

# Physics Analysis

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**The University of Michigan**

**CCAST Workshop @Tsinghua University**

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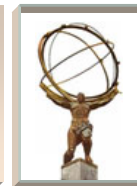
# Outline



- Collider physics analysis basics
- Monte Carlo Tools
- Physics Objects Reconstructions
- LHC Physics potential of ATLAS, CMS



# Basics



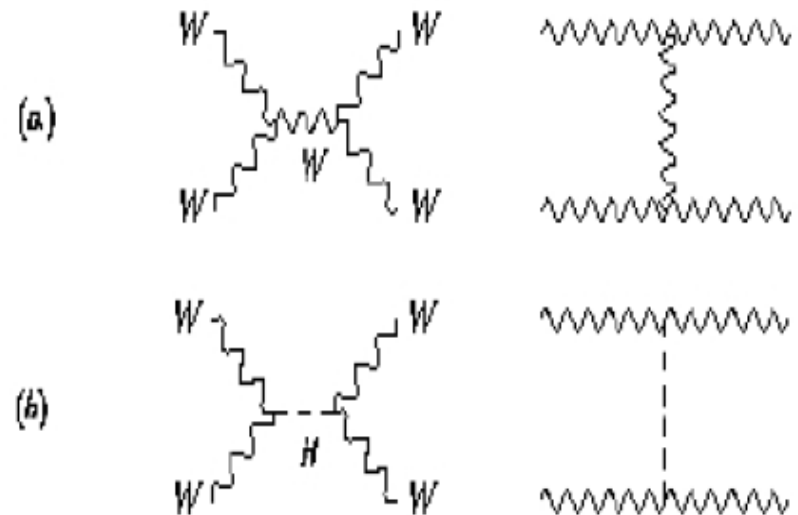
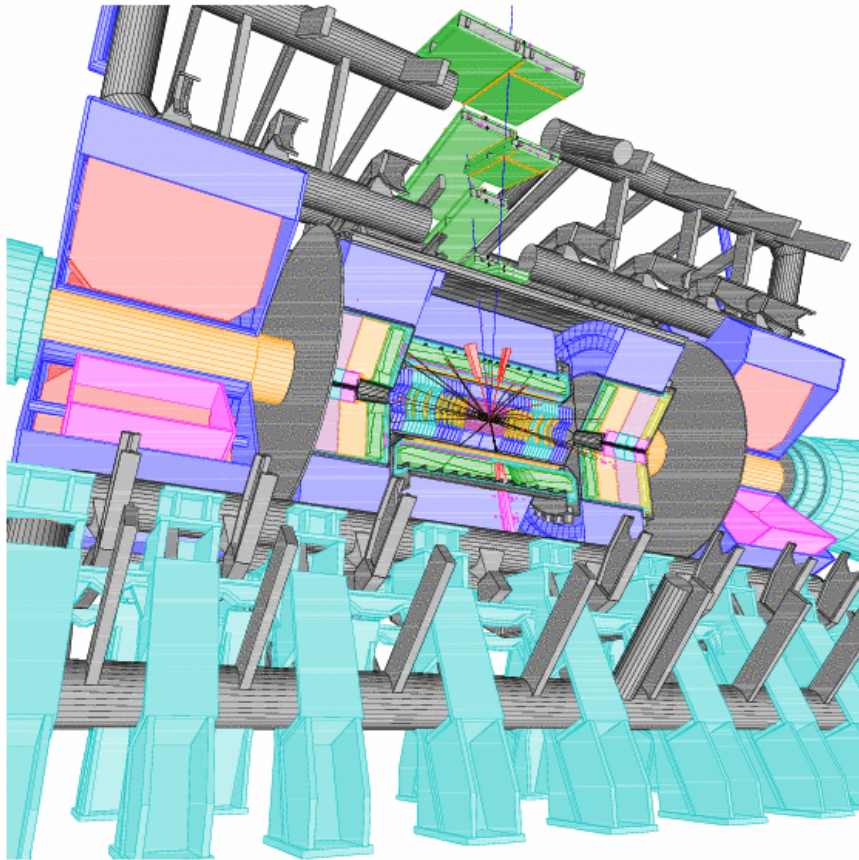
- Hadron Colliders - especially those that allow access to a new energy regime – are machines for discovery
  - In the case of the TeV scale, this is reinforced by the fact that the known SM forces and particles violate unitarity at  $\sim 1$  TeV: there must be something new (if only a SM Higgs)
- Discovery means producing convincing evidence of something new.
- New physics models: produce visible new phenomena
- Goals of analysis in this case: produce the evidence”
  - Separation of a signal from the background
  - Need to show that the probability of the ‘signal’ from known sources is small
    - **Must demonstrate that the backgrounds are well understood**
    - **Uncertainties: statistics and systematic**



# Physics Analysis Goal

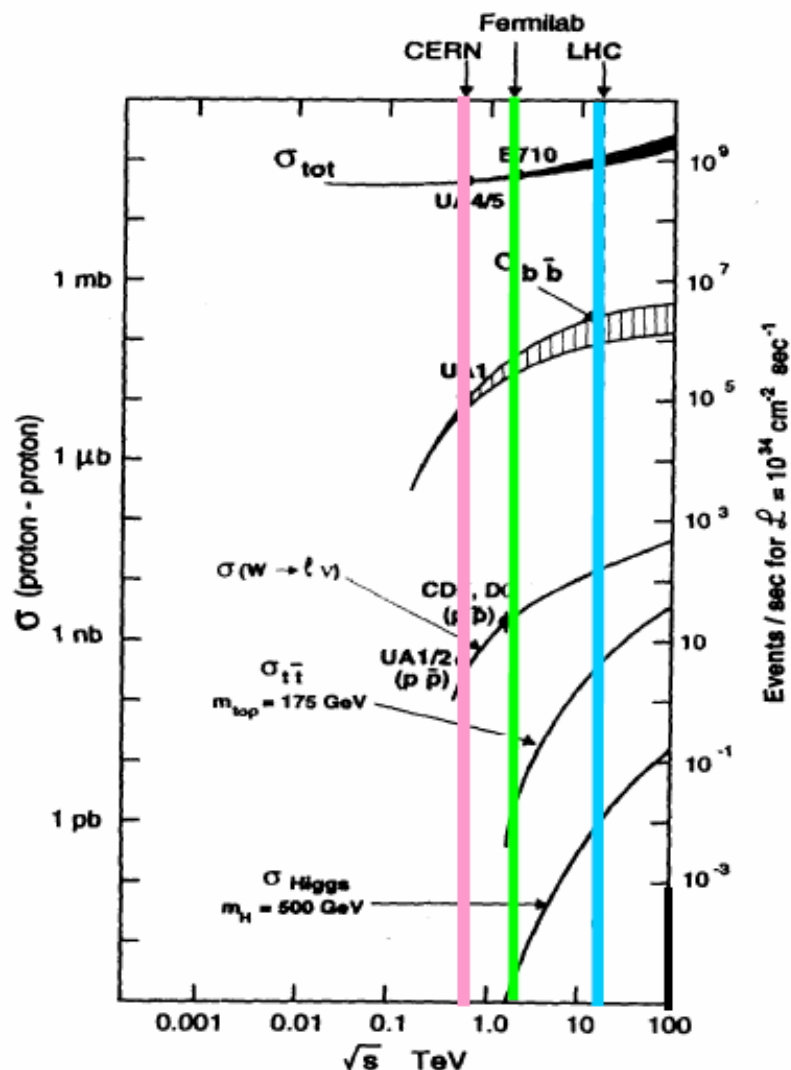
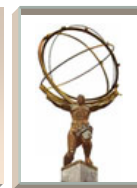


Translate from measurement To Physics process





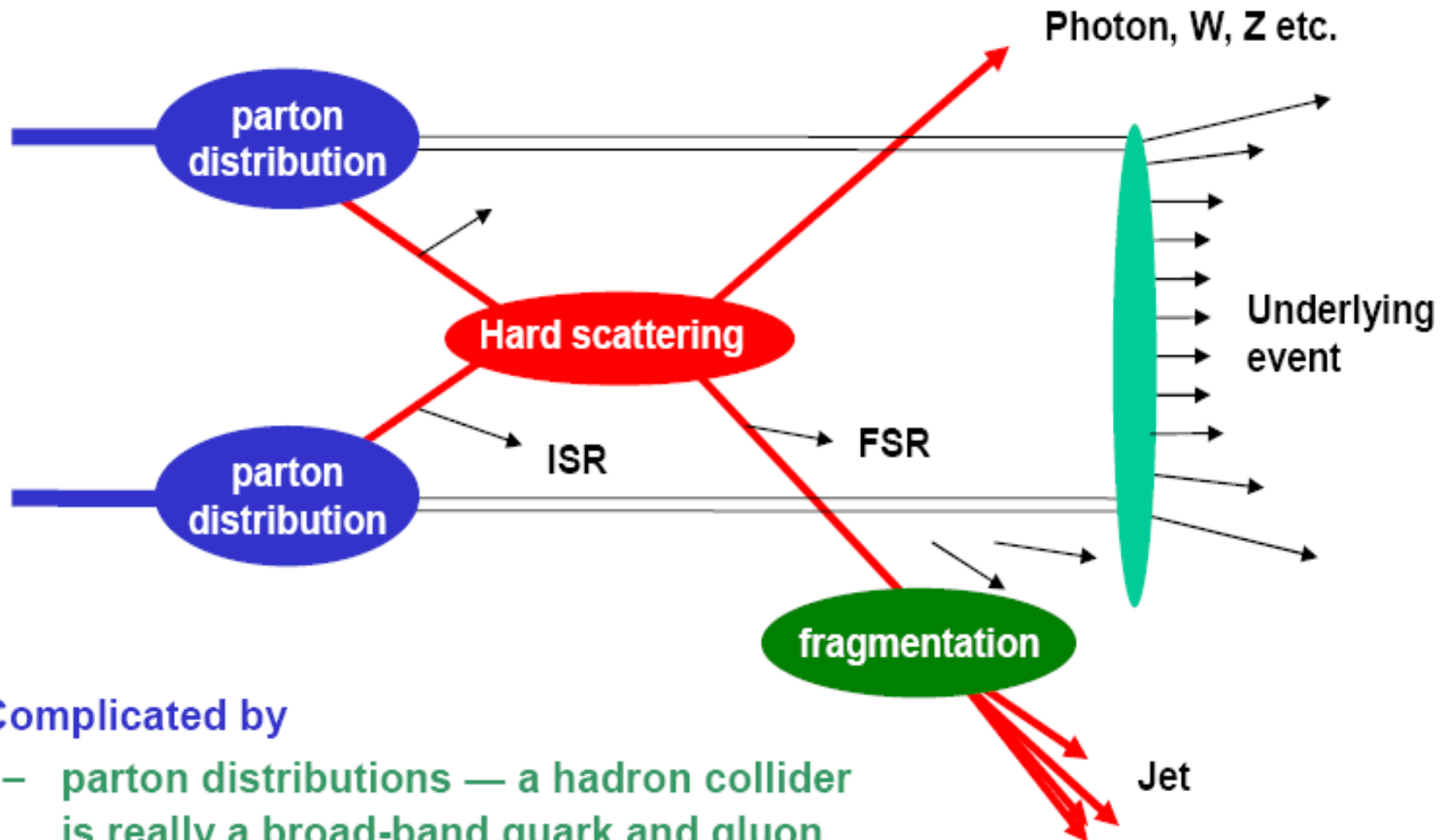
# Collider Basics: Cross section



- Rates Determined by:
  - Hard Scattering  $\sigma$
  - Parton Luminosity
- QCD Processes Dominate
  - EW rates lower by  $\alpha / \alpha_{\text{strong}}$
- Cross Sections Decrease Rapidly with  $s$ 
  - Heavy particles difficult to produce



# Basics: Hadron Collision Processes



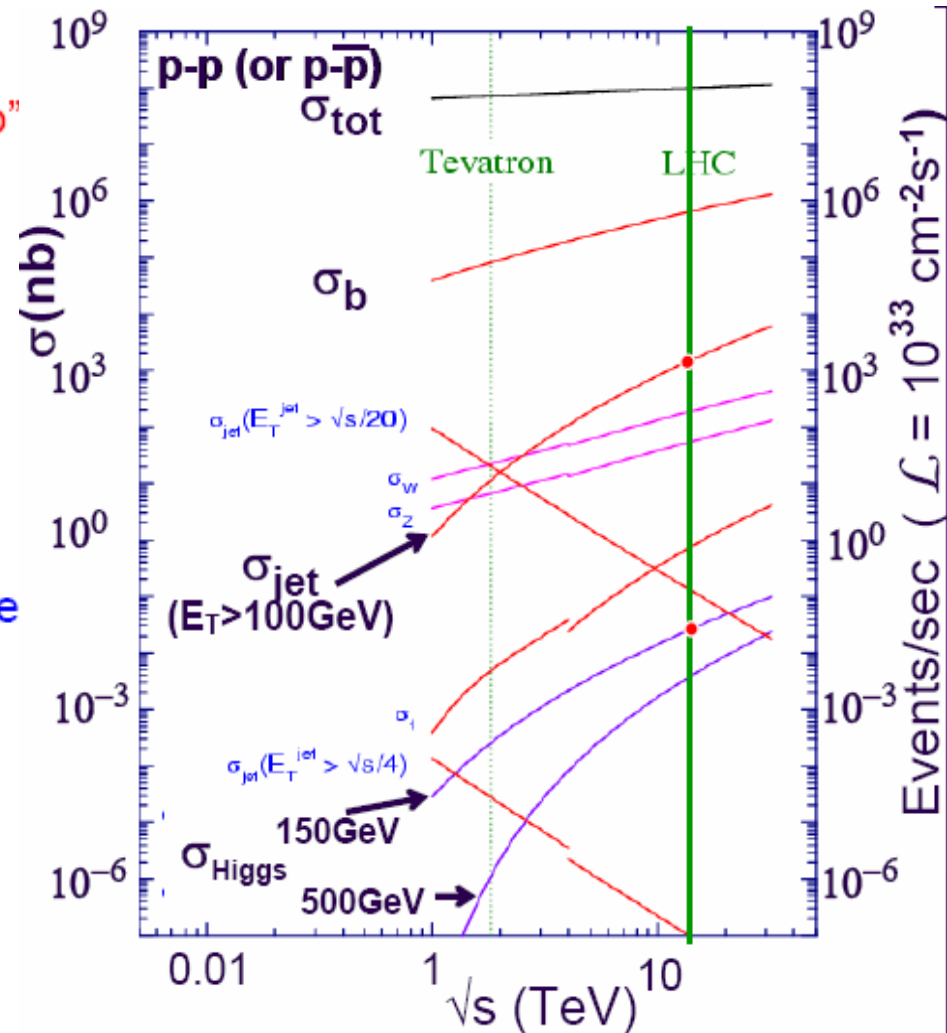
- Complicated by
  - parton distributions — a hadron collider is really a broad-band quark and gluon collider → different from  $e^+e^-$  collisions



# Implications at LHC



- Something happens every crossing
  - 25 inelastic evts/crossing at  $10^{34}$  “Pile-up”
- Must Select Events of Interest: **Trigger**
  - Must know what you've thrown out
  - Analysis must be trigger-aware
- Jets Dominate Hard Scattering Rate
  - Can isolate EW processes only they have something besides jets, eg leptons
  - Potential source of bckgnd “Fakes”
  - Detector mis-measurements can induce false signals
- W, Z: Bckgnd for Top, Higgs, SUSY

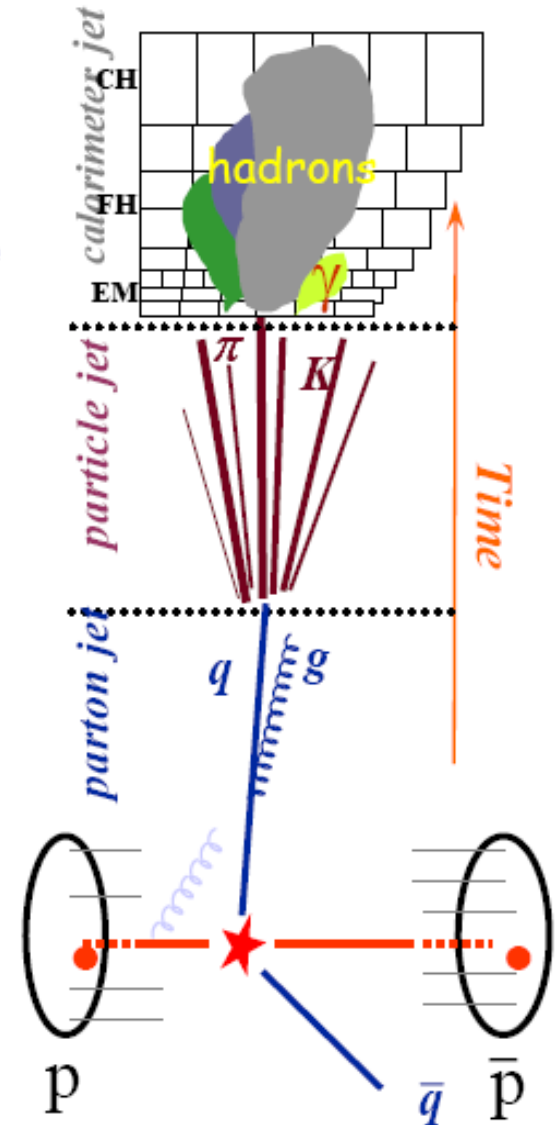




# Simulation Tools

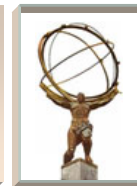


- A “Monte Carlo” is a Fortran or C++ program that generates events
- Events vary from one to the next (random numbers) — expect to reproduce both the average behavior and fluctuations of real data
- Event Generators may be
  - **parton level:**
    - Parton Distribution functions
    - Hard interaction matrix element
  - **and may also handle:**
    - Initial state radiation
    - Final state radiation
    - Underlying event
    - Hadronization and decays
- Separate programs for Detector Simulation
  - **GEANT** is by far the most commonly used





# Event Generators



- **LO: Pythia, Herwig** (include hardronization and 'soft jets')
- 'Underline' -- **Jimmy**
- **Improved generators (NLO)**
  - **MCFM, RESBOS, ALPGEN:** these go beyond the PYTHIA/HERWIG models and include real matrix element calculations. They are being improved today.
  - **MC@NLO, CKKW:** generate inclusive jets at NLO
  - **MadGraph/MadEvent, WhiZard/O'Mega, Sherpa/Amegic++ :** these are very complete and sophisticated event generators for MSSM processes.
    - spin amplitudes!
    - efficient generation of codes for reduced CPU load
    - multi-channel adaptive sampling of phase space
    - structure functions for incoming partons
    - interface to, e.g., PYTHIA for hadronization (Sherpa does this itself)
    - These three programs have been shown to agree! (hep-ph/0512012)



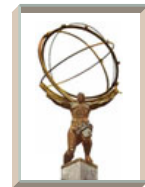
# How Much Can We Trust MC



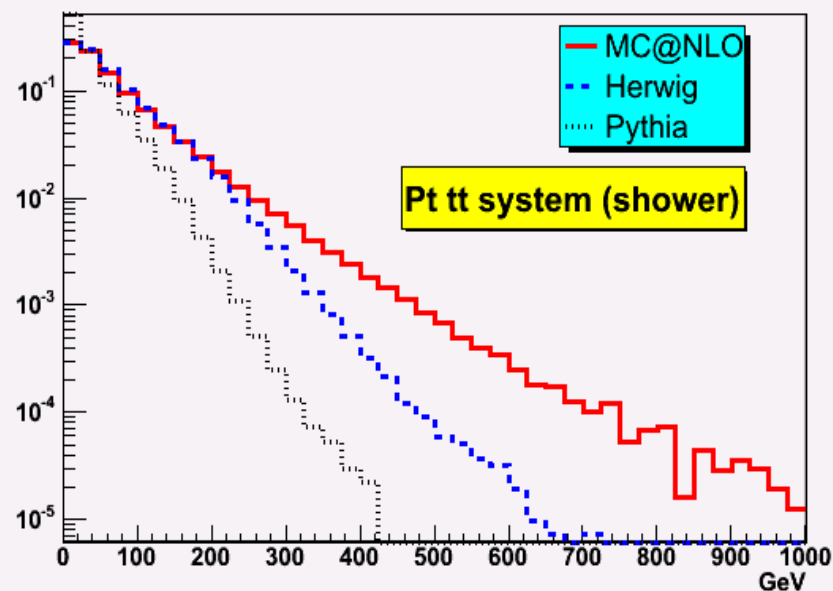
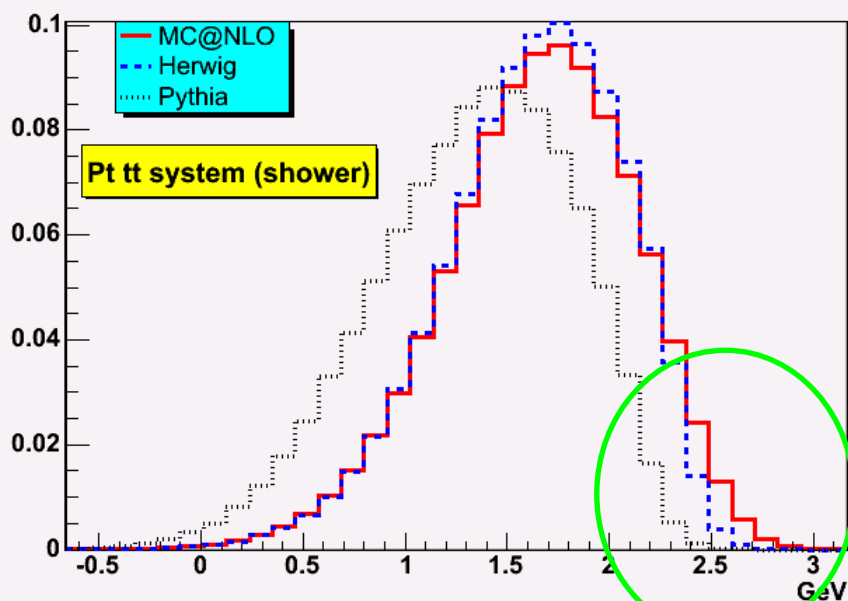
- PDF and Underline modeling need data to fit
- Event generators
  - May or may not generate additional jets through parton showering
  - May or may not treat spins properly
  - May or may not get the cross section right
    - NLO much better than LO – but sometimes no choice!
- It's not simple
  - One can not necessarily just add (for example) a  $W+1$  jet simulation and  $W+2$  jets and  $W+3$  jets to model  $W+n$  jets signal. Likely to be double counting.
  - One can not necessarily just run a  $W+1$  jet simulation and generate the extra jets through parton showering either...



