quantity				
$\frac{v^2}{M_{Pl}^2}$	18000	4200	1600	710
$\epsilon$	0.0085	0.0067	0.0045	0.0020
$\eta$	0.0062	0.00074	-0.0073	-0.022
ξ	-0.000053	-0.000077	-0.000079	-0.000050
r	0.14	0.11	0.072	0.031
$n_s$	0.961	0.961	0.958	0.945
$\alpha_s$	$-1.4 imes10^{-6}$	$-1.3 imes10^{-6}$	$-1.4 \times 10^{-6}$	$-1.1 \times 10^{-6}$
$\lambda$	$4.3 imes10^{-8}$	$1.0  imes 10^{-7}$	$2.1  imes 10^{-7}$	$4.1 \times 10^{-7}$

Good

inflation

& Mayromatos

• Numbers calculated for N = 50 e-folds



# Unofficial Combination of Higgs Data



### Inflationary Models in Light of BICEP2

- QUANTUM GRAVITY RADIATION!
- BICEP2 observations support for inflation

• "Large" tensor-to-scalar ratio:  $r = 0.20^{+0.1}_{-0.1}$ 



Starobinsky R<sup>2</sup> dead, quadratic inflation OK?

# It Walks and Quacks like a Higgs



So far, it looks like a (bog) Standard Model Higgs Boson

#### JE & Tevong You, arXiv:1303.3879

# Dixit Swedish Academy

Today we believe that "Beyond any reasonable doubt, it is a Higgs boson." [1] http://www.nobelprize.org/nobel\_prizes/physics/laureates/2013/a dvanced-physicsprize2013.pdf

# Gravitino Inflation

- Inflation driven by condensate of gravitinos?
- Possible with suitable non-perturbative parameters

 $e^{-1} \mathcal{V}_{\text{eff infl}} = \tilde{\kappa}^{-4} \left[ 0.4878 \,\sigma^4 \ln^2 \sigma^2 - 0.2549 \,\sigma^4 - 0.0022 \,\sigma^2 + 0.2574 \right]$ 

• Potential can yield good  $(n_s, r)$ :



• Need hypothesis on dynamics

# Scalars Come of Age

- A new boson discovered at the LHC
  - A scalar, beyond any reasonable doubt
  - Consistent with the Higgs of the Standard Model
- Circumstantial evidence for scalar inflaton
  - $-\Omega \sim 1$ , tilt in spectrum of scalar perturbations, ...
  - No sign of non-Gaussianity, strings, defects, ...
- Data compatible with Starobinsky  $(R + R^2)$  model
- Similar predictions in Higgs inflation: **BUT** M<sub>H</sub>?