# Generators: MC@NLO, Herwig, Pythia



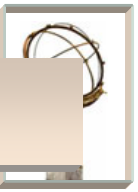
**Example:** distributions on top-anti-top characteristic –  $P_T$  of the whole system  
 $P_T$  of t-tbar system is balanced by ISR & FSR



•  $P_T(\text{tt system})$

- Herwig & MCatNLO agree at low  $P_T$ ,
- At large  $P_T$  MCatNLO 'harder'





## Underlying event (UE) / minimum bias

- Extremely difficult to predict the magnitude of the UE at LHC
  - Will have to learn much more from Tevatron before startup
- Various models exists
  - Herwig's UE and minimum bias shows much less activity compared to Pythia.
    - This has always been a problem in Herwig.
  - Jimmy is developed as alternative model for UE at ep collisions
    - Various 'tunings' exist – leading to wildly different results



## Top Events

Tune of  
Butterworth

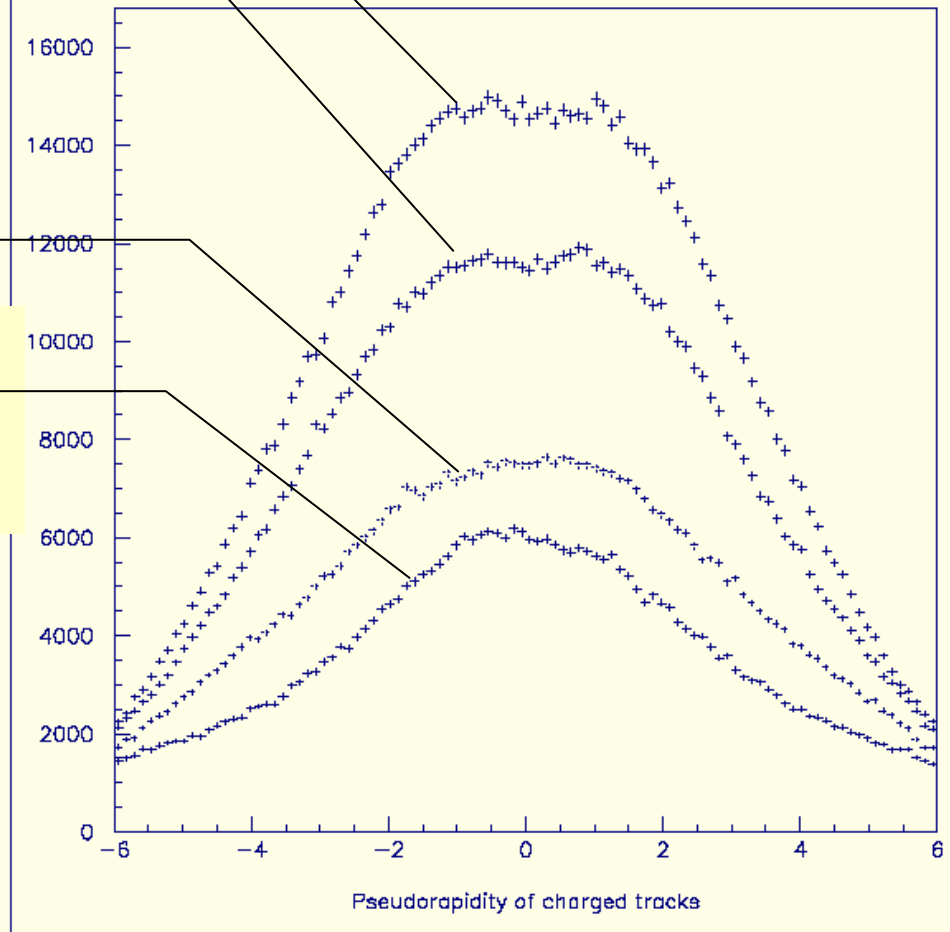
Standard  
Jimmy

Standard  
Pythia

Standard  
Herwig

Plot shows rapidity distribution of all  
charged tracks

Note that Jimmy has many more tracks





# How Well Can We Simulate Detector?



- Detector Simulation – is only as good as the geometrical modeling of the detector – are all the cables and support structure in your model?
- Short time structure in current detector adds another dimension – nuclear de-excitations, drift charge in argon can be slower than beam bunch crossing time!
- Do not blindly trust tails of the distributions or rare processes
  - Random number may not populate them fully
  - Modeling not verified at this level -- e.g. MC estimate the probability for a jet to be reconstructed as a photon – a  $10^{-3}$  or  $10^{-4}$  probability



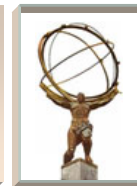
# Analysis Strategy at LHC



- Characterize Bulk of Cross Section “Soft Physics”
  - Tracks
- Identify Dominant  $2 \rightarrow 2$  QCD Processes
  - Jets
- Develop Strategies for Selecting EW Processes
  - $e, \mu, \tau, \nu, \gamma$
- Reconstruct Heavy Objects Produced Strongly
  - Top, SUSY(?)
- Understand Discovery Potential for Low Rate EW Processes
  - Higgs



# Offline Data Process for Analysis



- LHC reconstruction very complex
  - Many channels, many hits: Reconstruction is slow
- Large data collection rate and large event size
  - TB of storage required
- Cannot support bulk reconstruction by individual physicists
- Organized Production effort

RAW → RECO → AOD → TAG → NTUPL
- Results available through Data Delivery System
  - “Analysis” is performed on output of Offline reconstruction



# Reconstructed Physics Objects



- Collections of Candidate Objects
  - Tracks, Jets, Electrons,  $\mu$ ,  $\tau$ , vertices, missing energy, heavy flavor
- Selections Performed Using Loose Criteria
  - Can tighten during analysis phase
- No attempt to uniquely identify objects
  - Same energy deposition may appear as jet, e and  $\gamma$  candidate
- Support for multiple algorithms
  - Best jet algorithm for Top analysis may not be the best algorithm for QCD studies

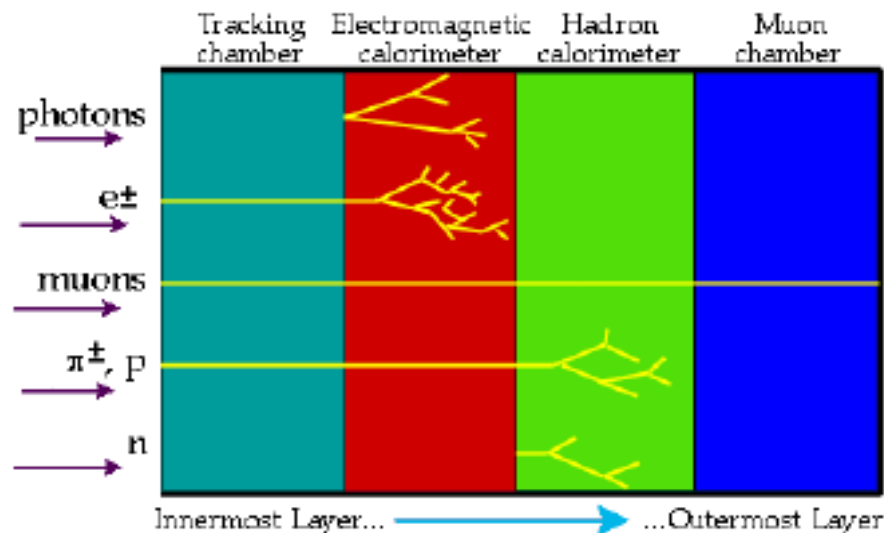
Physicists impose consistent interpretation during analysis phase



# Electron Identification



- Electrons signature:
  - Energy Deposition in EM Calorimeter
  - Track pointing at the energy deposition and with momentum consistent with calorimeter energy
  - Little or no energy in hadron calorimeter

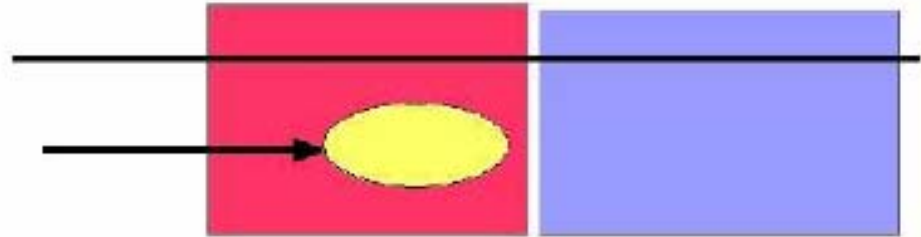




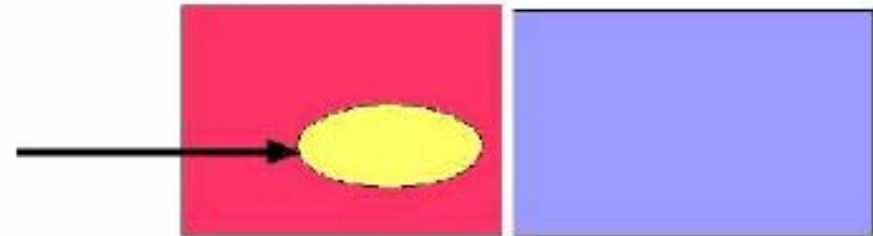
# Background for Electron ID



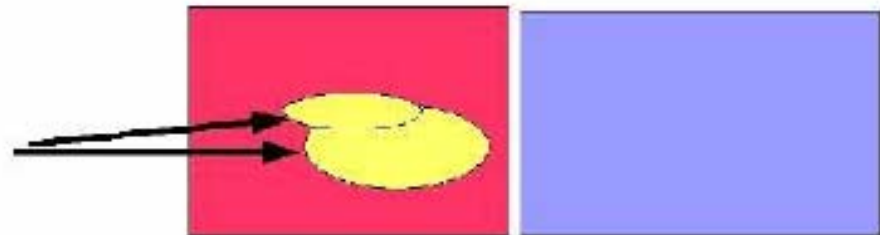
$\pi^0$  and non-interacting  $\pi^+$



Early showering  $\pi^+$



Photon Conversions





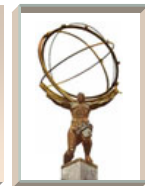
# Reject Electron Background (I)



Choice of variables depends on detector. Some possibilities:

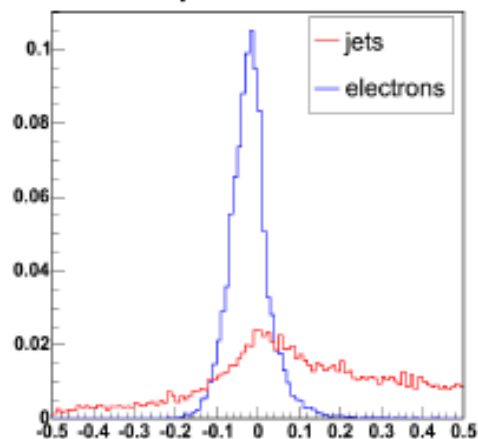
- Shower Shape Variables:
  - Longitudinal shape: ratio of energy in depth segments of calorimeter
  - Transverse shape: Hadron showers typically wider than electrons (also rejects  $\pi^0 \pi^+$  overlap)
  - Had/EM: Expect very little energy deposit in HAD calorimeter





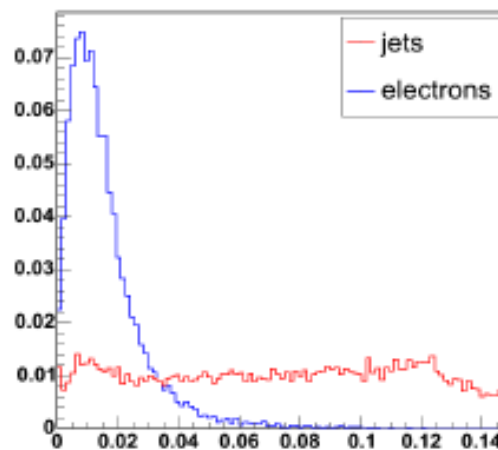
# Shower Shape Distributions: Electron vs. Jets

Shower profile

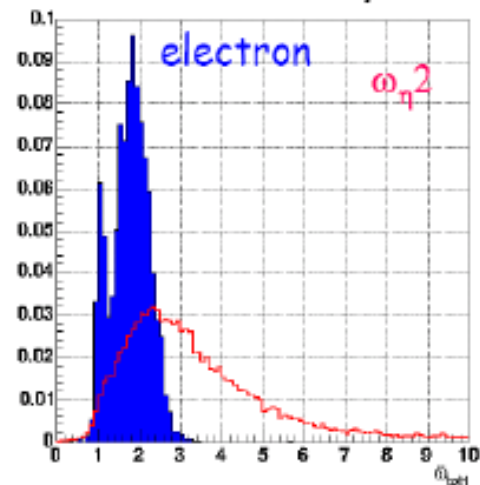


CDF

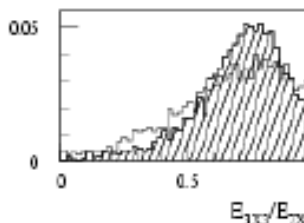
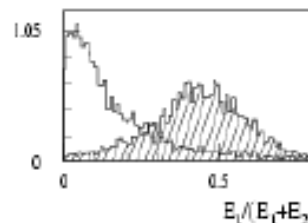
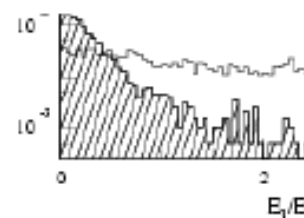
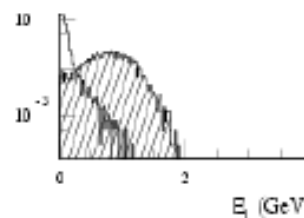
Had/em Ratio



Shower width in strips



ATLAS





# Reject Electron Background (II)



- Track-Shower Matching:
  - $E/P$ : Ratio of energy in calorimeter to momentum in tracker
  - Pointing: Compare extrapolated position of track to position of EM cluster

Caution:

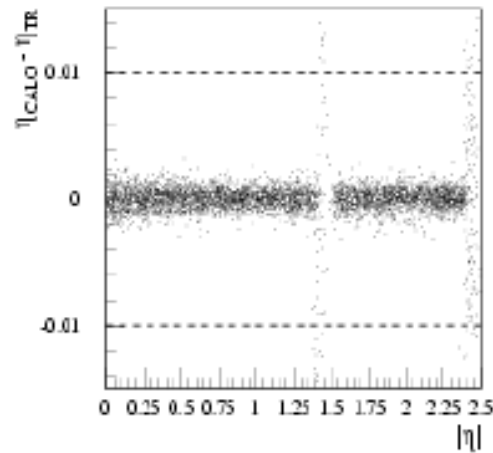
- Significant material in LHC trackers means electron bremsstrahlung
- Correct modeling of material distribution necessary both for defining selection criteria and for estimating efficiency



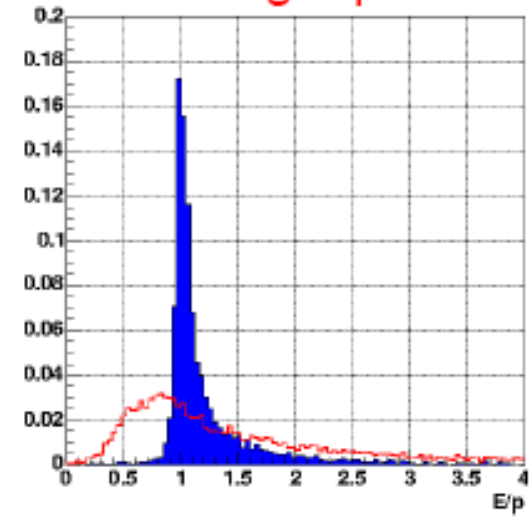
# Track Matching - E/P



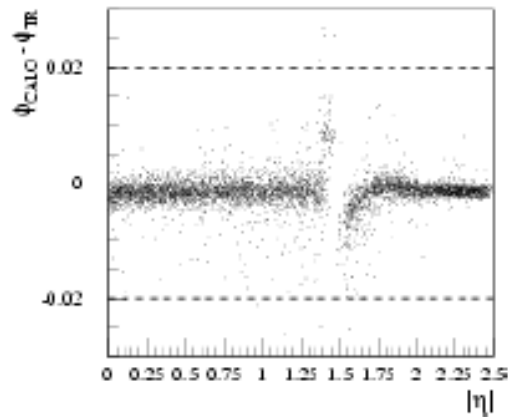
Matching in  $\eta$



Matching E/p



Matching in  $\phi$





# Reject Electron Background (III)



- Large amount of material also means photon conversions are an issue (photons from  $\pi^0$ )
  - Explicit removal of conversions:
    - Require hits in pixel layer (most of material outside this)
    - Look for second track from conversion: cut on reconstructed mass and angle



# Reject Electron Background (IV)



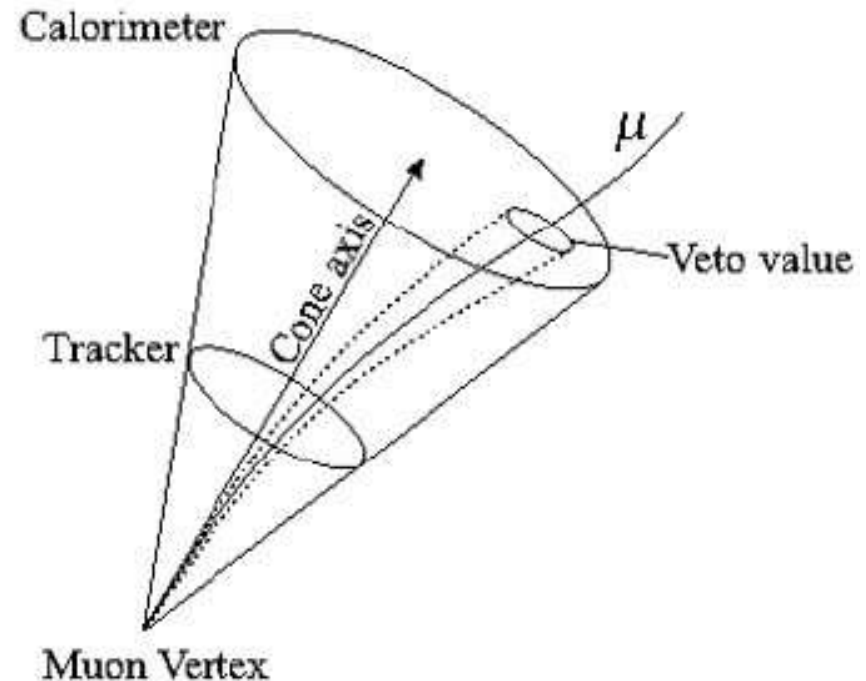
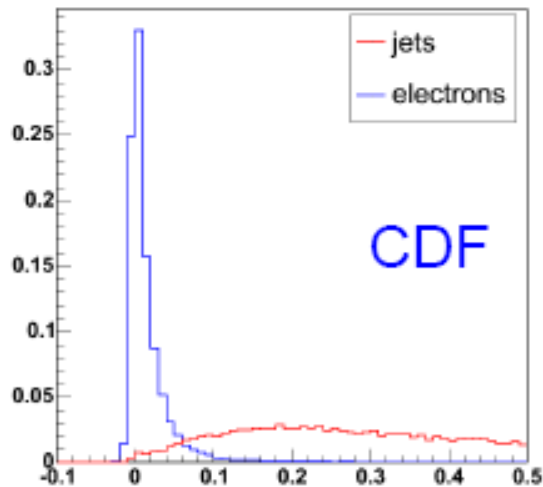
- Isolation:
  - Study ratio of energy in annulus round electron to energy of electron
  - As noted above: Does not work for all physics processes
- Transitions Radiation and  $dE/dx$ :
  - CDF drift chamber measures  $dE/dx$ : sensitive to particle velocity: helps for low momentum  $e$
  - Atlas tracker has TR function: Can require high energy deposition hit, at cost of efficiency



# Lepton Isolation



Calorimeter Isolation



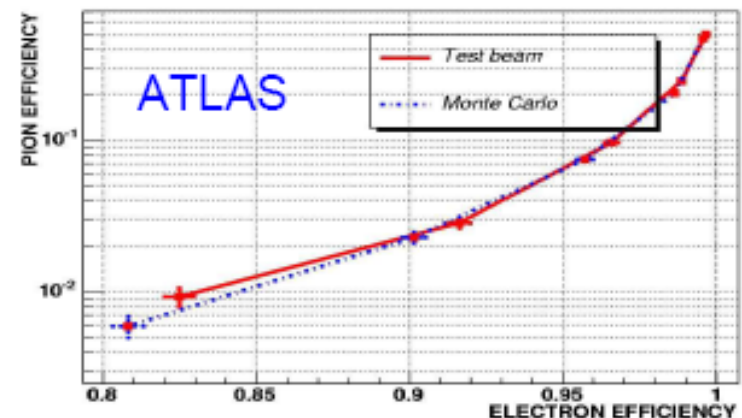
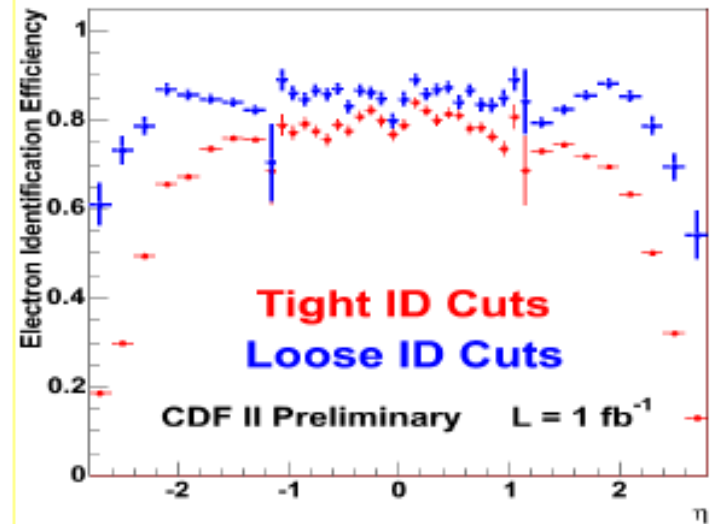
- Sum of  $E_t$  in calorimeter in a cone
- Sum of  $P_t$  in Tracker in a cone



# Overall Efficiency of Electron ID



- Measure when possible using real data:
  - **W** from no-track trigger to measure tracking efficiency
  - **Z** with one tight electron and with loose selection
- Use simulation to extrapolate kinematics and correct for environmental issues (eg isolation)





# Photon Identification



- Use same variables as for electron selection, with tighter cuts
  - Unconverted photons have track veto
  - Converted photons independently analyzed by looking for the second track
  - Emphasis on shower shape variables
    - Photons shower later than electrons
    - $\pi^0$  decay to  $2\gamma$  so probability of early shower twice as large
- Isolation is critical

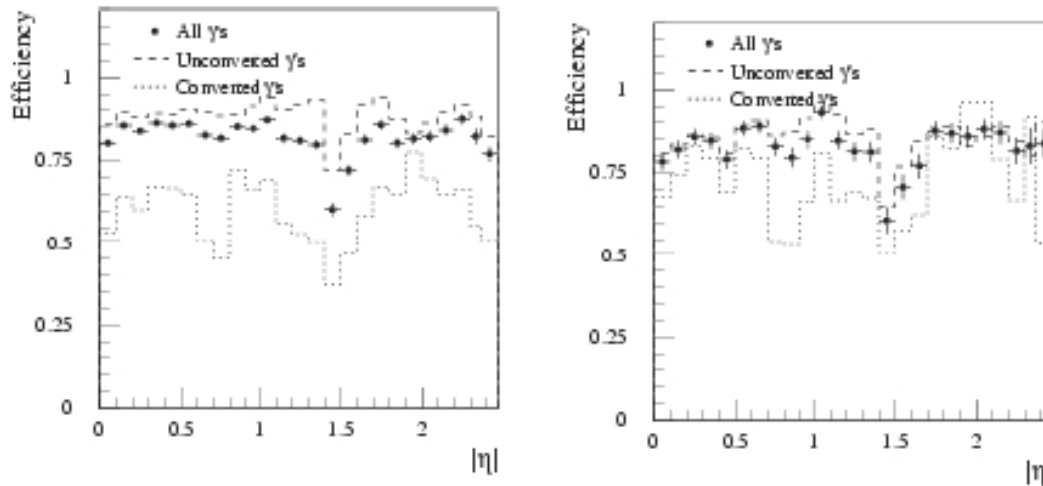
ATLAS and CMS have different emphasis due to different detector designs, but overall performance for Higgs similar



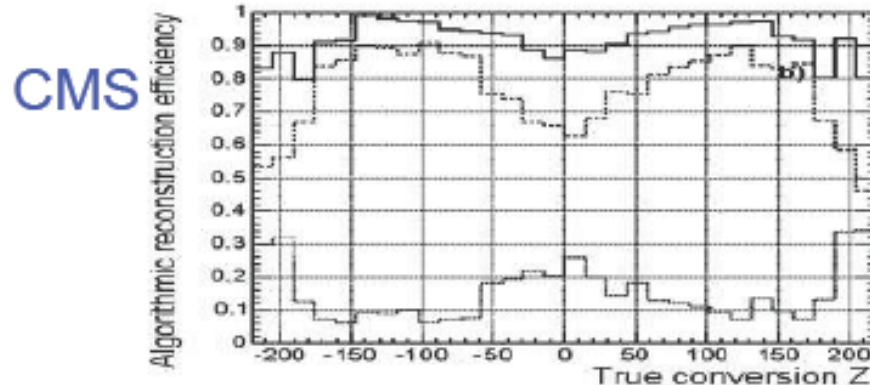
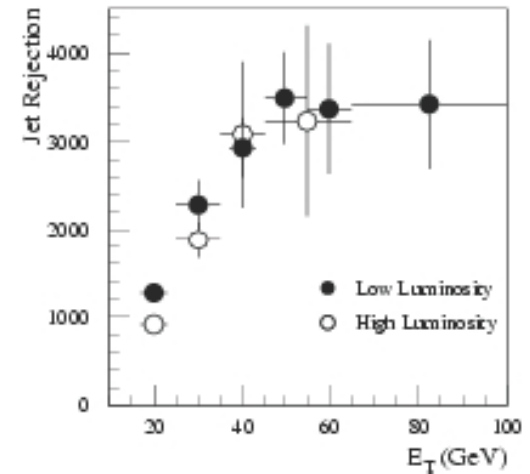
# Photon ID efficiency and Background Rejection



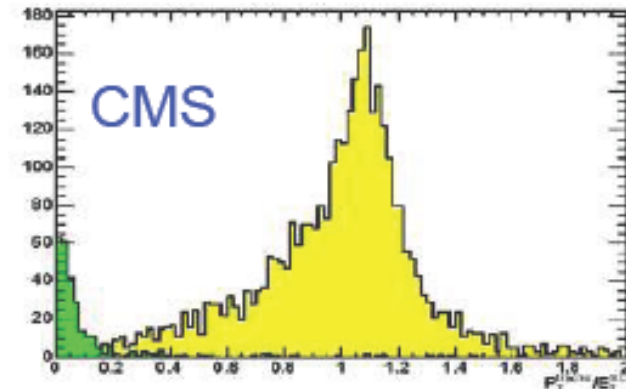
Atlas: Efficiency: Low and High Luminosity



Atlas Jet Rejection



Conversion Probability vs  $\eta$



Pt Tracks/Et ECAL

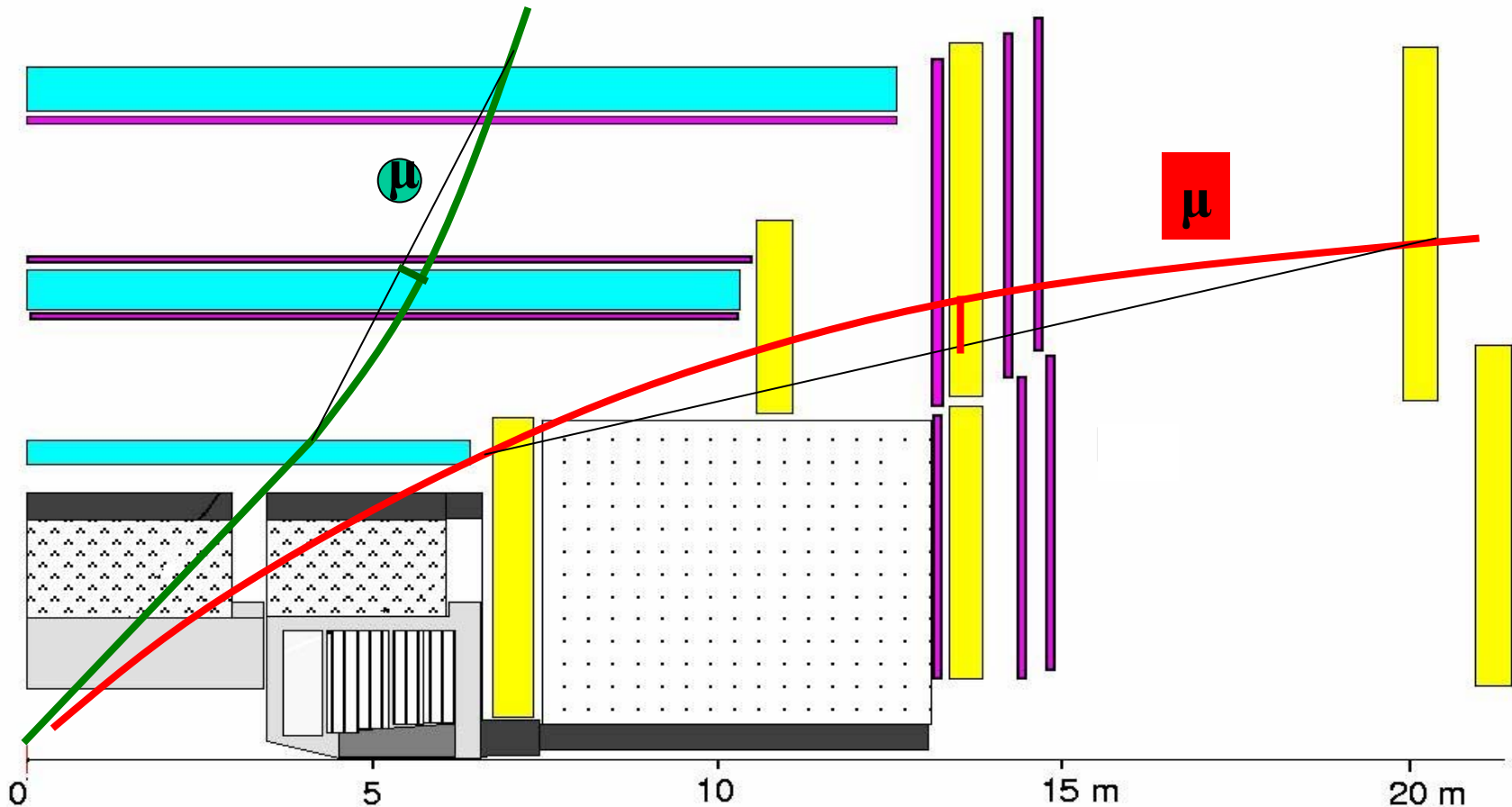


# ATLAS Muon Momentum Measurement



1 TeV muon sagitta  $\sim 500 \mu\text{m}$

Sagitta measurement precision  $\sim 50 \mu\text{m}$





# ATLAS Muon Reconstruction and Identification

**Muonboy** : Muon Reconstruction in the Muon System, +back-fit to IP

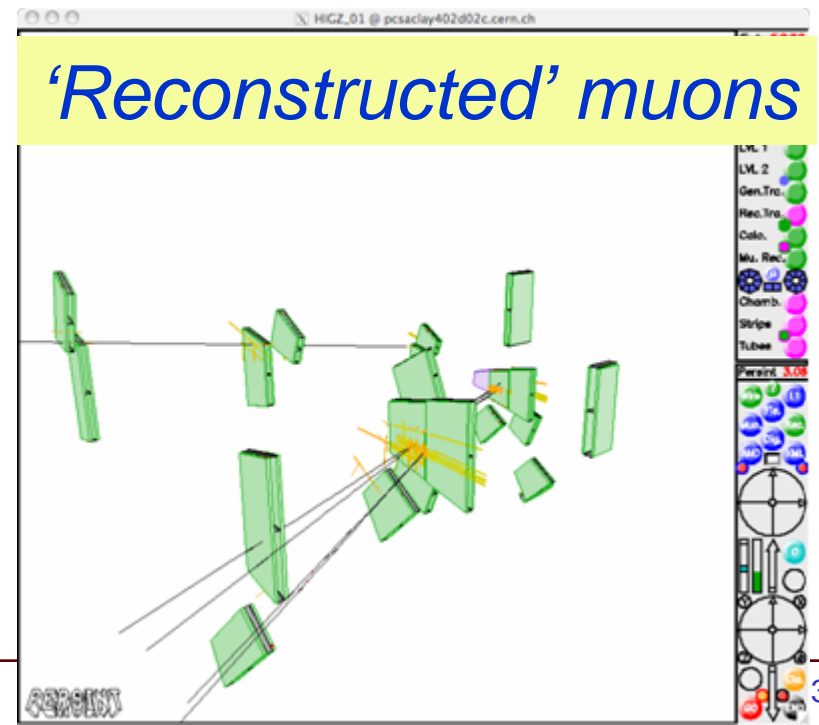
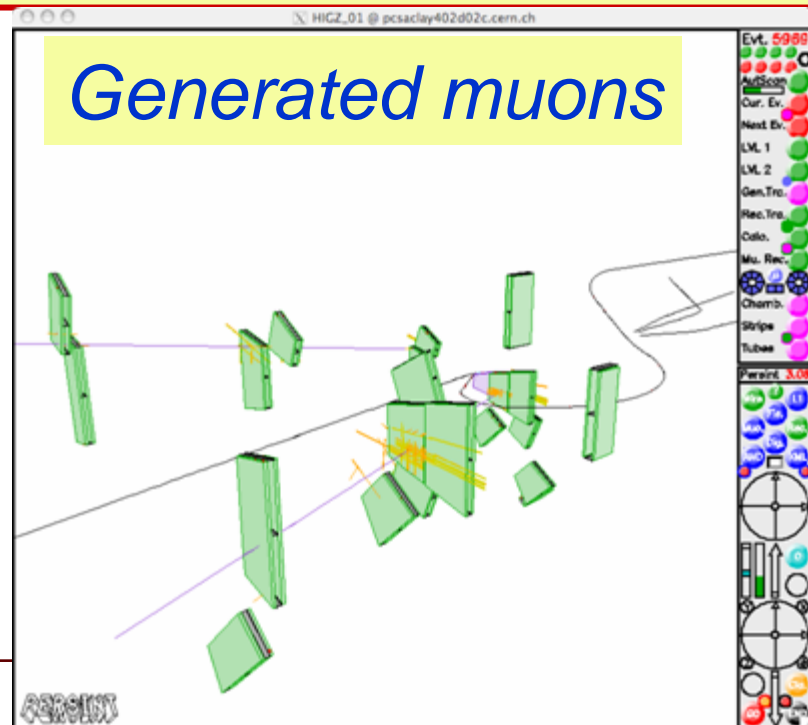
**STACO** : Statistical Combination of the Inner Detector and the Muon System

**MuTag** : Tag of ID tracks using Inner Stations Segments

**MOORE** : Muon Reconstruction in the Muon System

**MuonID Standalone** : back-fit from moore reconstruction to IP

**MuonID Combined**: Combine Inner tracker and the Muon System

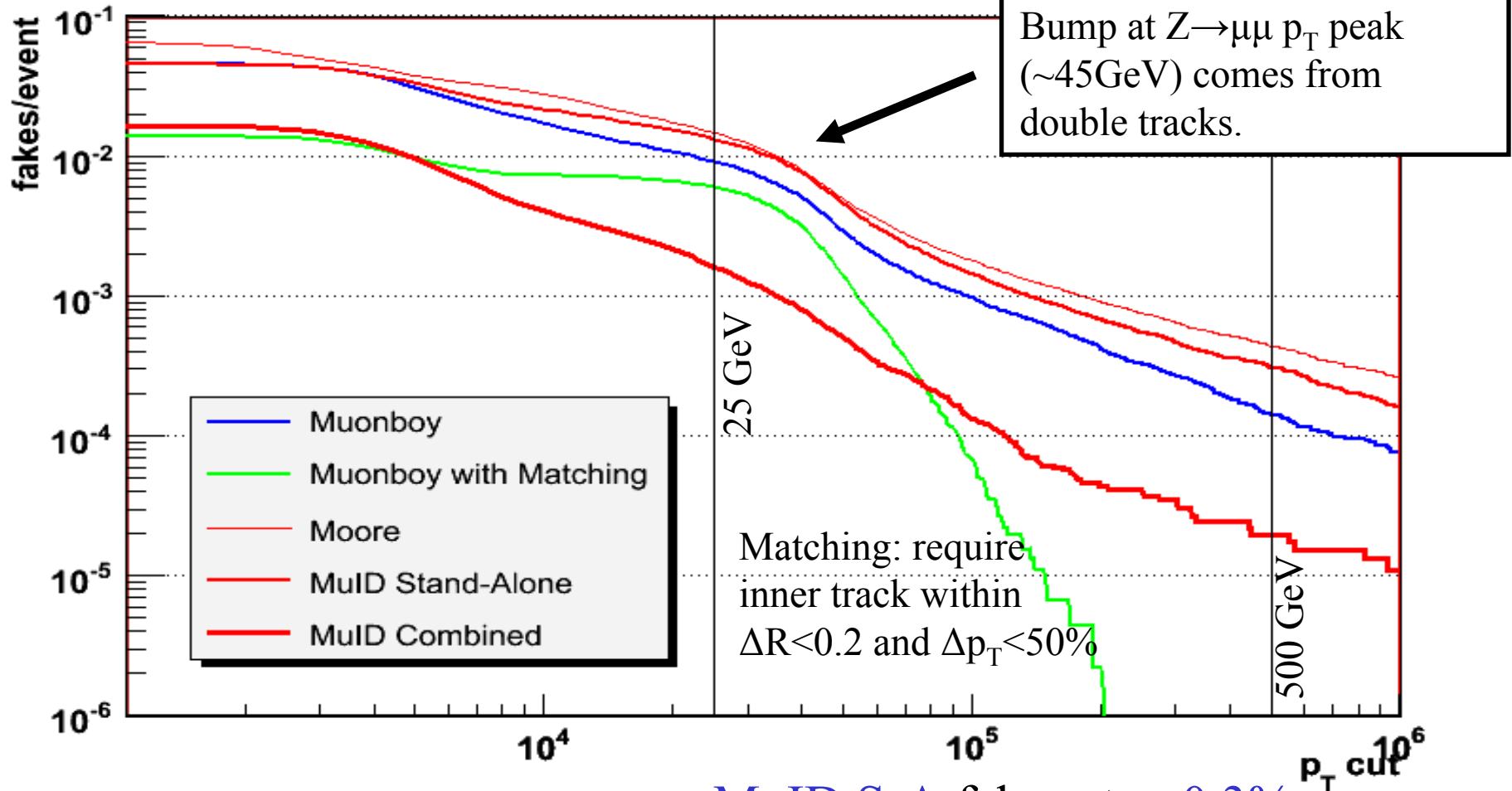




# Fake rate as a function of muon $p_T$ cut



Fakes per  $Z \rightarrow \mu\mu$  Event as  $p_T$  Cut



For  $p_T > 500$  GeV,  
MuonBoy fake rate  $\sim 0.15\%$

MuID S-A fake rate  $\sim 0.3\%$   
MuID Combined fake rate  $\sim 0.015\%$



# Tau Decay Modes



## $\tau$ Decays

### $\tau$ decay modes

#### Leptonical decay modes: ~35%

•  $\tau \rightarrow \nu_\tau + \nu_e + e$

(17.4%)

•  $\tau \rightarrow \nu_\tau + \nu_\mu + \mu$

(17.8%)

#### Hadronical decay modes: ~65%

##### 1 prong

•  $\tau \rightarrow \nu_\tau + \pi^\pm$

(11.0%)

1 track, impact parameter

77% •  $\tau \rightarrow \nu_\tau + \pi^\pm + \pi^0$

(25.4%)

Shower shape, energy

•  $\tau \rightarrow \nu_\tau + \pi^\pm + \pi^0 + \pi^0$

(10.8%)

sharing

•  $\tau \rightarrow \nu_\tau + \pi^\pm + \pi^0 + \pi^0 + \pi^0$

(1.4%)

•  $\tau \rightarrow \nu_\tau + K^\pm + \nu\pi^0$

(1.6%)

##### 3 prong

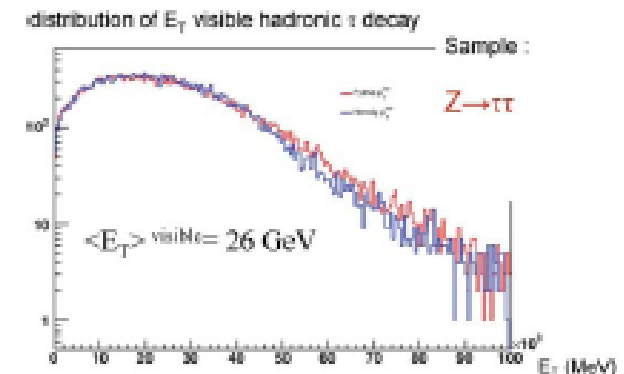
23% •  $\tau \rightarrow \nu_\tau + 3 \pi^\pm + \nu\pi^0$

(15.2%)

3 tracks, impact parameter

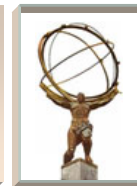
Secondary vertex

Shower shape, energy sharing





# ATLAS Tau Identification



## $\tau$ -jet Reconstruction

- **Tau Candidate Reconstruction**

- Identification

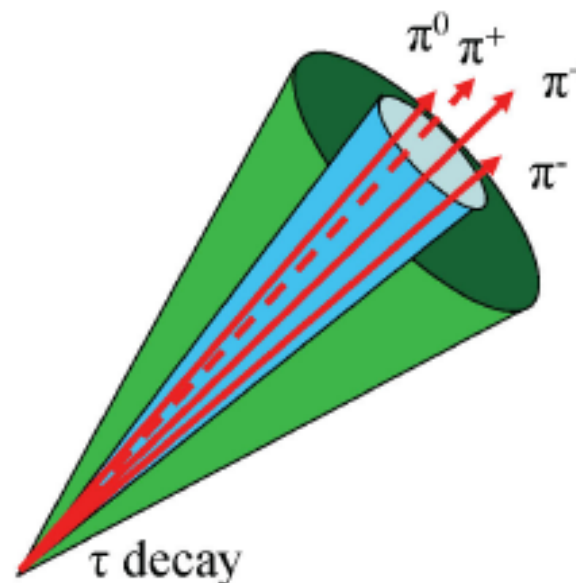
- Well-collimated Calorimeter Cluster with a small number of associated Charged Tracks

- Distinguishing variables

- $R_{em}$  = EM jet radius

- $\Delta E_T^{12}$  = fraction of  $E_T$  in EM/hadronic calorimeters within  $0.1 < \Delta R < 0.2$

- $N_{tracks}$ , Charge, Impact Parameter, Width of the energy deposition in the  $\eta$ -Strips



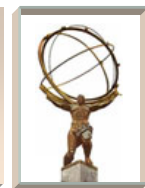
- **Backgrounds misidentified as Taus**

- QCD Jets

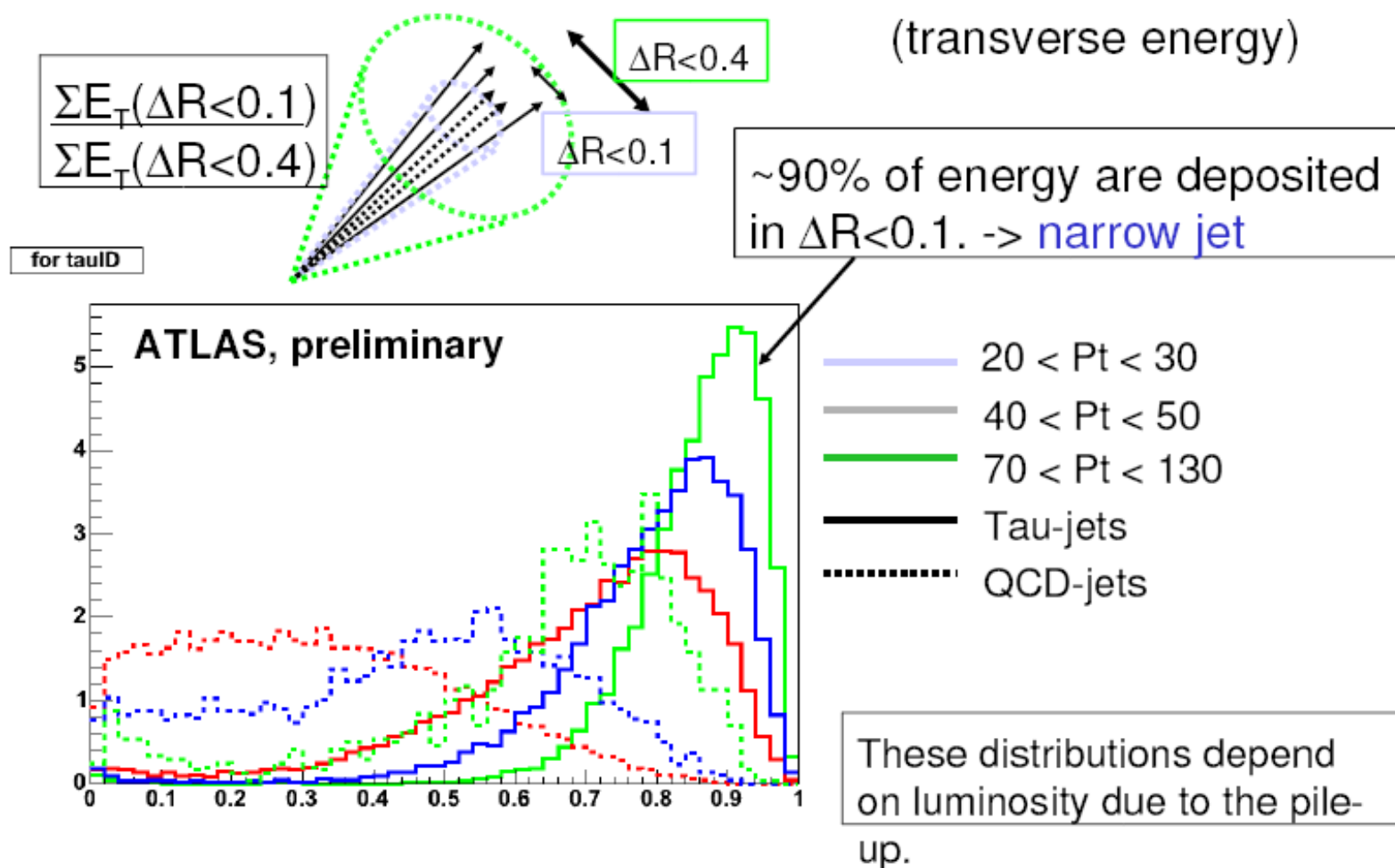
- Electrons that shower late or with strong Bremstrahlung

- Muons interacting in the Calorimeter





# Tau: Fraction of Energy in $\Delta R < 0.1$

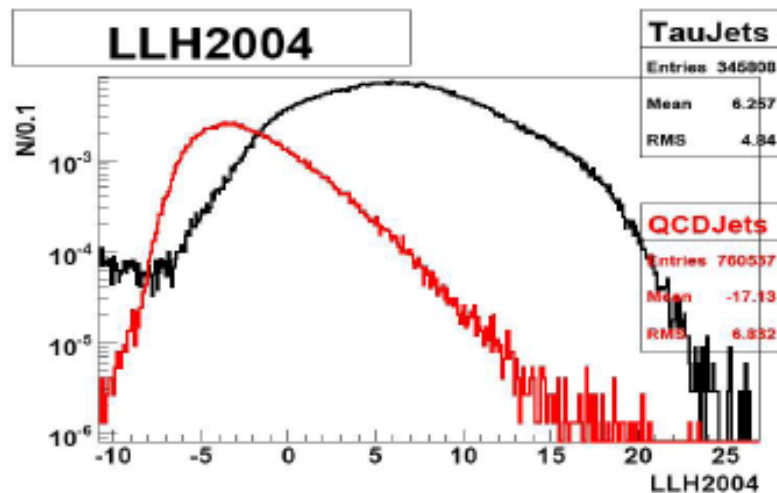




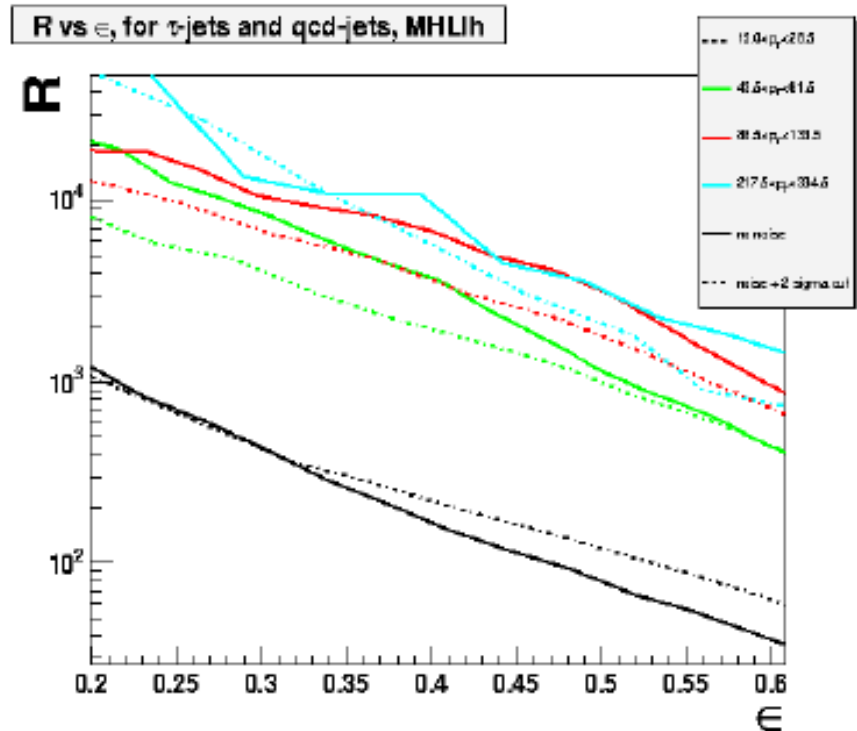
# A Likelihood approach to $\tau$ -ID



- Construct variable that combines all cut variables
- Compare signal and bckgnd
- Can vary cut to get need rejection



ATLAS

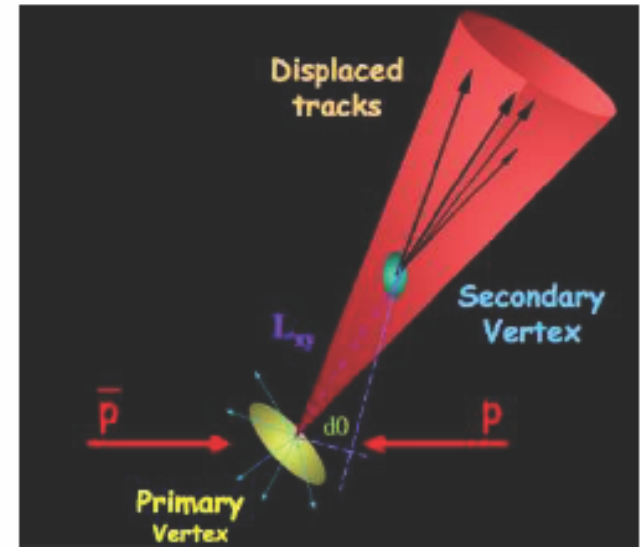




# B-jet Identification

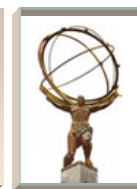


- Characteristics of B decays:
  - B lifetime long:  $c\tau \sim 460 \mu$
  - Semileptonic BR 10% per lepton species
- Two methods of b-tagging
  - Displaced vertex (or track from it)
  - “Soft” leptons close inside jets
- Vertex tagging has higher efficiency and better purity
  - But can combine both techniques

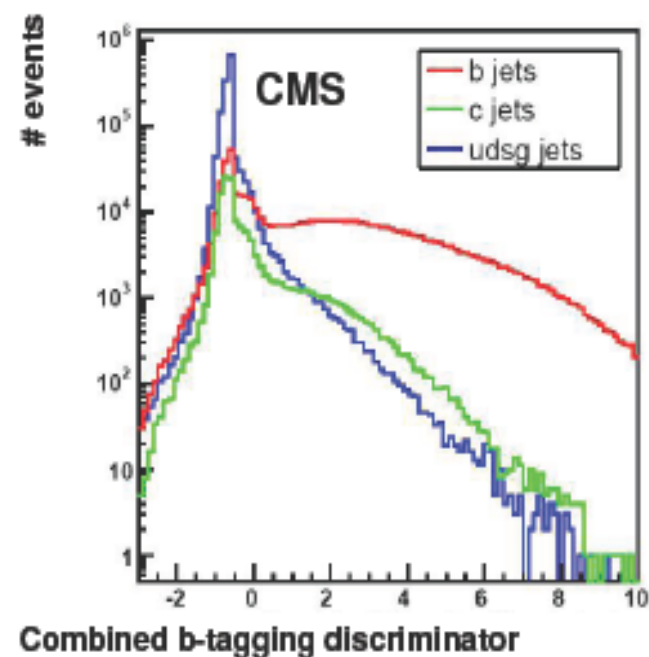
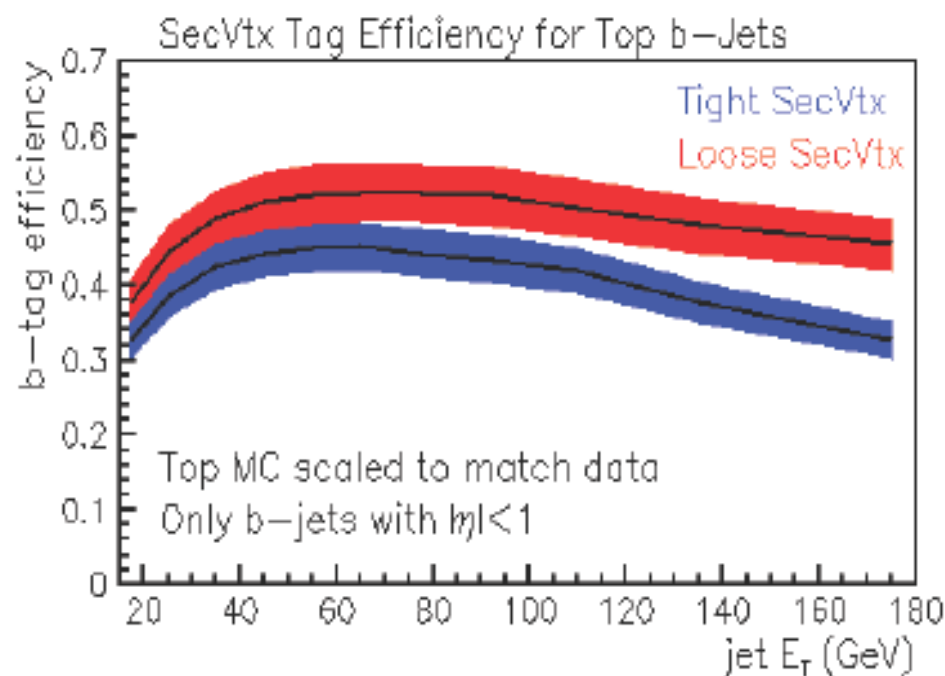




# B-tagging efficiency Depends on Background

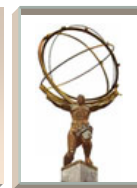


- Charm also long-lived: less rejection
- Performance  $E_T$  dependent

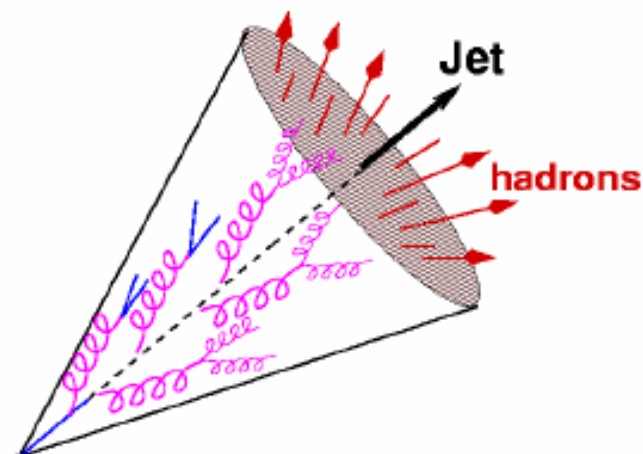
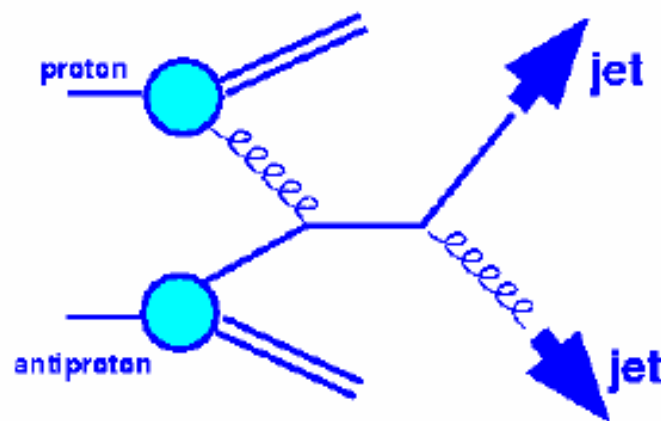




# QCD Jet



- $2 \rightarrow 2$  elastic scattering of quarks and gluons



- Strategy
  - Calorimeter based pattern recognition
  - Associate tracks with the jets after calorimeter jet found
  - Primary vertex needed to calculate  $p_T$



# A ‘Cone’ Jet Reconstruction Algorithm



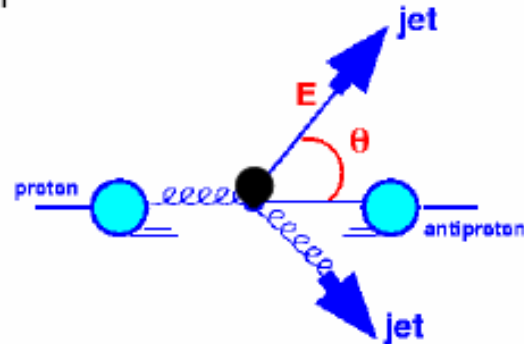
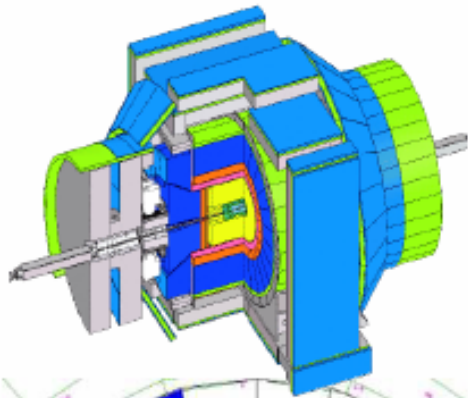
- Jets are circles when projected in  $\eta - \phi$  space
- To reject fluctuations in underlying event and pileup:
  - Start with a “seed” tower above fixed  $E_T$  ( $E_{\text{Seed}}$ )
  - Draw a circle in  $\eta - \phi$  space (Cone Size: 0.4 to 1)
  - Include all towers with above a fixed  $E_T$  ( $E_{\text{tmin}}$ )
  - Calculate  $E_T$  centroid
  - Iterate list of towers until stable
- This is the “pattern recognition” phase



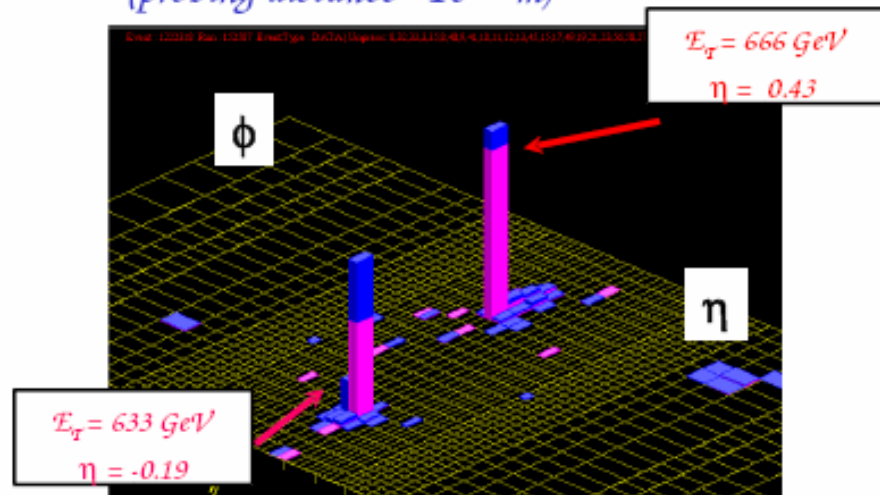
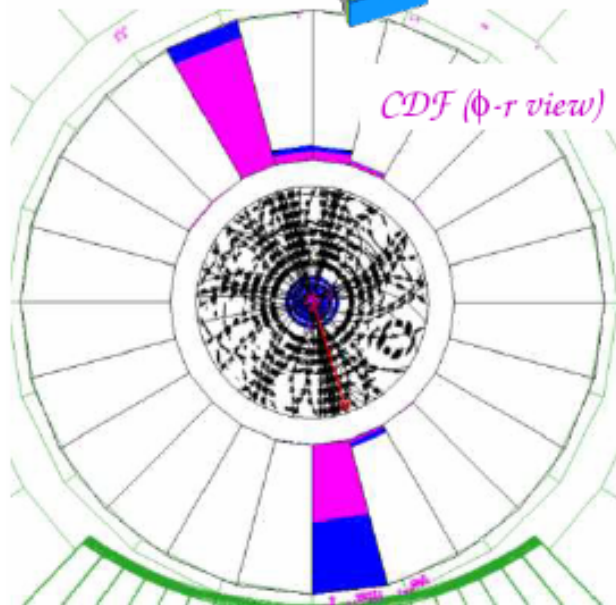
# What Do Jets Look Like ?



(the highest  $P_T$  jet event from CDF)



*Dijet Mass = 1.36 TeV*  
*(probing distance  $\sim 10^{-19}$  m)*





# Jet Energy Scale



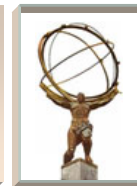
The jet energy scale is the dominant uncertainty in many measurements of the top quark.

CDF and DØ use different approaches to determine the jet energy scale and uncertainty:

- CDF:  
Scale mainly from single particle response (testbeam) + jet fragmentation model  
Cross-checked with photon/Z-jet  $p_T$  balance
  - ~3% uncertainty, further improvements in progress.
- DØ:  
Scale mainly from photon-jet  $p_T$  balance.  
Cross-checked with closure tests in photon/Z+jet events
  - Run II calibration uncertainty ~ 2%



# Missing $E_T$ ( $\nu$ ) Reconstruction



- Use same technique as for jets
  - Create a grid of calorimeter towers
  - Treat each tower as a massless particle with momentum direction normal to the tower



Calorimeter “Tower”

detector



**Define  $\vec{E}_T$  (2 vector)**

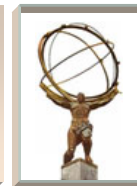
$$\begin{aligned}\vec{E}_T &= -\sum_{\text{Towers}} E_{iT} \hat{n}_i \\ &= -\sum E_i \sin \theta_i \hat{n}_i\end{aligned}$$

**Similarly total  $E_T$**

$$\begin{aligned}E_T &= \sum_{\text{Towers}} |\vec{E}_{iT}| \\ &= \sum |E_i| \sin \theta_i\end{aligned}$$



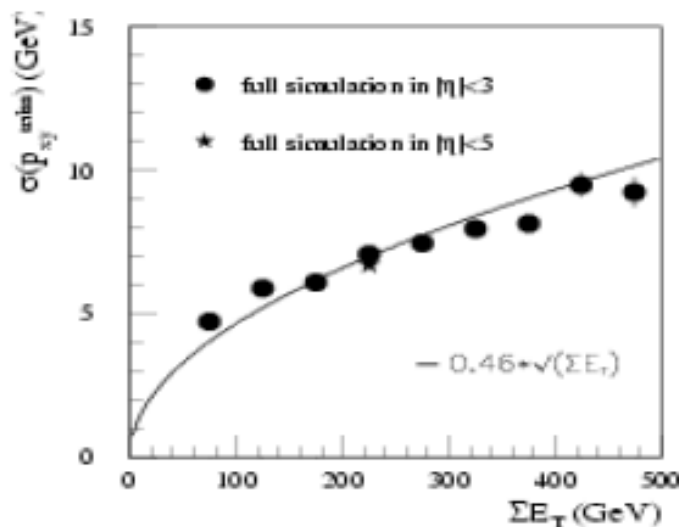
# Missing $E_T$ Resolution



- Calorimeter resolutions depend on energy deposition

$$\sigma_{E_T} \propto \sqrt{E_T}$$

- Measurement is also sensitive to detector imperfections (cracks) and noise
- Degrades with pile-up

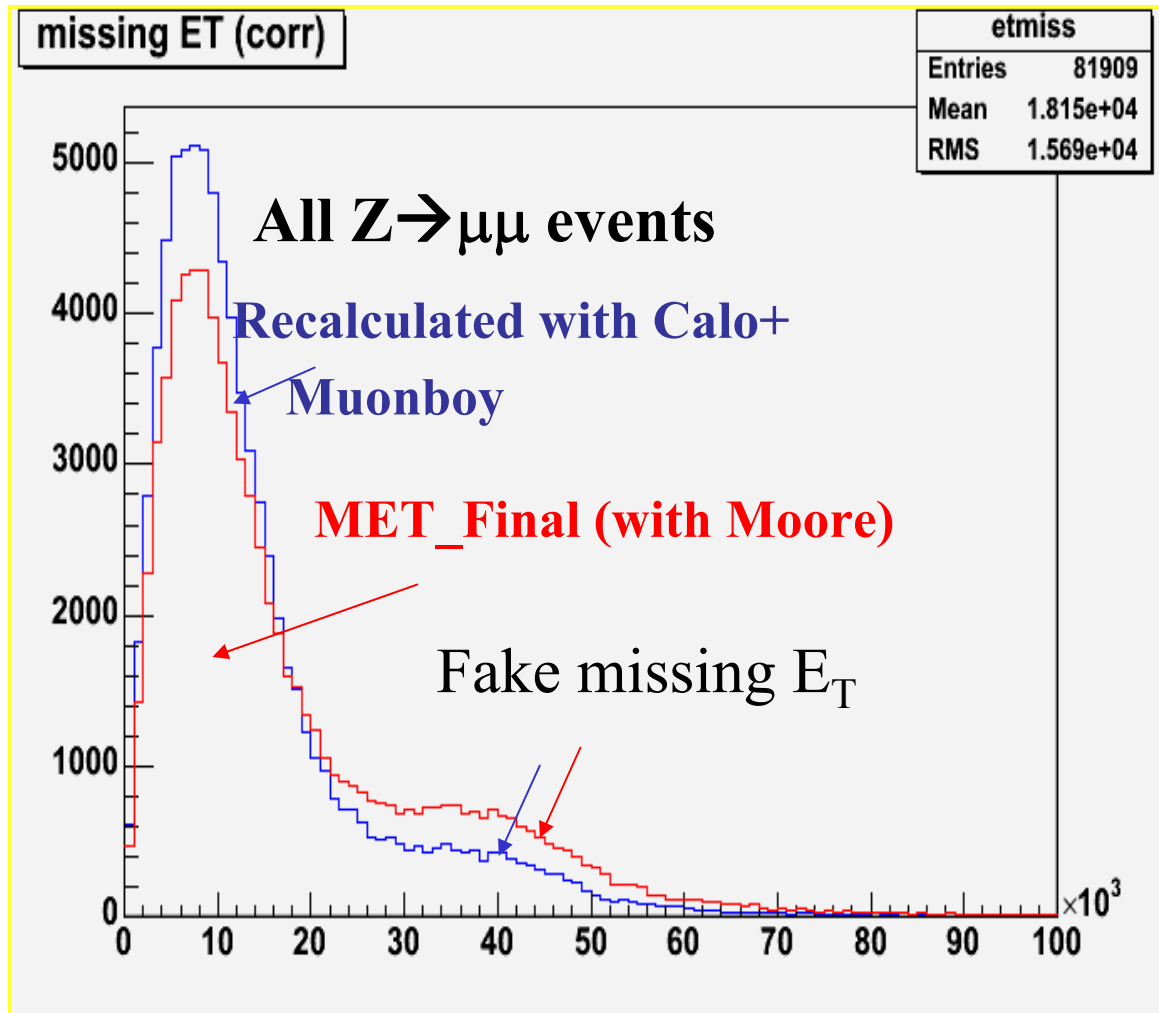


ATLAS TDR



# Comment: Muon ineff. -> Fake Missing $E_T$

(ATLAS:  $Z \rightarrow \mu\mu$  MC sample, recon 10.0.1)

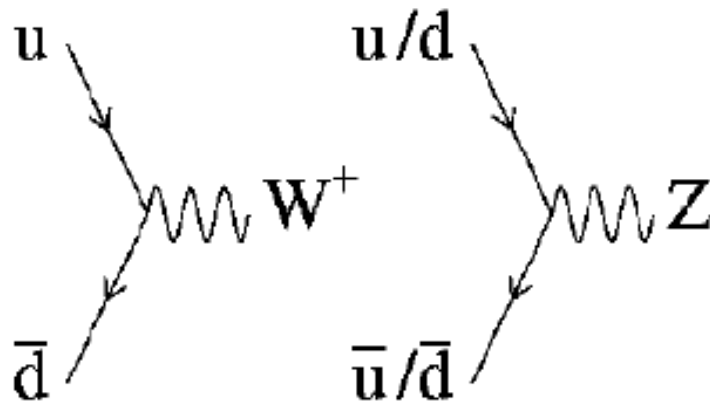




# W and Z Production in Hadron Collider



- Lowest order diagram: quark annihilation
- W and Z obtain  $P_T$  via initial state gluon emission



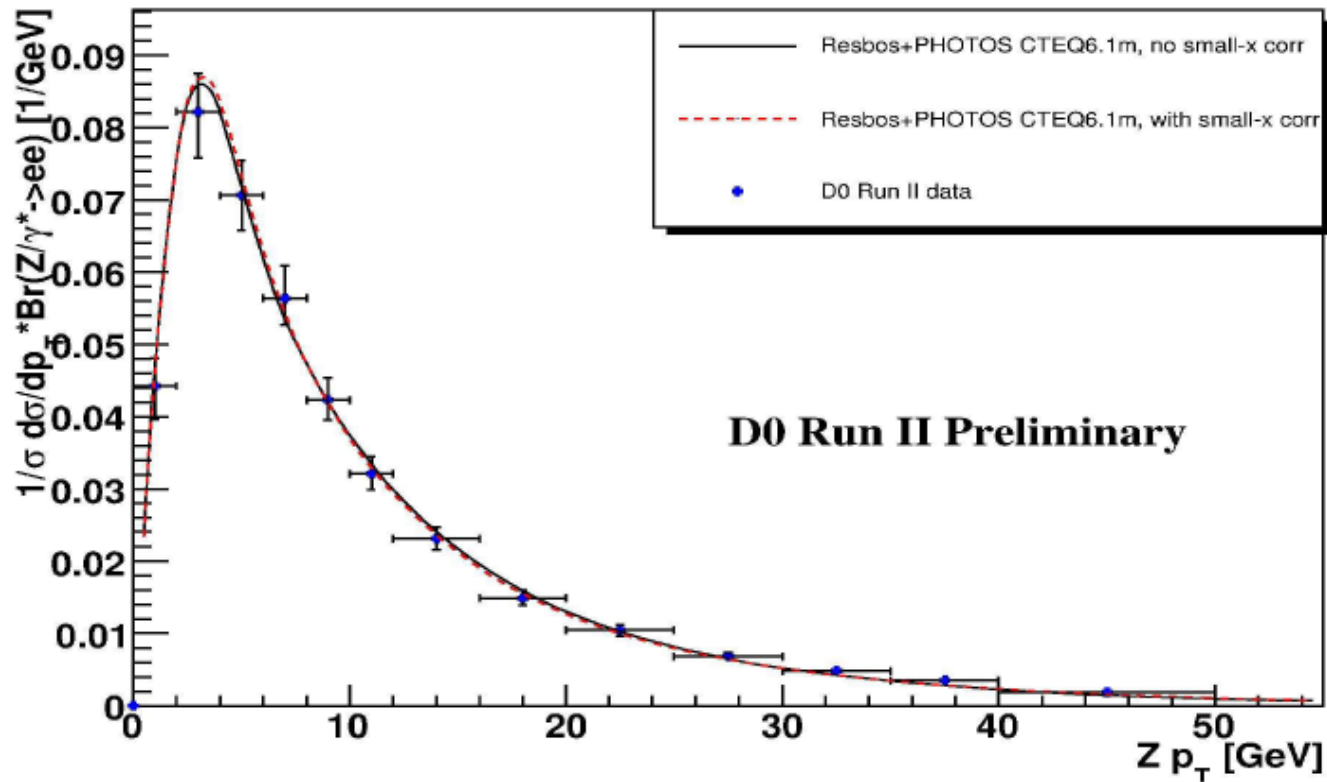
Lowest order production:  
W and Z produced with 0  $P_T$



# Full QCD Calculation: V-Boson $p_T$ Remains Small



Z boson  $p_T$  after unfolding



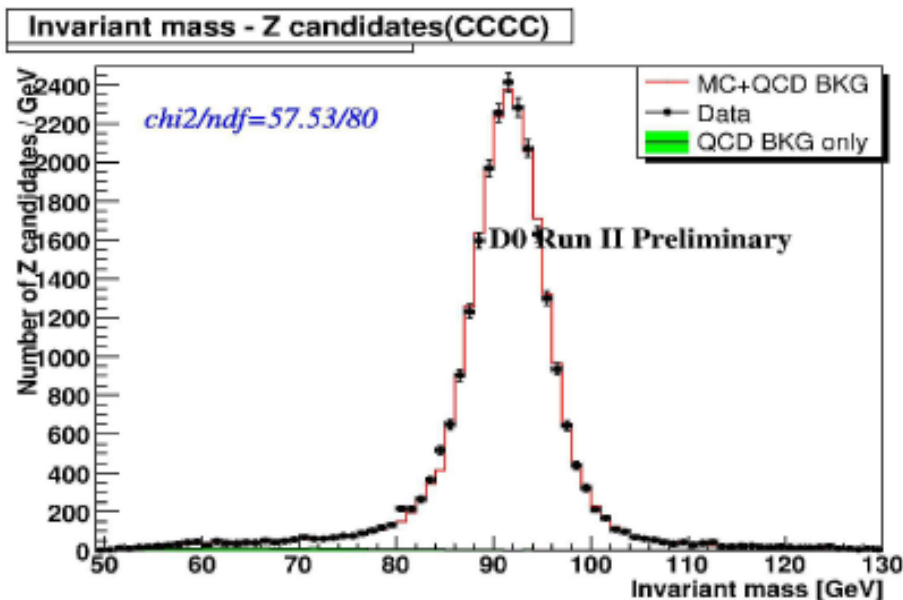
Distribution dominated by multiple soft gluon emission



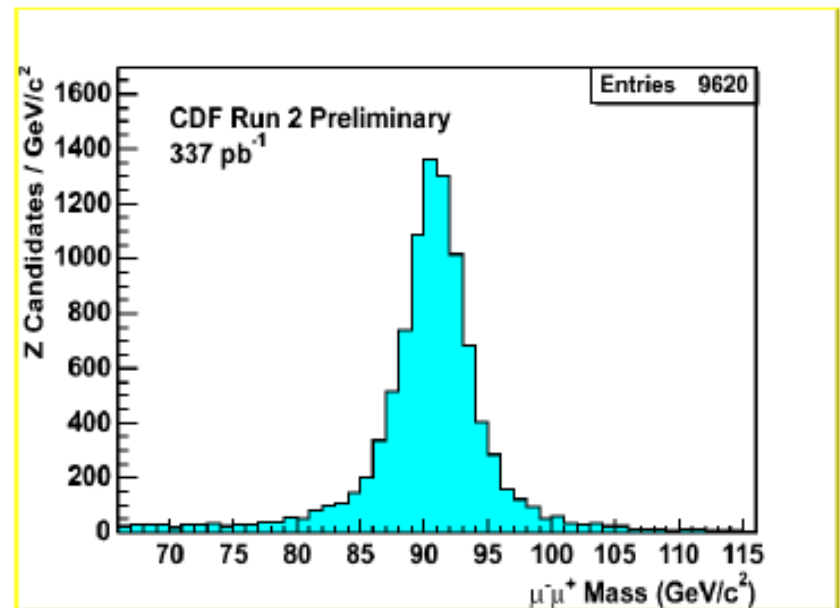
# Reconstruction of Z



- Limited to leptonic modes unless you trigger on b-jets
- Two high  $P_T$  leptons, nearly back-to-back in  $\phi$
- Reconstruction straightforward, background small



$Z \rightarrow ee$



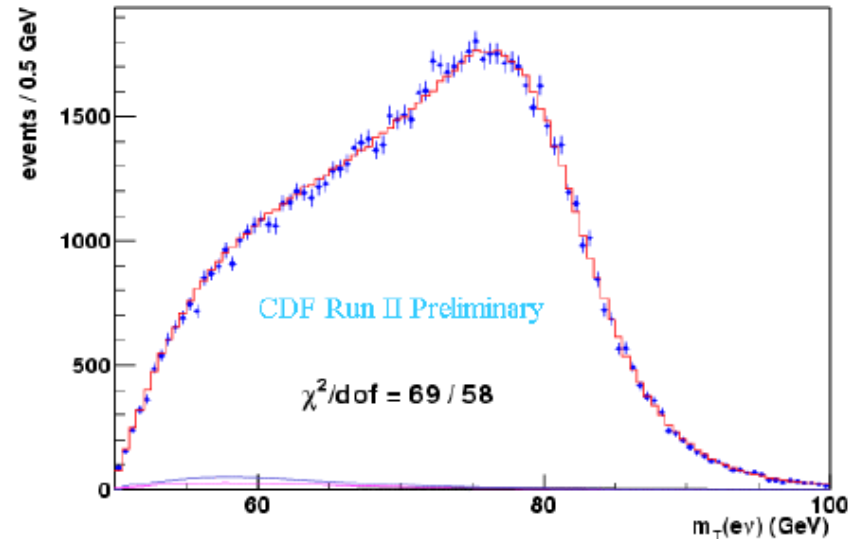
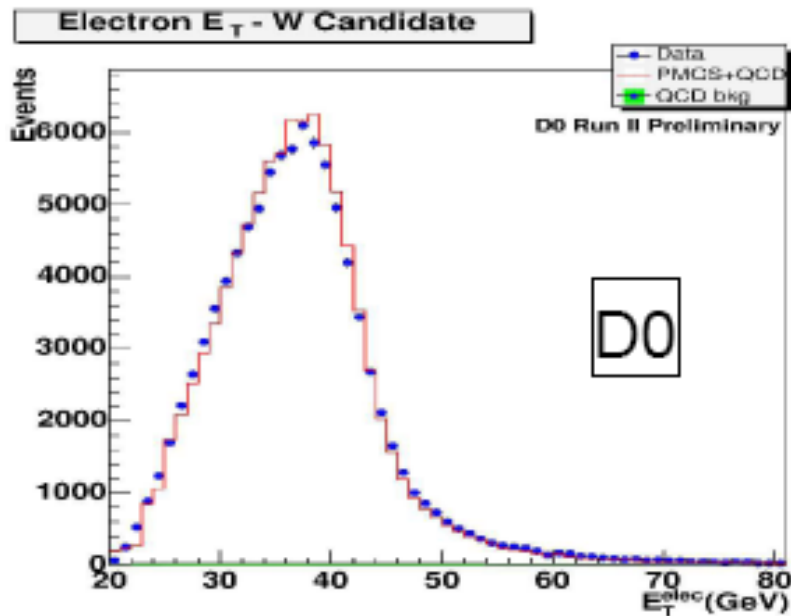
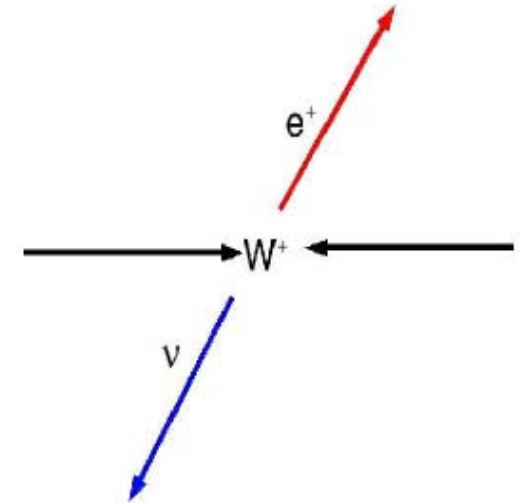
$Z \rightarrow \mu\mu$



# Reconstruction of $W \rightarrow \text{lepton} + \text{neutrino}$



- Select W event:
  - isolated high  $P_T$  e or muon
  - large missing  $E_T$  ( $> 20$  GeV)

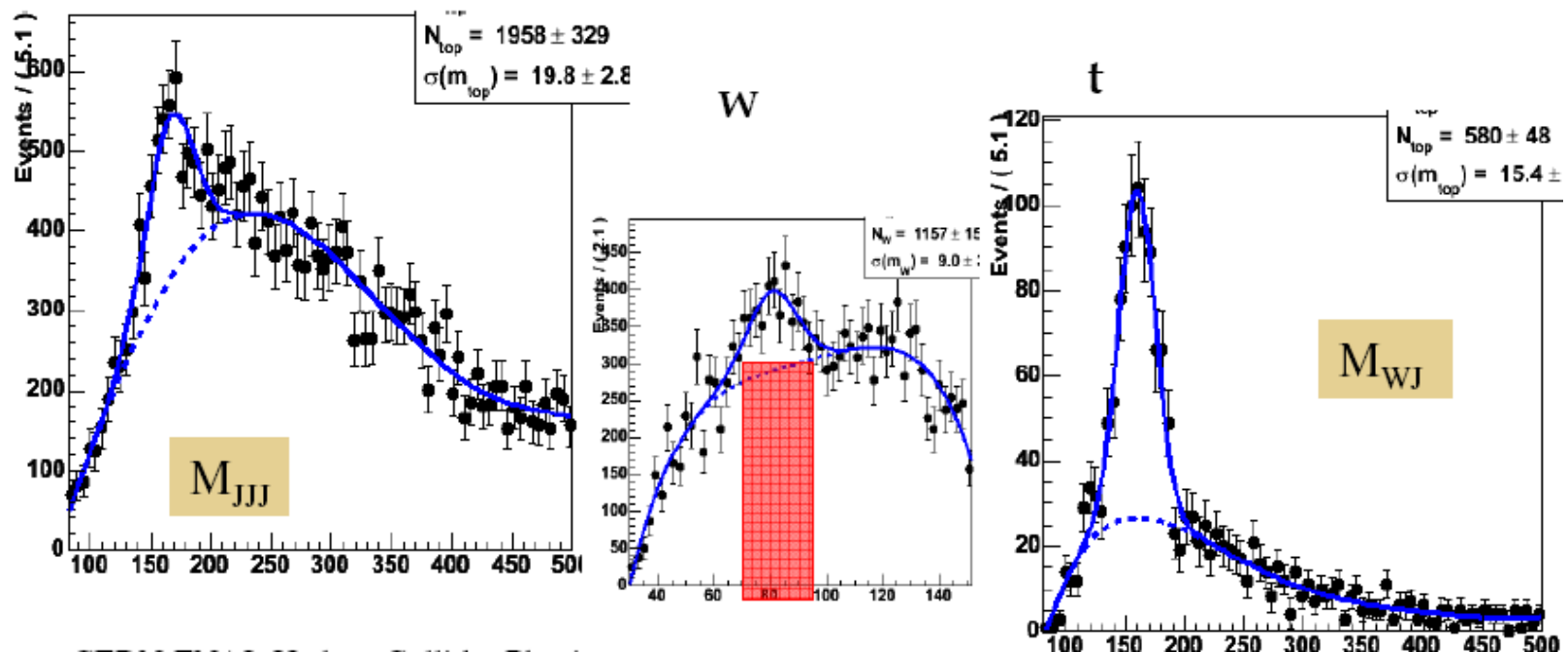




# Top Identification



Three jets with highest vector-sum  $P_T$  as the decay products of the top – lepton trigger. Two jets in hadronic top with highest momentum in reconstructed J+J+J C.M. frame. Lumi = 300 /pb. Top mass with cut on W in  $M_{JJ}$ .





# **Physics Discovery Potential at LHC**

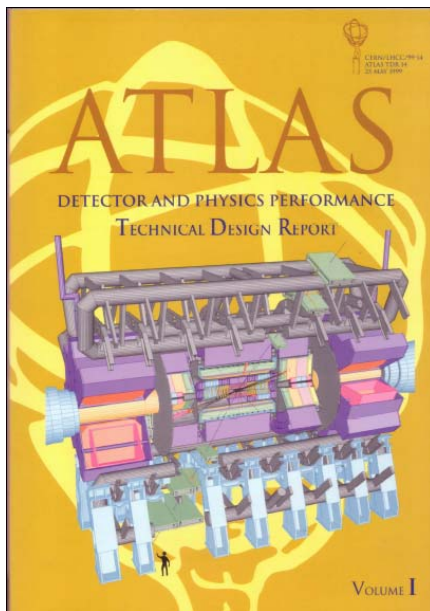
**Bing Zhou**

**The University of Michigan**

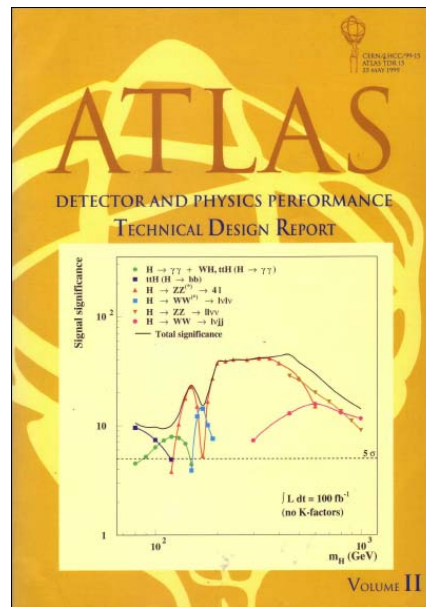
**CCAST Workshop @Tsinghua University**

**Nov. 6-10, 2006**





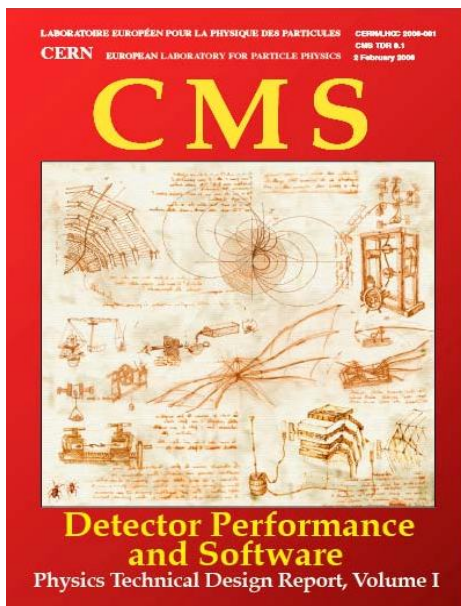
May 1999



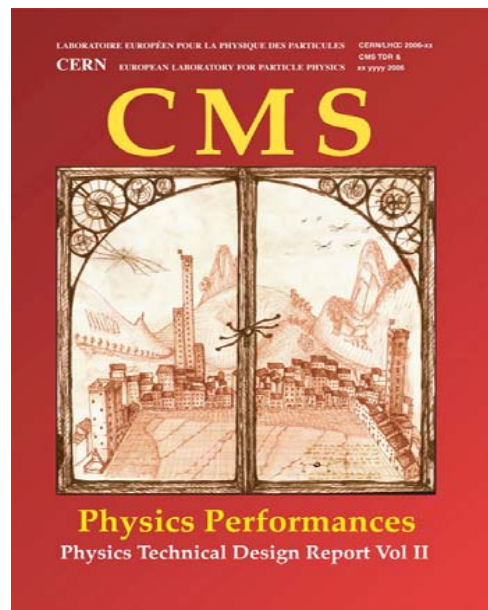
May 1999

## Physics TDRs

CSC (Computing System Commissioning) notes are to be produced in spring 2007, covering software and physics analysis validation for the early physics run with  $0.1 \text{ fb}^{-1}$  and  $1 \text{ fb}^{-1}$ .



Feb. 2006



Jun. 2006

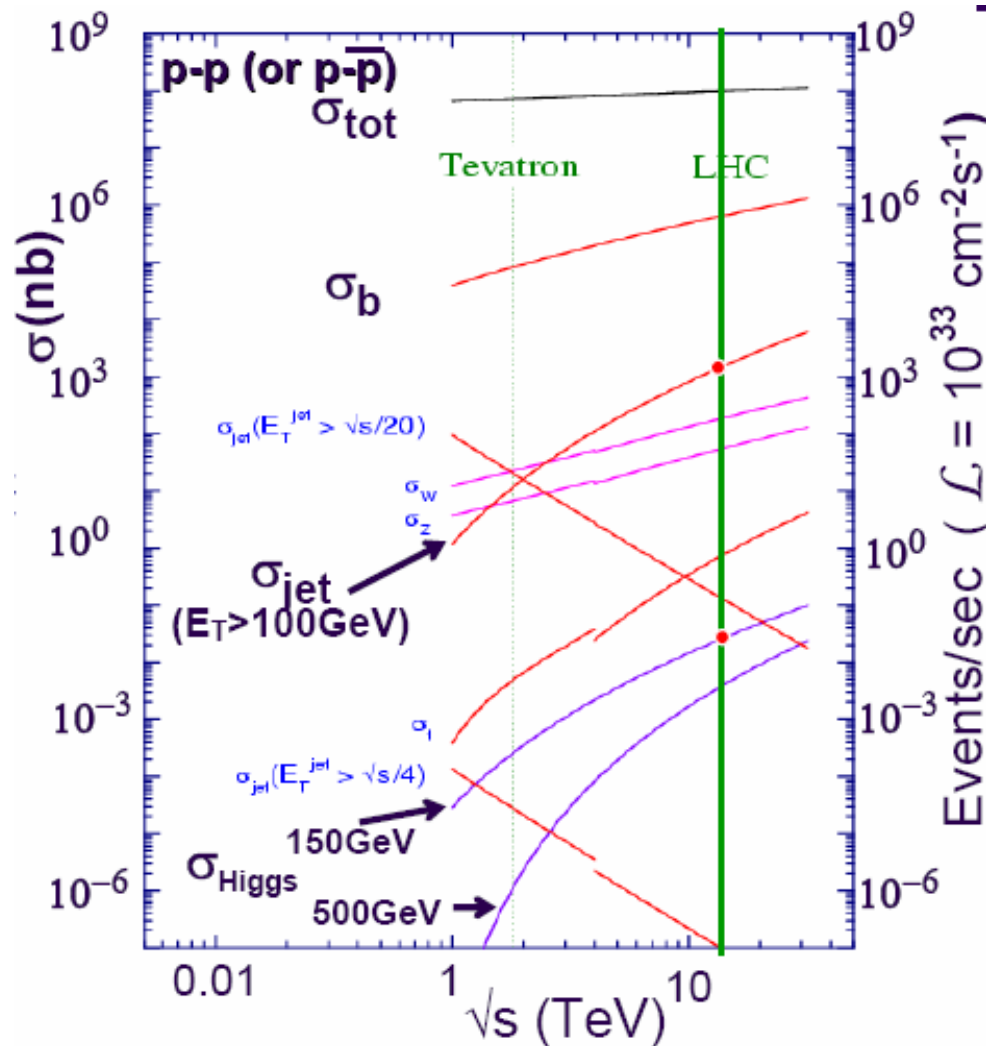
Instead of 3-rd vol. of TDR, short notes on startup will be submitted to LHCC in summer 2007, along with the very early physics reach with  $0.1 \text{ fb}^{-1}$  and  $1 \text{ fb}^{-1}$ .



# First Step of Campaign at the LHC



## 1) MC Models and Re-establish the SM processes



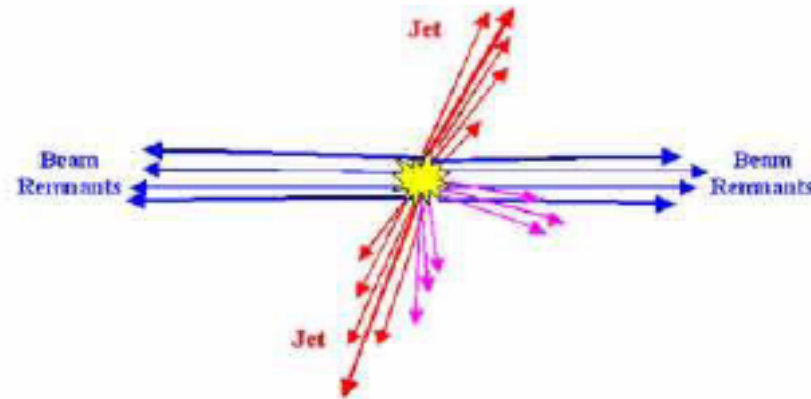
Channel	Recorded [1 fb <sup>-1</sup> ]
$W \rightarrow \mu\nu$	$7 \times 10^6$
$Z \rightarrow \mu\mu$	$1 \times 10^6$
$t\bar{t} \rightarrow \mu + X$	$0.1 \times 10^6$
Jets $p_T > 150 \text{ GeV}$ (if 10% bandwidth)	$\sim 10^6$
Min Bias (10% bandwidth)	$\sim 10^6$ (can be larger)
$gg$ ( $M \sim 1 \text{ TeV}$ )	$10^2 - 10^3$



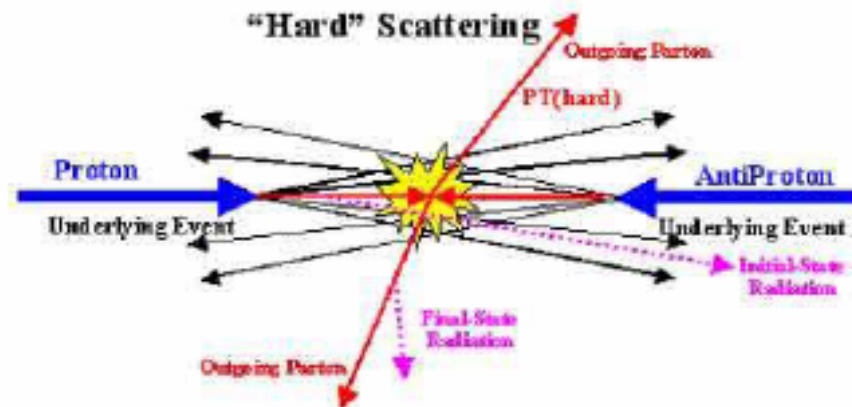
# Underlying Event and ISR Modeling



- Hard Collision leaves remnants of incoming p's moving in Beam Direction



- “Initial State” gluon radiation largely co-linear with incoming partons: same basic structure



Soft particles distributed  
uniformly in  $\eta$

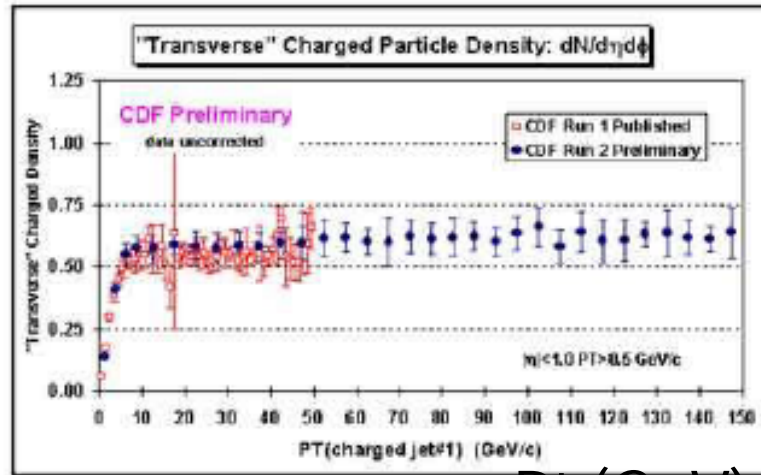


# Track Distributions from Underlying Events



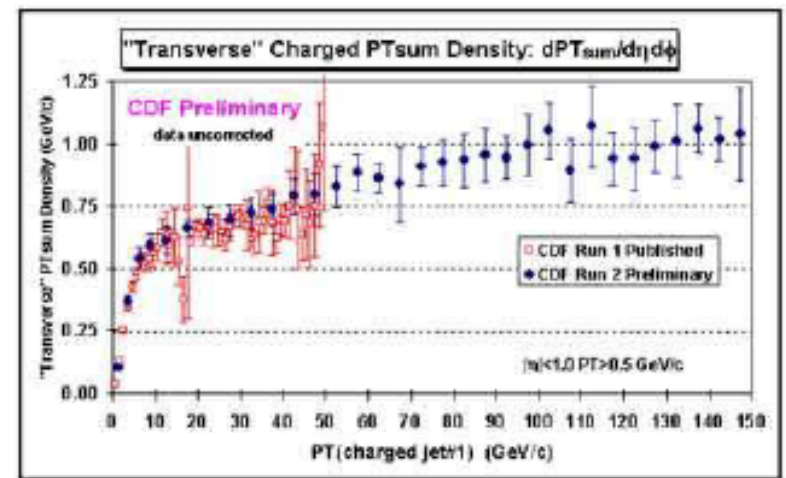
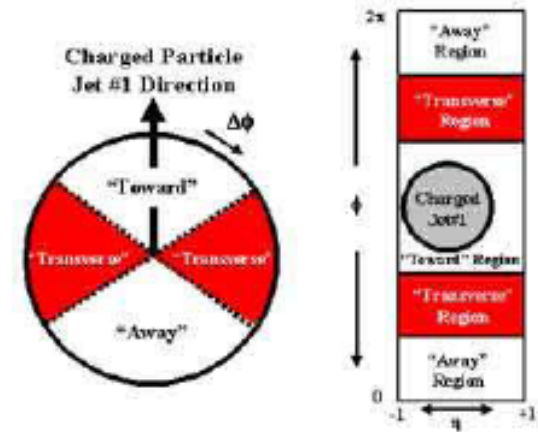
- Look at  $90^\circ$  from jet direction

$dN/d\eta d\phi$



Pt (GeV)

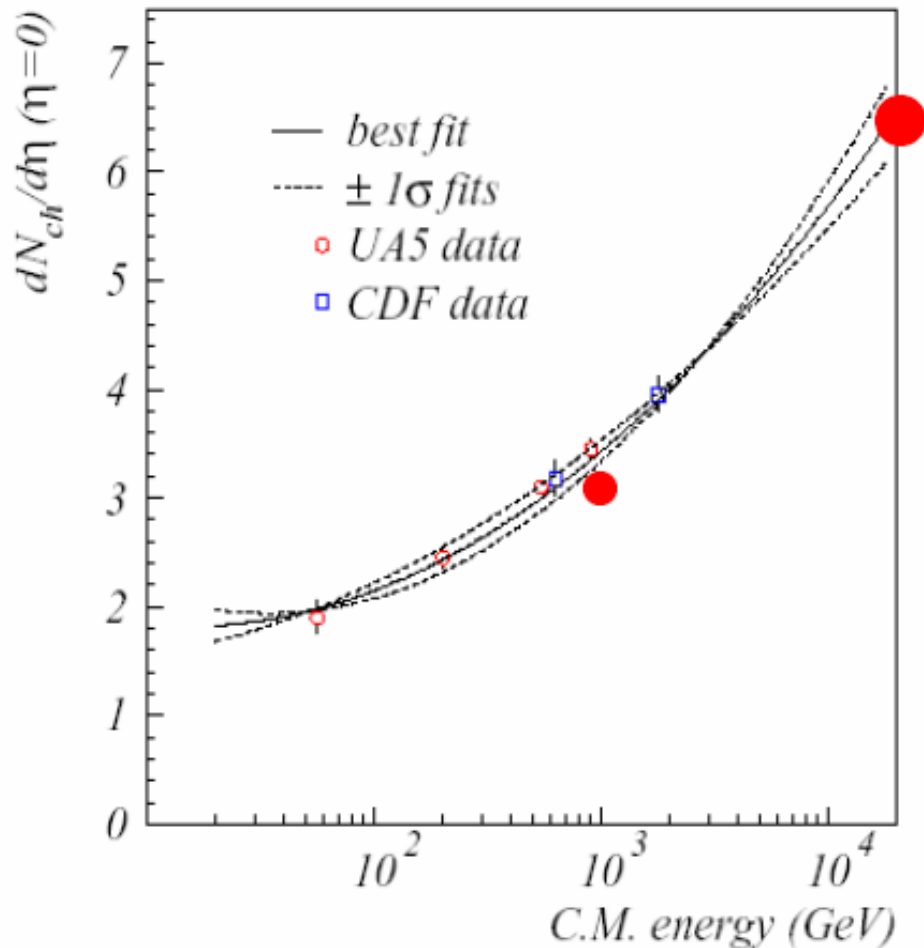
- Approx constant Particle multiplicity
- Energy density increases with hard scattering scale



Pt (GeV)



# MinBias Density Extrapolate



Using data from 0.2 to 1.8 TeV to extrapolate the plateau rapidity density. For all pions expect

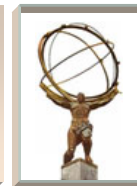
$$\rho = 1/\sigma(d\sigma/d\eta) \sim 9$$

$$\pi^+ = \pi^- = \pi^0$$

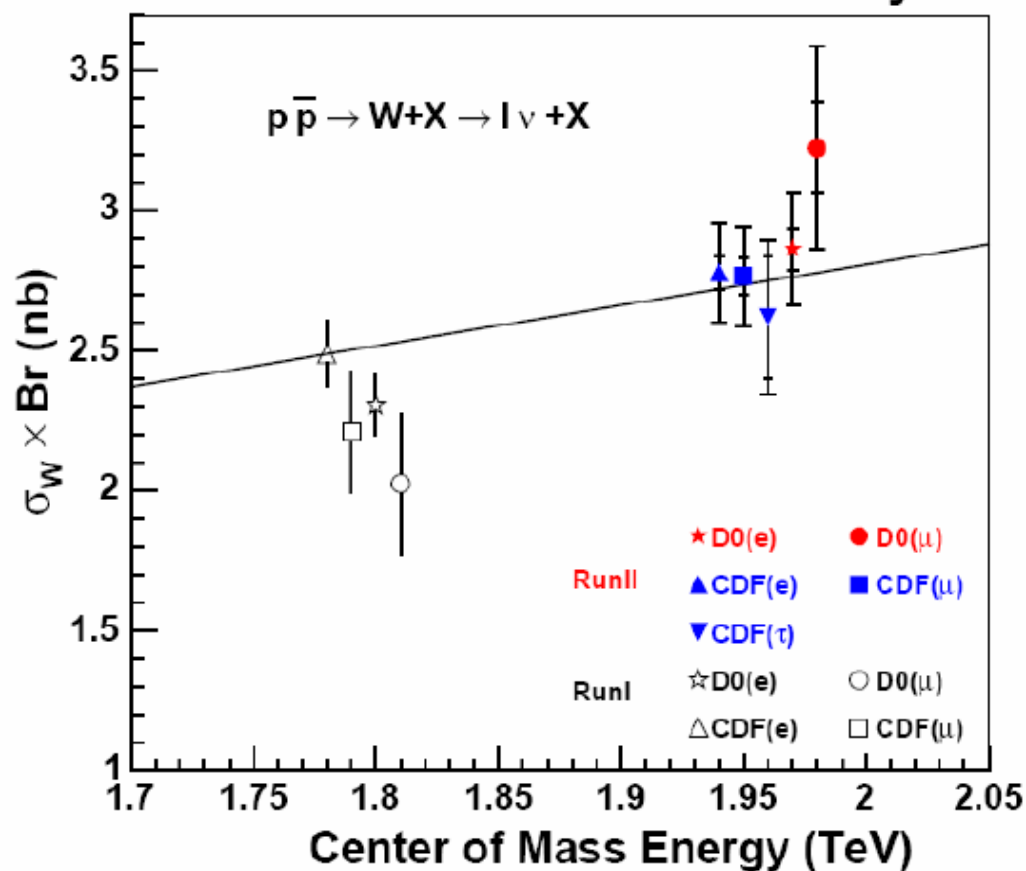
**Note - 2x  
extrapolation from  
0.9 TeV**



# Standard W Candle



**CDF and D0 RunII Preliminary**



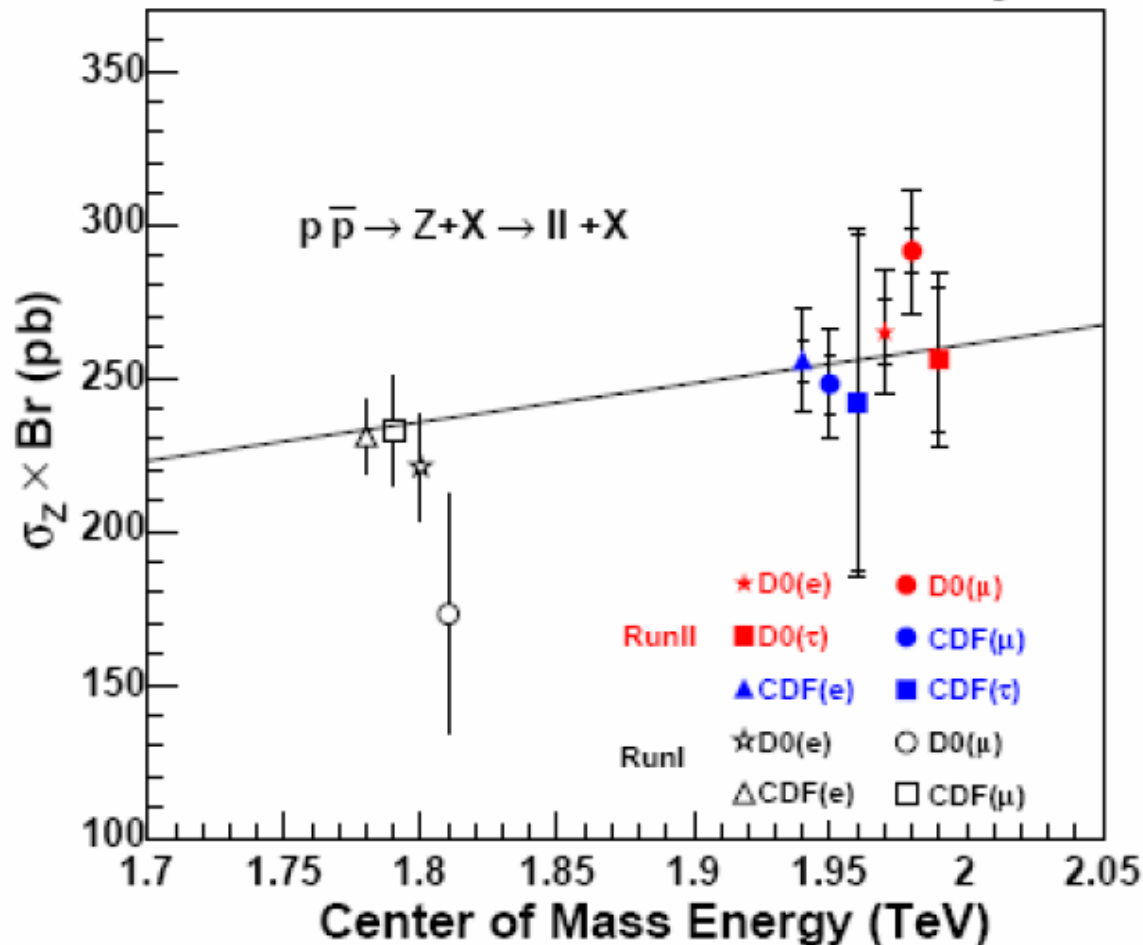
Use  $W \rightarrow \mu + \nu$  as a “standard candle” to set the LHC luminosity? Expect  $\sim 2\%$  accuracy on the predicted cross section. Cross section for  $W \rightarrow \mu + \nu$  with  $|\eta| < 2.5 * Pt > 15$  GeV is  $\sim 10$  nb. In isolated muon triggers look for MET and for Jacobean peak indicating cleanly identified W D-Y production. Once established, look at transverse mass tail in isolated leptons. In new territory above Run II mass reach – start a discovery search in 2008.



# Standard Z Candle



CDF and D0 RunII Preliminary



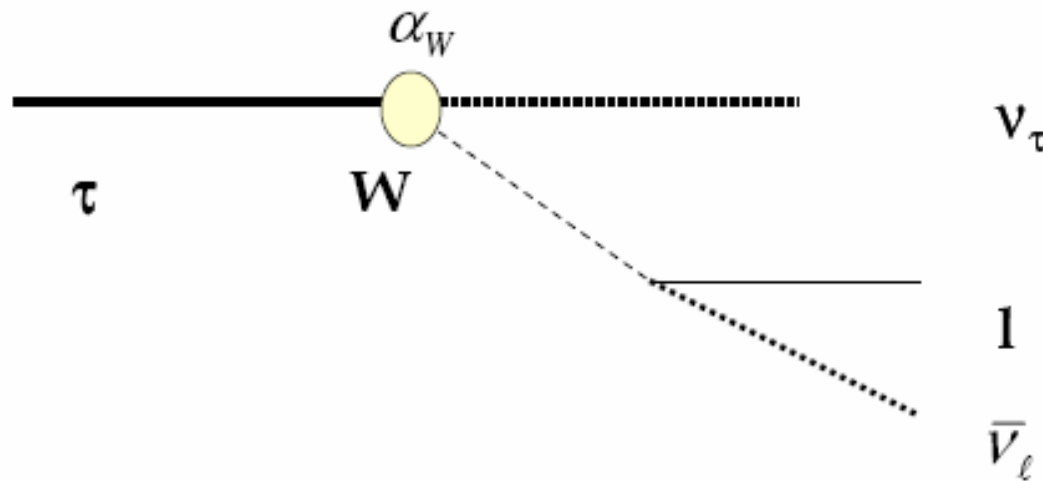
Use  $Z \rightarrow \mu + \mu$  as second “standard candle” to determine the LHC luminosity. Expect  $\sim 2\%$  accuracy in cross section prediction. Find cross section for  $Z \rightarrow \mu + \mu$  decay with  $|\eta| < 2.5$  \*  $P_t > 15$  GeV is  $\sim 600$  pb. Note slow rise from Run II cross section values.



# Establish $Z \rightarrow \tau\tau$



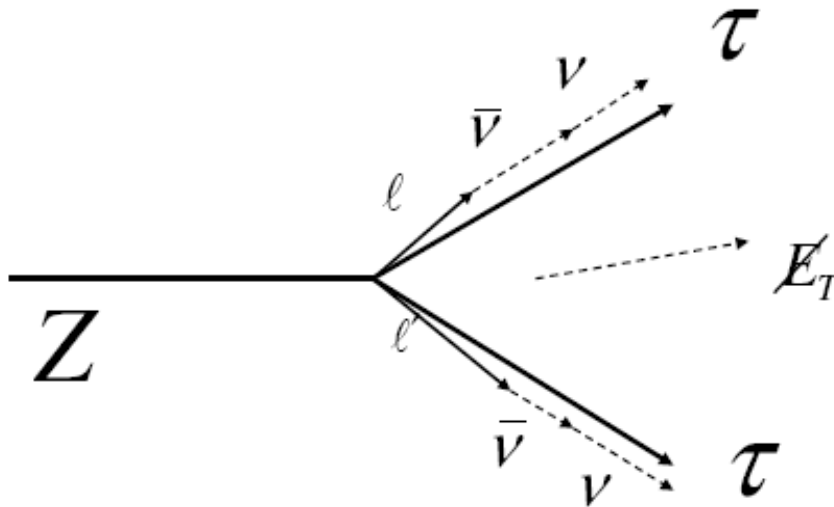
$Z \rightarrow \tau + \tau$  decays as with muon and electron pairs.



Decays appear in dilepton trigger stream. BR is 35% for tau into muon or electron.



# $Z \rightarrow \tau\tau$ : Colinear



$$M_{\tau\tau} \sim \sqrt{2(E_{\ell 1} + E_{\nu\nu 1})(E_{\ell 2} + E_{\nu\nu 2})(1 - \cos \theta_{\ell 1 \ell 2})}$$

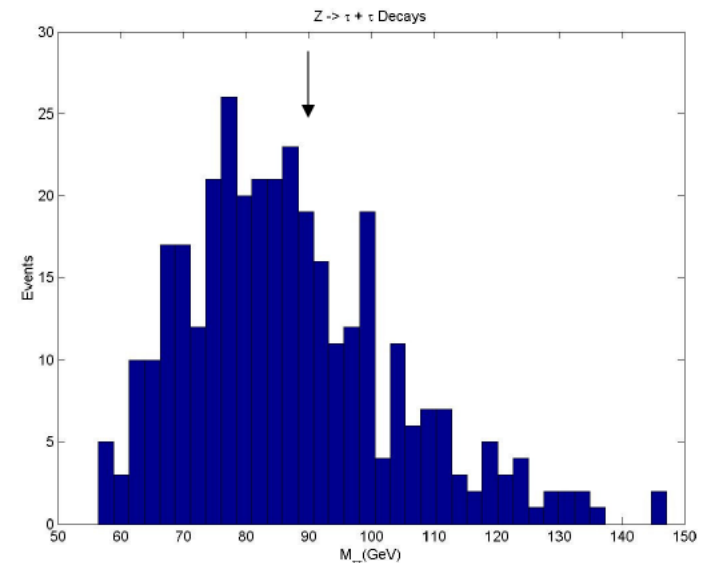
$$MET_x = E_{\nu\nu 1} \alpha_{x\ell 1} + E_{\nu\nu 2} \alpha_{x\ell 2}$$

$$MET_y = E_{\nu\nu 1} \alpha_{y\ell 1} + E_{\nu\nu 2} \alpha_{y\ell 2}$$

Assume collinear neutrinos. Then have 2 Eqs in 2 unknowns. Must cut on determinant

$$\det = \sin \theta_{\ell 1} \sin \theta_{\ell 2} \sin(\phi_{\ell 1} - \phi_{\ell 2})$$

$|\det| > 0.005$  is  $\sim 70\%$  efficient after cuts on Pt of the leptons and MET.





# Establish SM as a Baseline in Early LHC (10 – 500 pb<sup>-1</sup>)



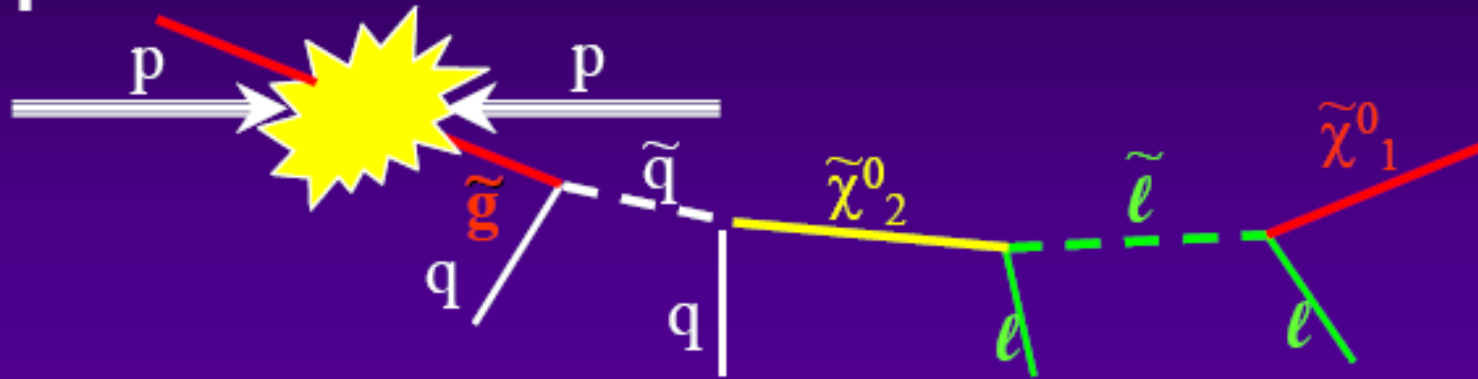
- Use SM processes to understand the detector and calibrate detector in early running
  - Lepton energy scale from Zs
  - Jet energy calibration from g-j; b-j
  - Luminosity from W+Z rates
  - Tuning of Monte Carlos
- Understand SM background processes
  - Minimum bias and underlying event
  - Establish limits of knowledge of SM processes
- Aim to study SM Physics topics
  - Cross section measurements (W, Z, Di-jets, Di-boson)
  - Precision EW measurements:  $M_W$ ,  $\sin^2\theta_w$  and gauge couplings
  - QCD measurements:  $\alpha_s(Q^2)$  and PDFs



# Quick Discovery of SUSY Particles?



## Typical SUSY event at LHC:



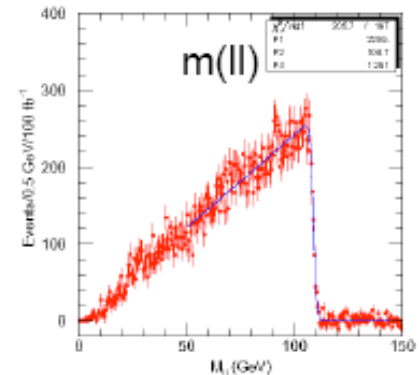
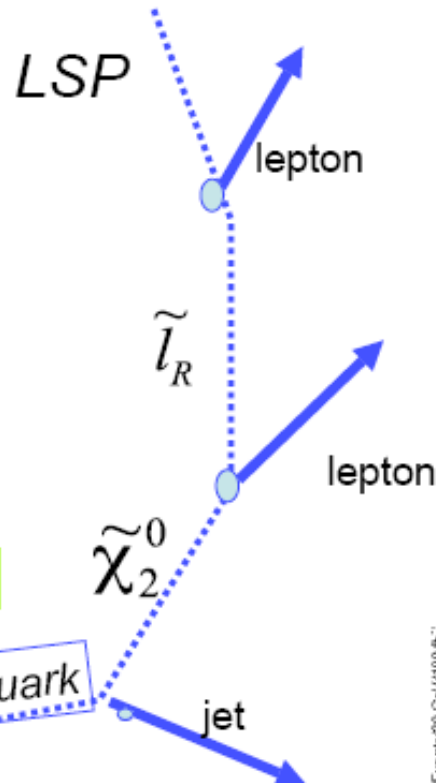
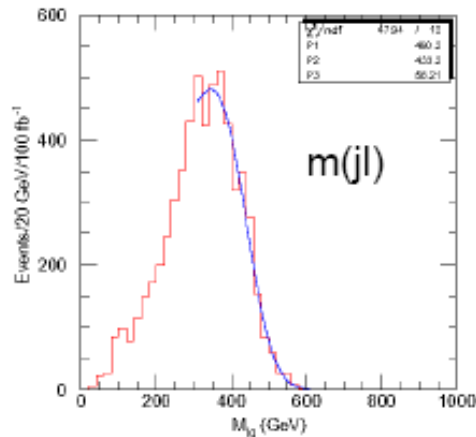
- ◆ Strongly interacting sparticles (squarks, gluinos) dominate production
  - ◆ Can have high cross-sections  $\Rightarrow$  good candidate for early discovery
- ◆ sleptons, gauginos etc.  $\tilde{g}$  cascade decays to LSP.
- ◆ Long decay chains and large mass differences between SUSY states
  - ◆ Many high  $p_T$  objects observed (leptons, jets, b-jets).
- ◆ If R-Parity conserved LSP stable and sparticles pair produced.
  - ◆ Large  $E_{T\text{miss}}$  signature
- ◆ Closest equivalent SM signature  $t \rightarrow Wb$  with  $W \rightarrow \ell \nu$



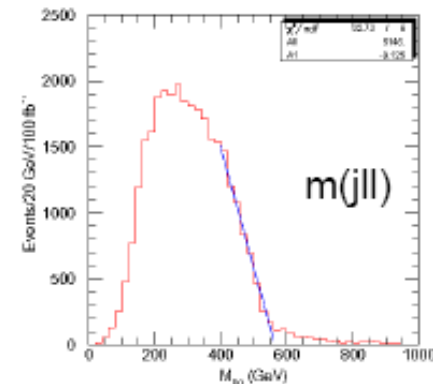
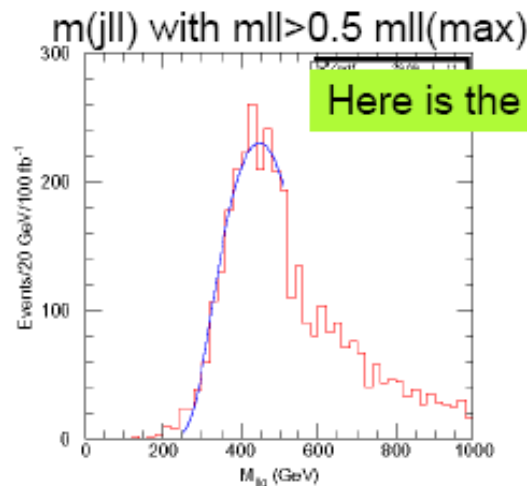
# Mass Reconstruction



determination of the boundary of phase space  
for mass determination.

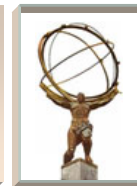


$ee+\mu\mu-e\mu$  subtraction  
is effective to select  
single channel

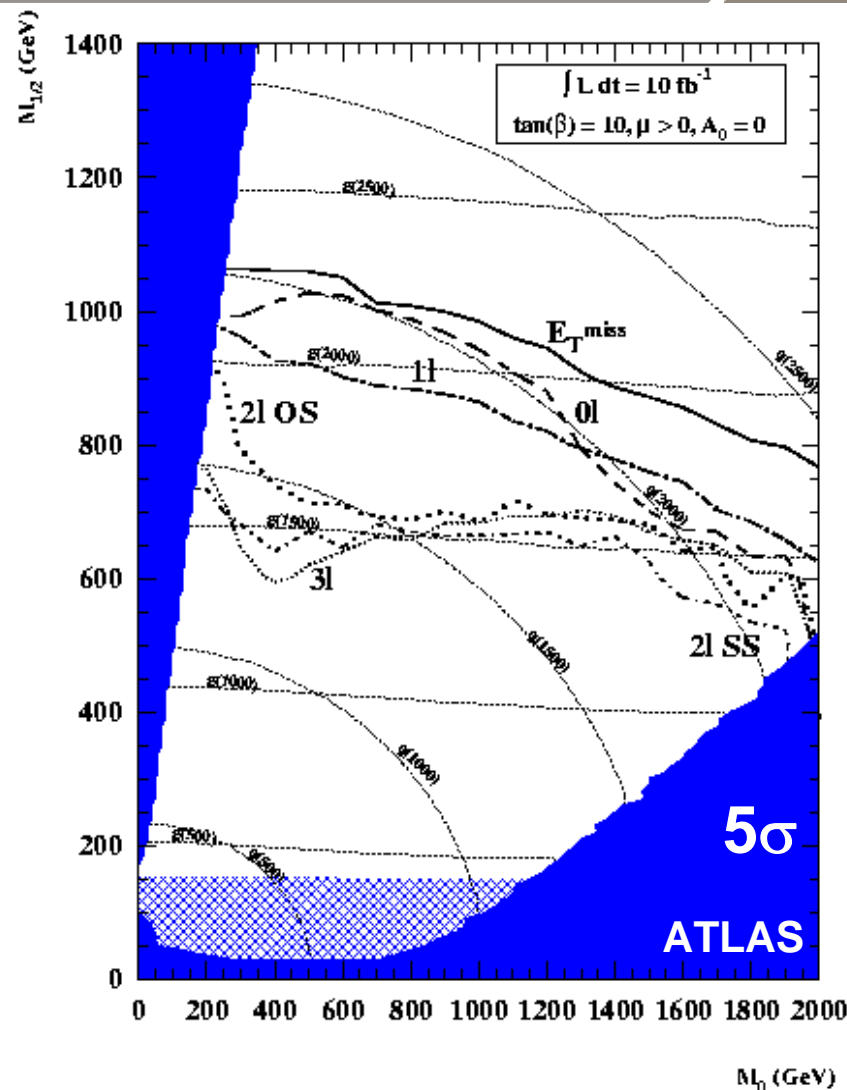




# SUSY Search (ATLAS)



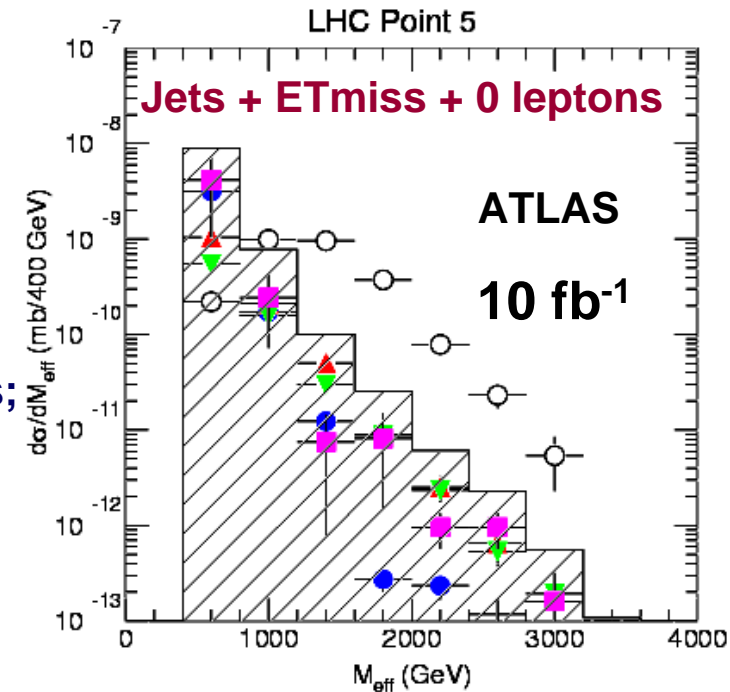
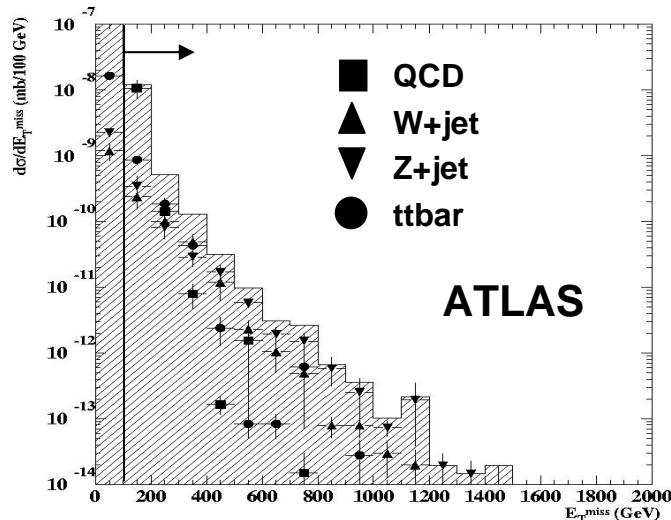
- R-Parity conserving SUSY search channels:
  - Large  $E_T^{\text{miss}}$ ;
  - Large jet multiplicity;
  - Large  $E_T^{\text{sum}}$ .
- Will need convincing estimates of backgrounds with as little data as possible.
- Background estimation techniques will change depending on integrated lumi.
- Ditto optimum search channels & cuts.
- Aim to use combination of
  - Fast/'brisk'-sim;
  - Full-sim;
  - Estimations from data.
- Use comparison between different techniques to validate estimates and build confidence in (blind) analysis.





# Strategy

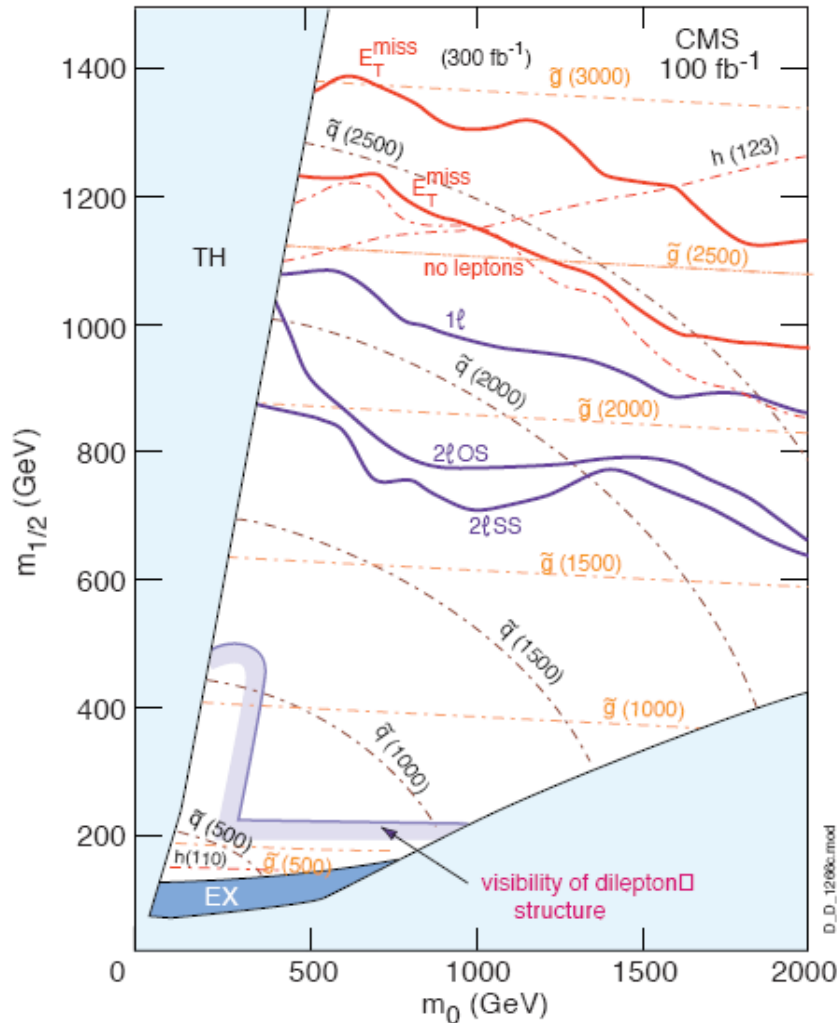
- **Main backgrounds:**
  - Z + n jets
  - W + n jets
  - ttbar
  - QCD
- **Also:**
  - Single top
  - WW/WZ/ZZ
- **Generic approach :**
  - Select low  $E_T^{\text{miss}}$  background calibration samples;
  - Extrapolate into high  $E_T^{\text{miss}}$  signal region.



- **Used by CDF / D0**
- **Extrapolation non-trivial.**
  - Must find variables uncorrelated with  $E_T^{\text{miss}}$
- **Several approaches developed.**



# SUSY Discovery Potential (CMS)

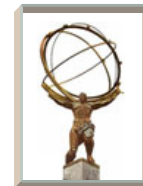


CMS  $5\sigma$  reach for **R-parity conserving mSUGRA** in various inclusive channels:

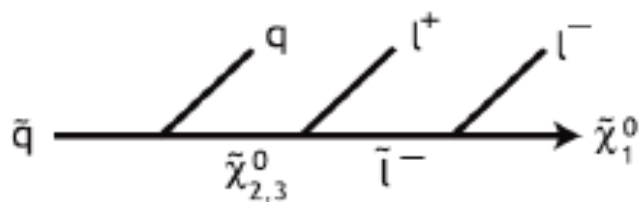
- $\cancel{E}_T$
- $\cancel{E}_T$  with lepton veto
- One lepton channel
- Two opposite sign (OS) leptons
- Two same sign (SS) leptons



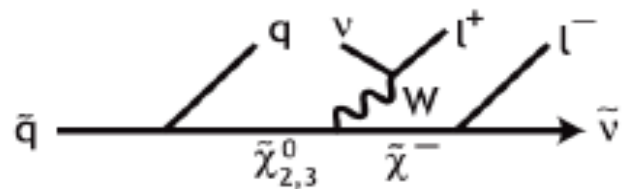
# Do we understand what the new physics is?



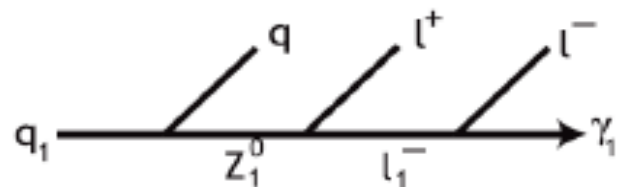
Many ways to produce a 'signature' in  $l^+l^-$  jet +  $\cancel{E}_T$  (dark matter signature)



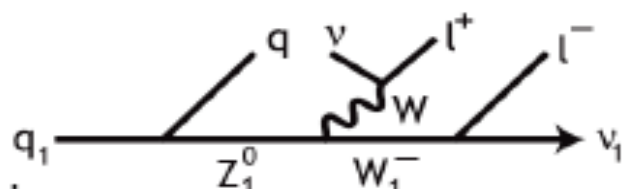
**SUSY LSP = neutralino**



**SUSY LSP = scalar neutrino**

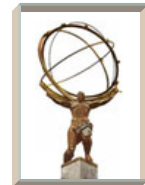


**X – dimension model variants**

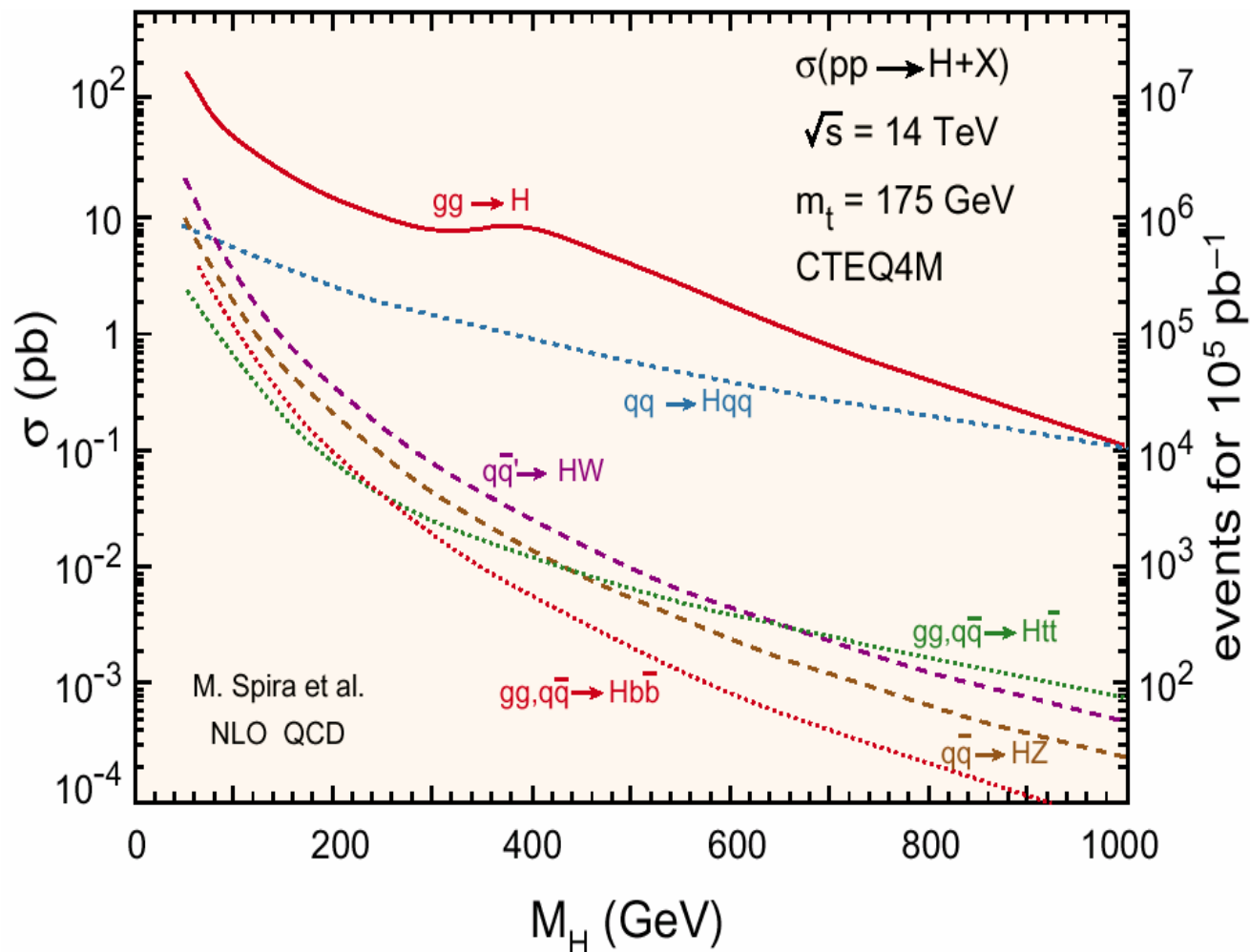
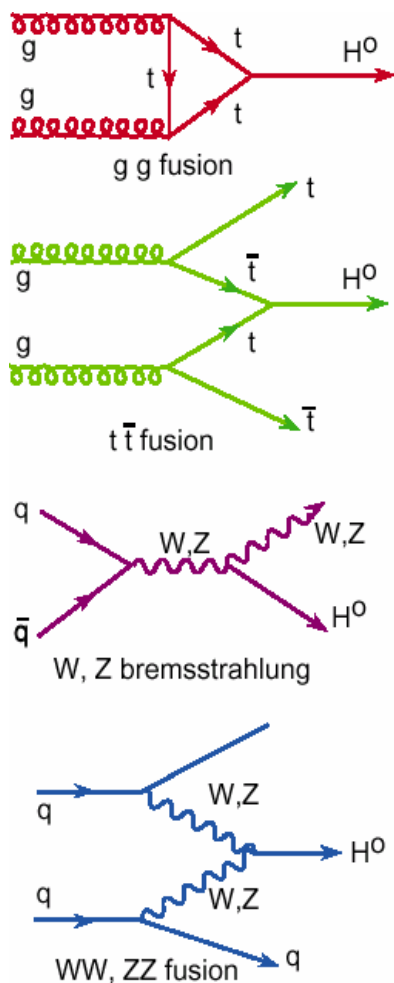


LHC can't distinguish these interpretations. At ILC, the cross-sections and angular distributions for different initial state polarizations tell us which is happening.





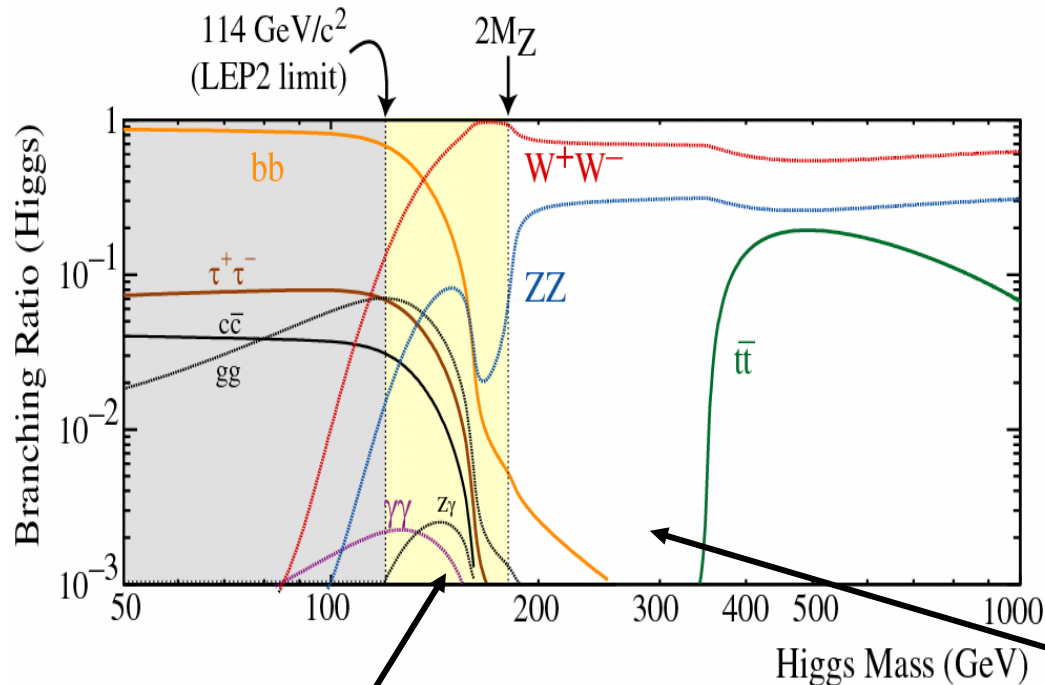
# Higgs Production Mechanism @ LHC



4 production mechanism  $\rightarrow$  key to measure H-boson parameters



# Higgs Discovery Channels at LHC



Dominant BR for  $m_H < 2m_Z$ :

$\sigma(H \rightarrow bb) \approx 20 \text{ pb}$ ;

$\sigma(bb) \approx 500 \mu\text{b}$

for  $m(H) = 120 \text{ GeV}$

→ no hope to trigger  
or extract fully  
hadronic final states

→ look for final  
states with  $\ell, \gamma$   
( $\ell = e, \mu$ )

Low mass region:  $m(H) < 2 m_Z$  :

$H \rightarrow \gamma\gamma$  : small BR, but best resolution

$H \rightarrow bb$  : good BR, poor resolution →  $ttH$ ,  $WH$

$H \rightarrow ZZ^* \rightarrow 4\ell$

$H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$  or  $\ell\nu jj$  : via VBF

$H \rightarrow \tau\tau$  : via VBF

$m(H) > 2 m_Z$  :

$H \rightarrow ZZ \rightarrow 4\ell$

$qqH \rightarrow ZZ \rightarrow \ell\ell \nu\nu^*$

$qqH \rightarrow ZZ \rightarrow \ell\ell jj^*$

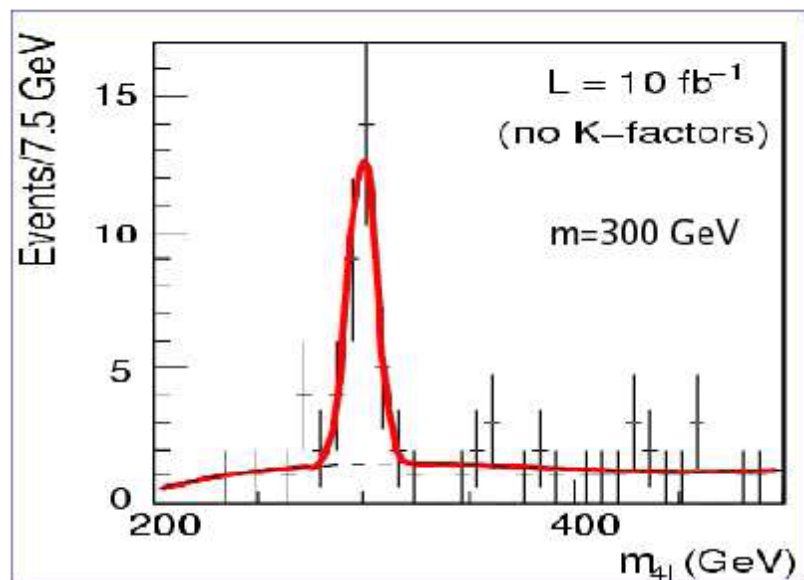
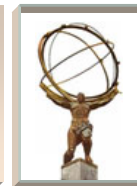
$qqH \rightarrow WW \rightarrow \ell\nu jj^*$

\* for  $m_H > 300 \text{ GeV}$

forward jet tag



# Light SM Higgs Search



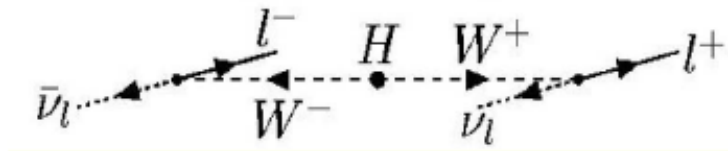
- The  $ZZ \rightarrow 4\ell$  very clean signal.
  - Also use  $ZZ \rightarrow \ell\ell qq, \ell\ell\nu\nu\dots$
- All of the Higgs boson's decay products are reconstructed.
- Very good sensitivity in this channel for  $m_H \geq 130$  GeV.
- For  $m_H > 2m_Z$  this becomes the “gold plated channel”.
- Dominant background from continuum  $ZZ$  production.
  - Can be estimated from sidebands in data.
- Most important backgrounds:
  - $ZZ$  production. Irreducible.
  - $t\bar{t}$  and  $Zbb$ , with two semi-leptonic  $b$ -quark decays.



# H $\rightarrow$ WW

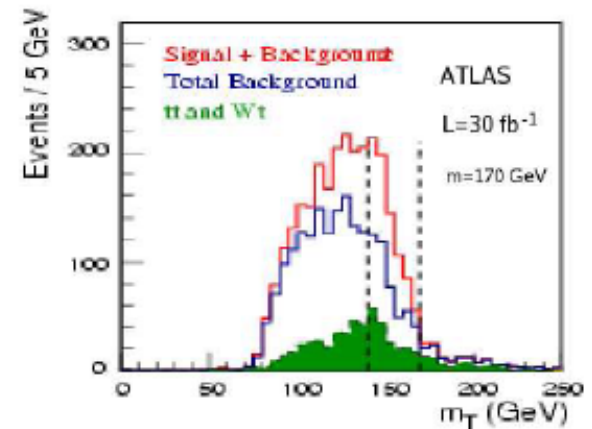
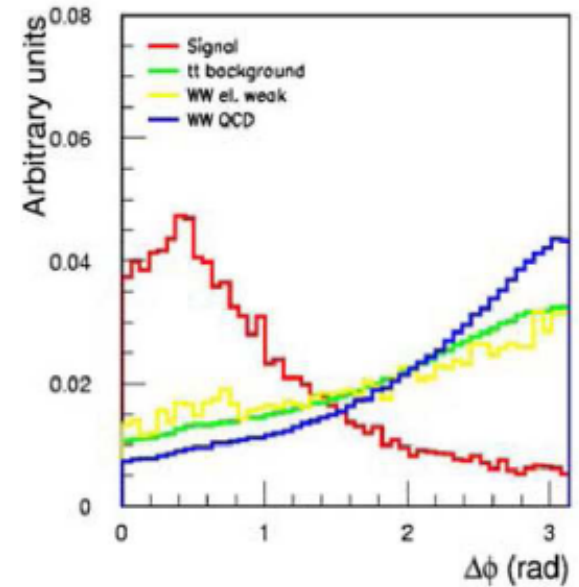


- $BR(H \rightarrow WW)$  is nearly 98% for a Higgs boson with  $m_H \approx 160$  GeV.
- Backgrounds from  $WW$ ,  $t\bar{t}$ ,  $WZ$ .
- Use the lepton spin correlations:



- No mass peak, have to use  $m_T$ :

$$m_T = \sqrt{2p_T^{\ell\ell} E_T (1 - \cos \Delta\phi)}$$

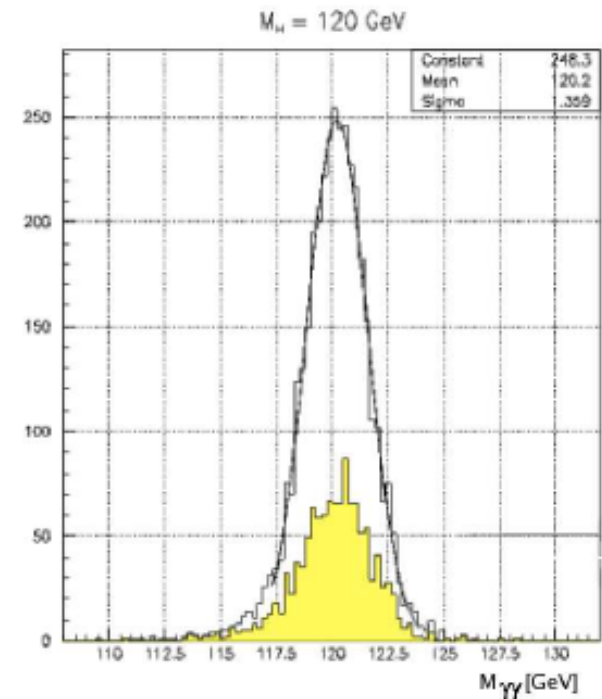




$$H \rightarrow \gamma \gamma$$



- Typical event selection criteria:
  - Two isolated photons with  $p_T > 40$  and  $p_T > 25$ .
  - Both photons within  $|\eta| < 2.4$ , excluding gaps.
  - Invariant mass of the photons:  $M_H - 2 < M_{\gamma\gamma} < M_H + 2$  GeV.
- Mass resolution  $\sigma(M_{\gamma\gamma}) \approx 1.36$  GeV.
- Photons conversions taken into account ( $\sim 40\%$  of the events).
- EM calorimeter is crucial:
  - Energy and angle resolution.
  - Photon acceptance.
  - $\gamma/\text{jet}$  and  $\gamma/\pi^0$  rejection.

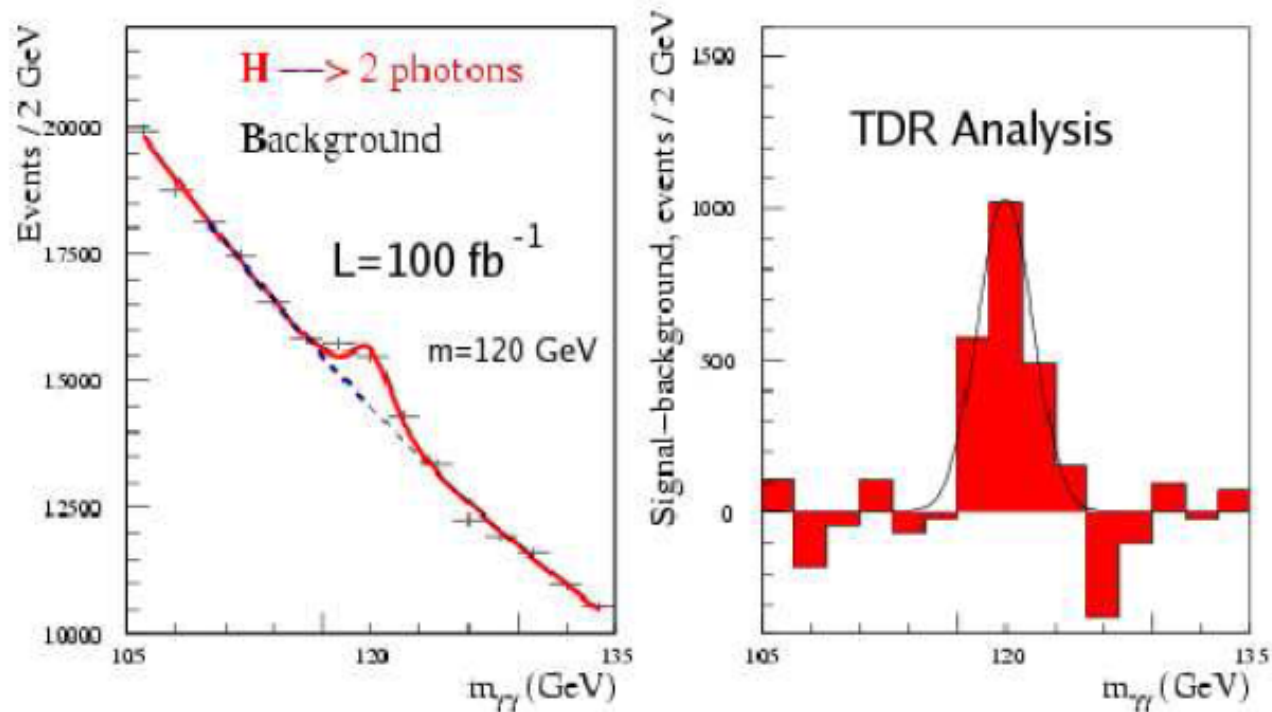




# H $\rightarrow$ $\gamma\gamma$ Background



- The largest background comes from irreducible  $\gamma\gamma$  production.  $\sigma_{\gamma\gamma} \approx 125 \text{ fb/GeV}$ . Need mass resolution  $\sigma(M_{\gamma\gamma})/M_{\gamma\gamma} \approx 1\%$ .
- Reducible background from jets misreconstructed as photons.  $\sigma_{\gamma j} \approx 8 \cdot 10^2 \cdot \sigma_{\gamma\gamma}$  and  $\sigma_{jj} \approx 2 \cdot 10^6 \cdot \sigma_{\gamma\gamma}$ . Need rejection  $R_j > 10^3$ .



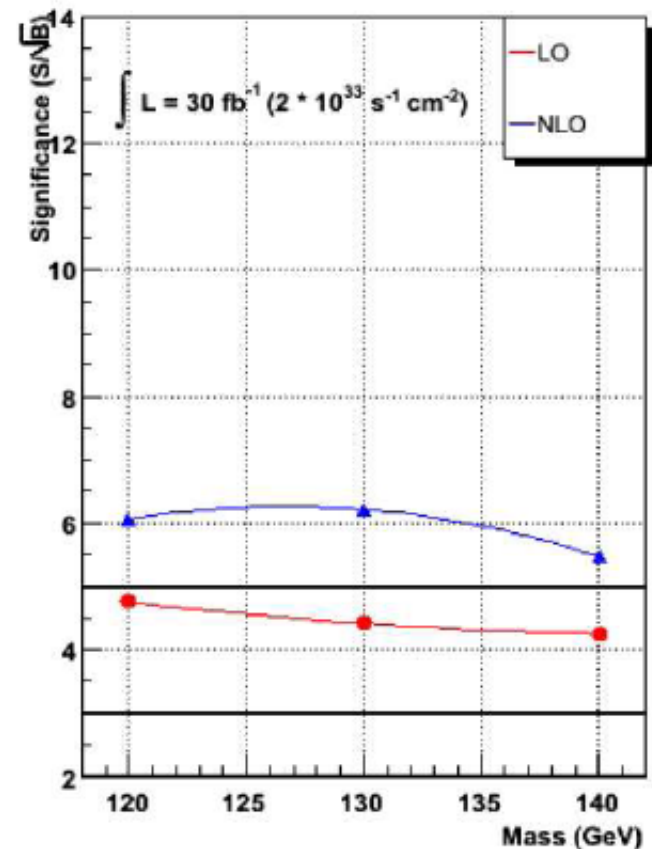


# Significance with NLO Calculation



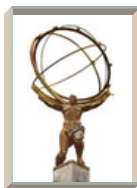
- Cross sections calculated at NLO:
  - $K \sim 1.8$  for  $gg \rightarrow H$ .
  - The full  $\mathcal{O}(\alpha_S)$  calculation, plus box diagram, for the  $\gamma\gamma$  background.
  - $K \sim 1.7$  for  $\gamma j$  and  $j j$  backgrounds.

Event yield (at NLO) for $30 \text{ fb}^{-1}$		
$m_H$	120 GeV	130 GeV
$H \rightarrow \gamma\gamma$	815	758
$\gamma\gamma$	14100	9552
$\gamma j, j j$	3967	3396
$S/\sqrt{B}$	6.06	6.22

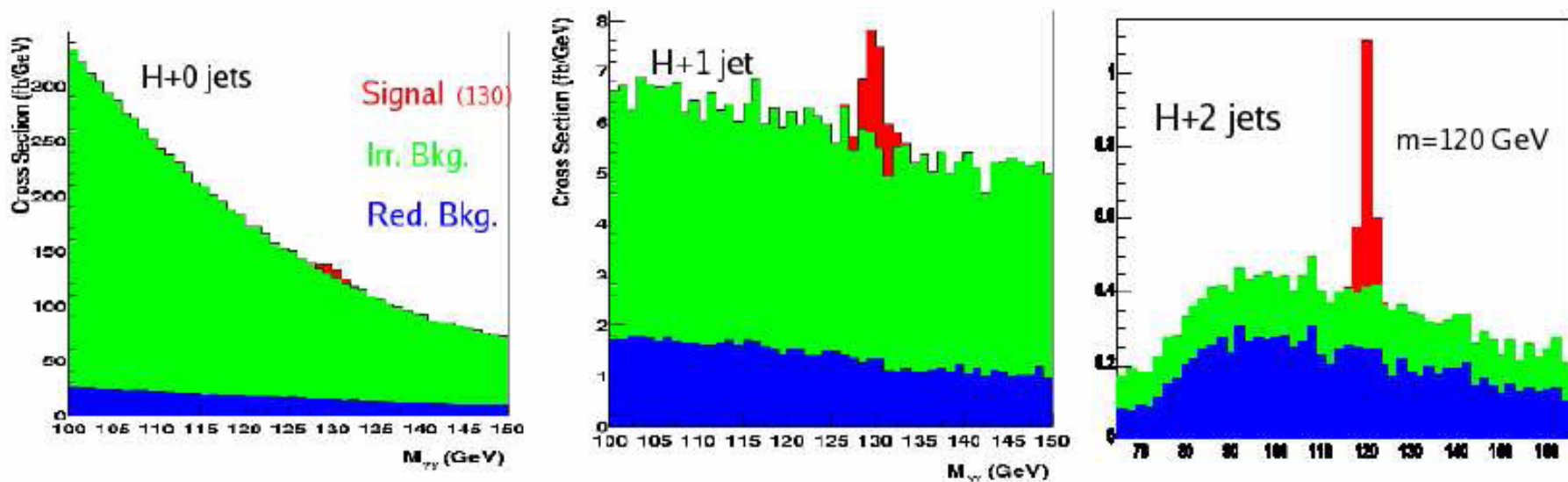


- Gives dramatic improvement in sensitivity.





# Combination of H + n-jets Analysis



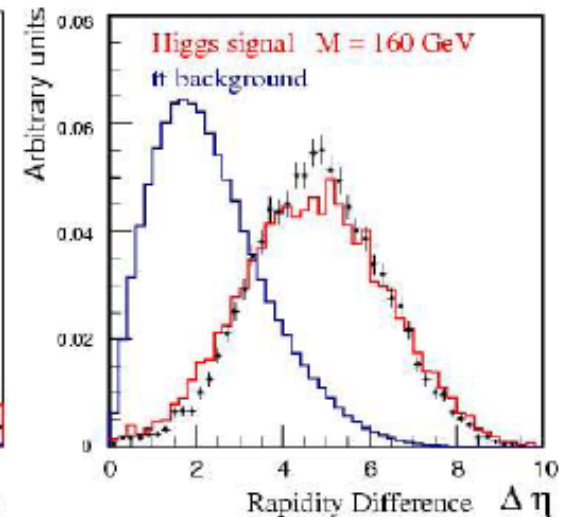
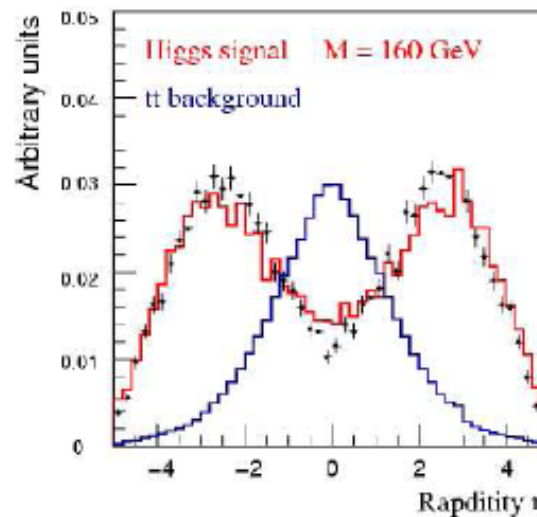
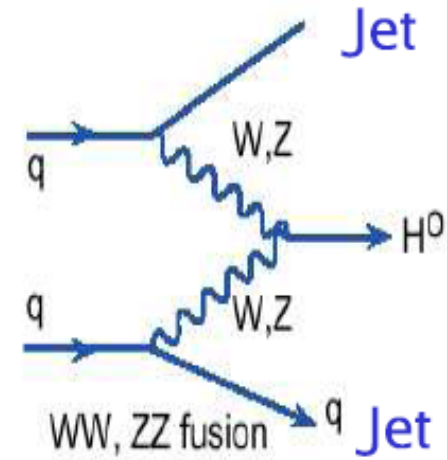
- Do inclusive analysis. Use a combination of all jet multiplicities.
  - $H + 0$  jets from  $gg \rightarrow H$ .
  - $H + 1$  jet at NLO, plus VBF production with one lost jet.
  - $H + 2$  jets from VBF.
- Decrease in signal, increase in S/B with jet multiplicity.
  - Decrease in systematic uncertainty?
- This is a preliminary study.
  - Looks promising.



# Signal with VBF Events



- Typical selections for tag jets:
  - $p_T > 40$  and  $p_T > 20$  GeV.
  - $\Delta\eta > 3.8$  between the tag jets.
- Apply a central jet veto:
  - No central jets with  $p_T > 20$  GeV.
  - $H$  decay products between tag jets.
- Powerful way to reduce backgrounds.
- Uncertainty from underlying and overlapping event, pile-up.

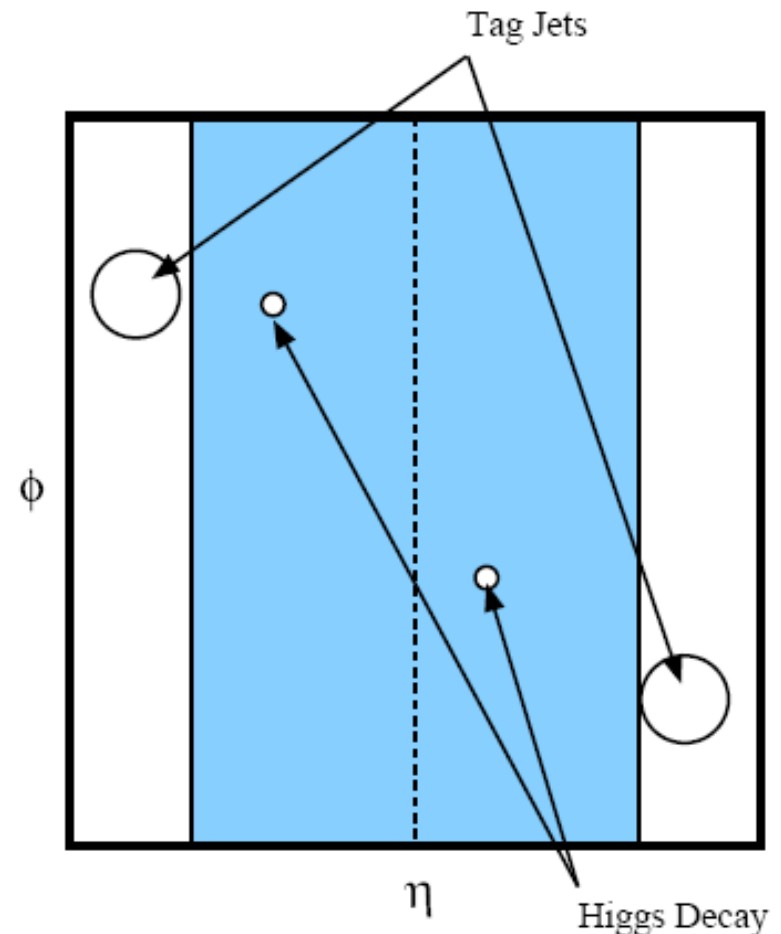




# $H \rightarrow \tau\tau$ and $H \rightarrow WW$ in VBF production

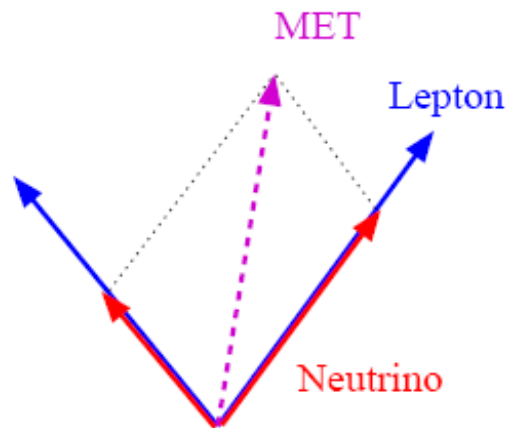


- At low Higgs masses the largest sensitivity search channels are found in the vector boson fusion production mode.
- At least one of the  $W/\tau$ 's have to decay leptonically.
- Main backgrounds are  $t\bar{t}$ ,  $Wt$ ,  $WW$ +jets,  $\gamma^*/Z$ +jets.
- Some selection criteria ( $e\mu$ ):
  - $p_T^e > 15 \text{ GeV}$ ,  $p_T^\mu > 10 \text{ GeV}$ .
  - $|\eta_e| < 2.5$  and  $M_{e\ell} < M_H/2$ .
  - Tag jet cuts, central jet veto.
- $\tau$  reconstruction provide extra sensitivity or rejection.

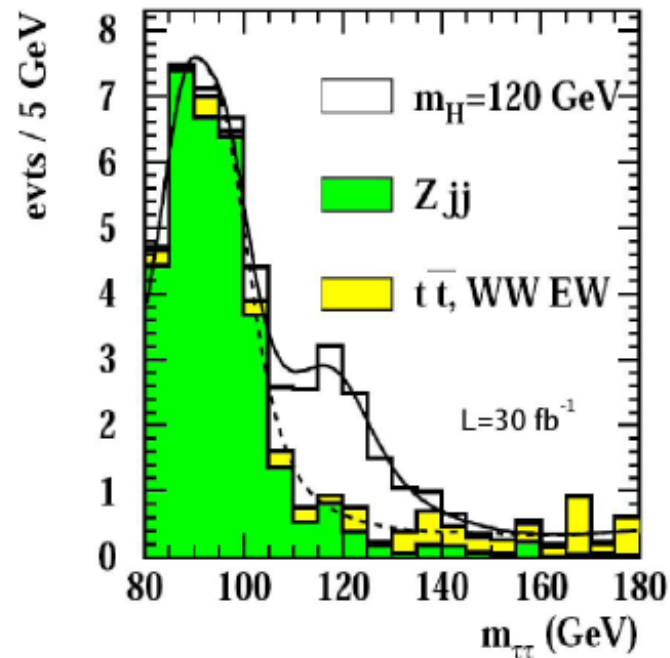
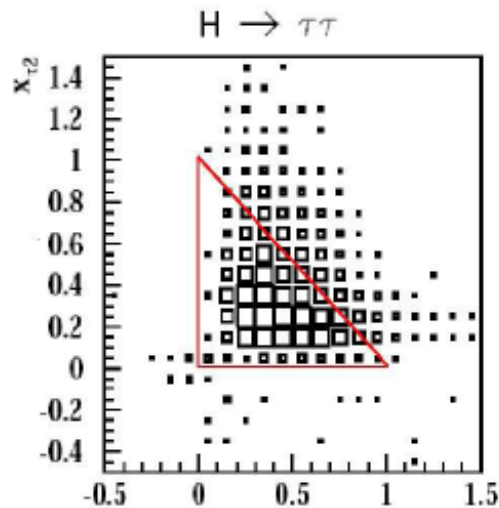




# VBF $H \rightarrow \tau\tau$ Reconstruction



- Assume  $\ell$  and  $\nu$ 's from  $\tau \rightarrow \ell\nu\nu$  are collinear.
- Label visible energy fractions  $x_{\tau_1}$  and  $x_{\tau_2}$ .
- Assume  $\vec{p}_T$  vector comes from the  $\nu$ 's, and solve the equations for  $x_{\tau_1}$  and  $x_{\tau_2}$ .

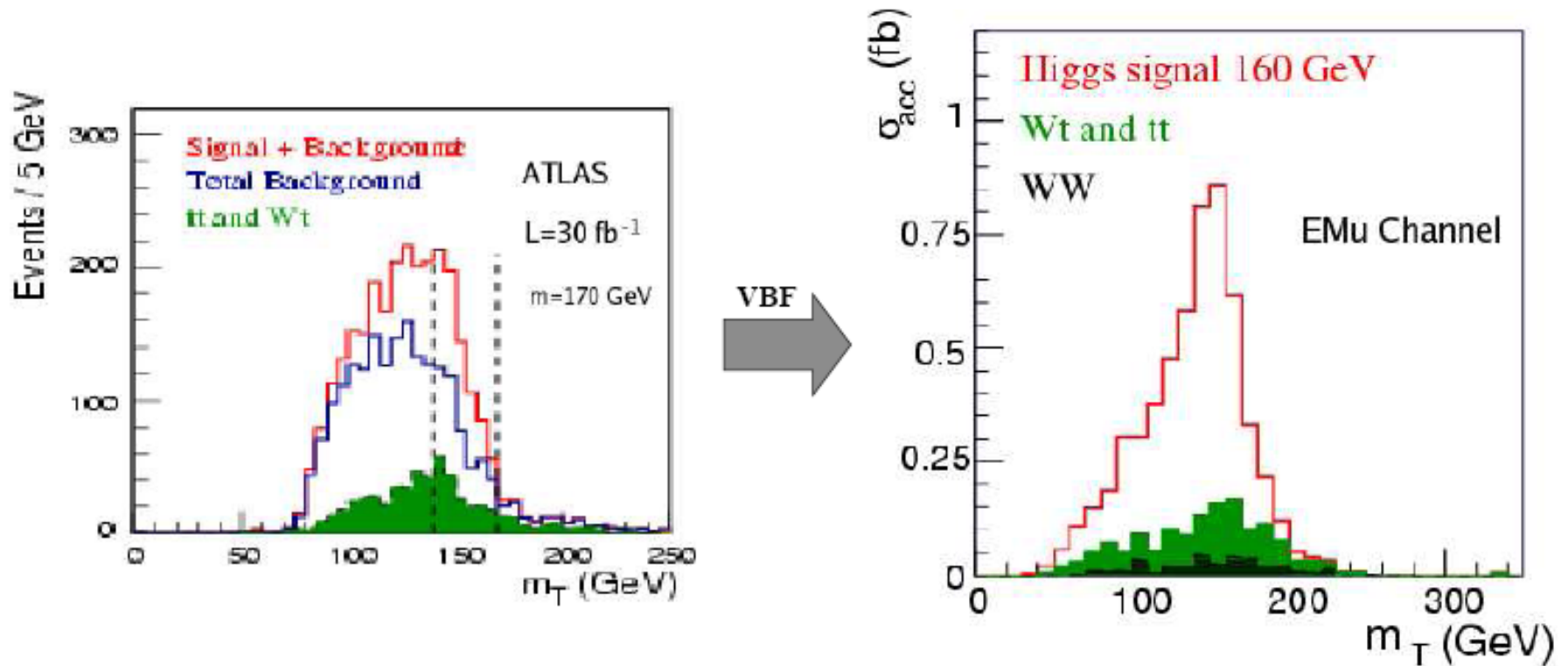




# VBF $H \rightarrow WW$ Reconstruction



- As in the inclusive  $WW$  analysis, spin correlations between the leptons are used to enhance the signal fraction.



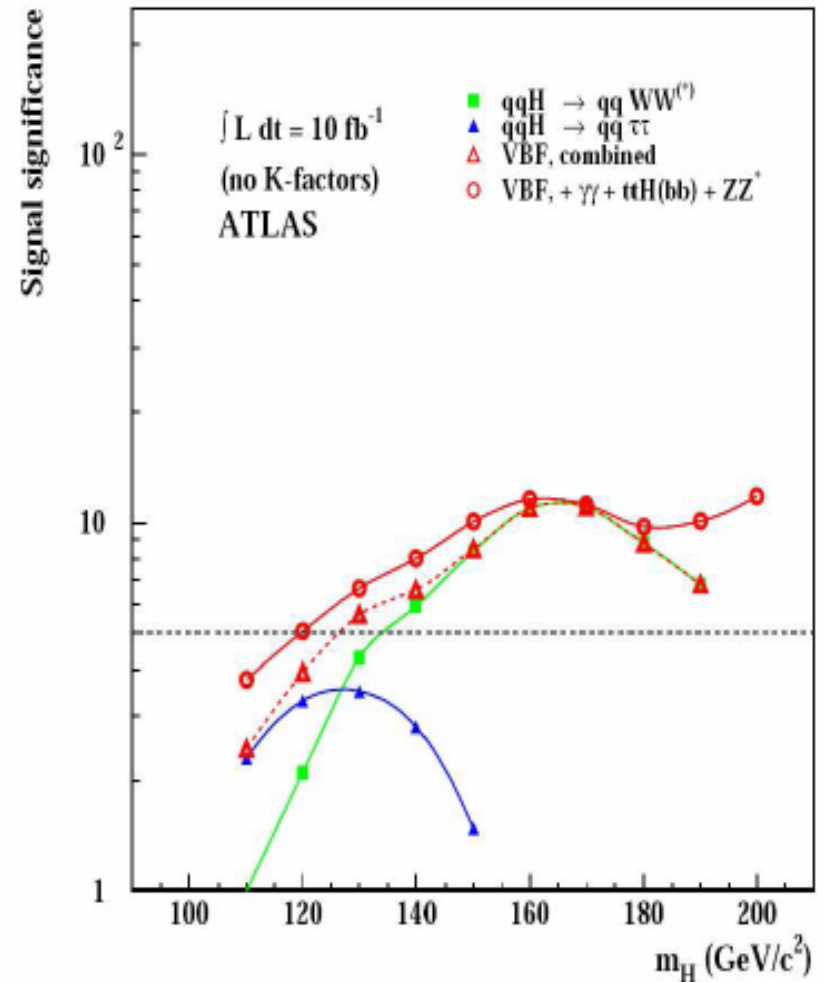
- Signal to background in the VBF channel increases to  $\sim 3.6$ .



# Sensitivity of VBF Channels

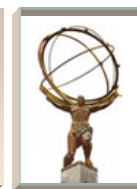


- Both  $H \rightarrow \tau\tau$  and  $H \rightarrow WW$  are sensitive for low Higgs masses.
- $H \rightarrow \tau\tau$ :
  - $m_H \lesssim 145 \text{ GeV}$ .
- $H \rightarrow WW$ :
  - $125 \leq m_H \leq 200 \text{ GeV}$ .
- If the two channels are combined, a  $5\sigma$  discovery is possible for  $m_H \gtrsim 130 \text{ GeV}$  with  $10 \text{ fb}^{-1}$ .





# Most Challenge Channel $H + t\bar{t}$



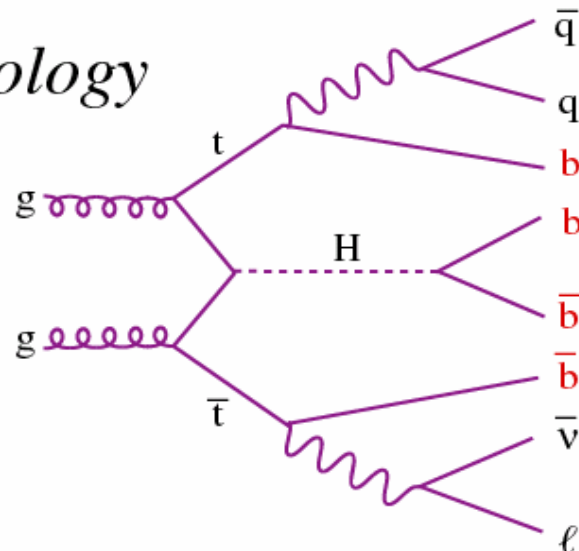
*Challenging and complex topology*

4 b-jets, 2 jets, 1 lepton

$H \rightarrow b\bar{b}$

$t \rightarrow bq\bar{q}'$

$t \rightarrow b\ell\nu$



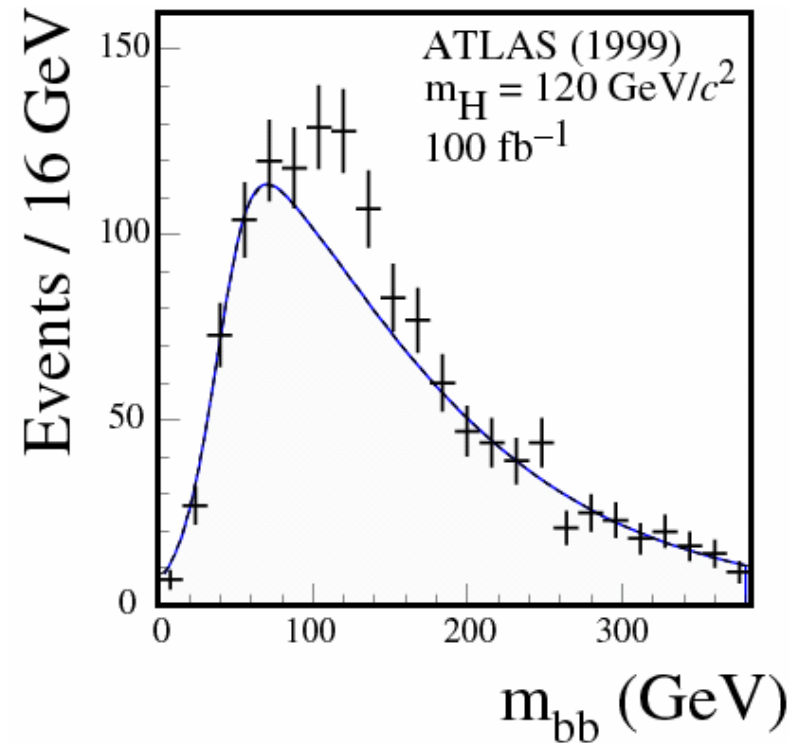
- Spectacular, very energetic, final state!
- The  $\ell$  from the  $t$  decay allows for triggering, even though  $H \rightarrow b\bar{b}$ .
- Dominant background is from non-resonant  $t\bar{t}b\bar{b}$  production.
  - Smaller backgrounds from  $t\bar{t}Z$ ,  $t\bar{t}jj$ ,  $WWb\bar{b}jj$  ...



# Need a lot of Luminosity for H ttbar detection



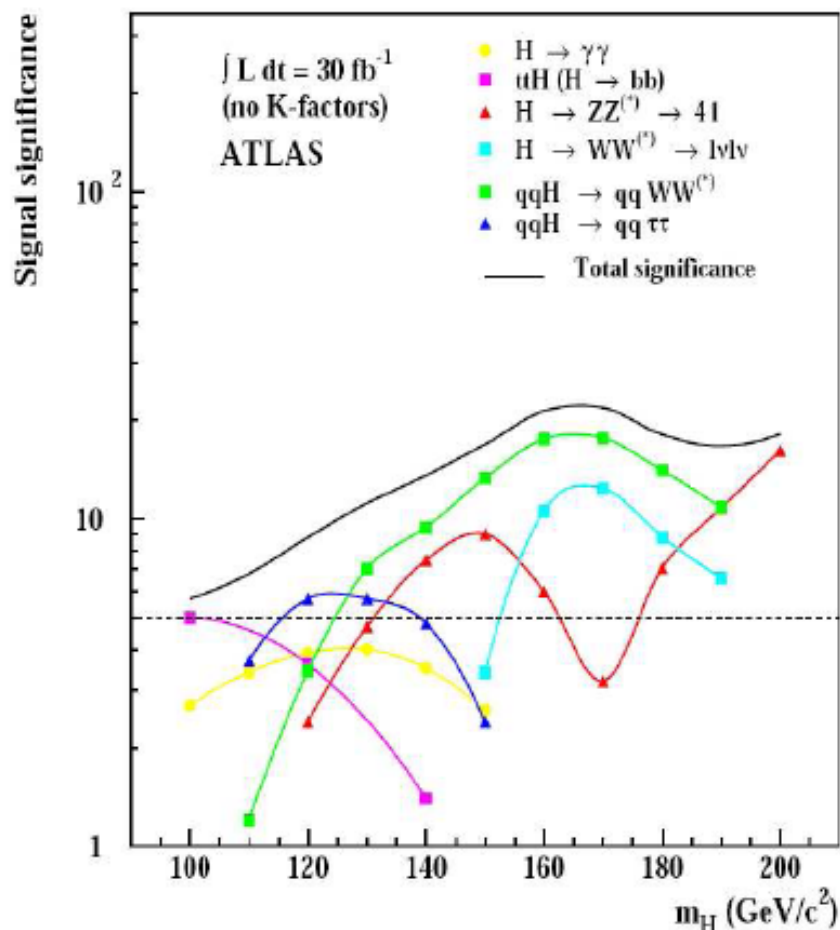
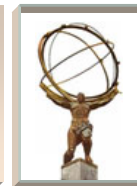
- To extract the signal:
  - Reconstruct all six jets.
  - Exactly 4  $b$ -tagged jets.
  - Reconstruct the two  $t$  quarks.
  - Determine invariant mass of  $b$ -jets from the Higgs decay.



Signal significance ( $5\sigma$ ) :  
 $m_H < 120 \text{ GeV}$  needs  $100 \text{ fb}^{-1}$



# Combined Sensitivity



- Only LO results shown in plot.
- Full mass range covered after a few years of running.
- Several channels available for any given Higgs mass.
- VBF channels play important role for low Higgs masses.
- For  $m_H \leq 120 \text{ GeV}$  three complementary channels:
  - $H \rightarrow \gamma\gamma$  ( $M_{\gamma\gamma}$  resolution).
  - $t\bar{t}H$  ( $b$ -tagging).
  - $qqH \rightarrow qq\tau\tau$  (large  $|\eta|$  jets).



# Search for MSSM Higgs(es)



Complex analyses; 5 Higgses:  $h_0$ ,  $H_0$ ,  $A_0$ ,  $H^\pm$

- *At tree level, all masses and couplings depend on:  $M_A$  and  $\tan\delta$*
- **Large variety of observation modes**
  - if SUSY particles heavy
    - **SM-like:**  $h \rightarrow \gamma\gamma, bb; H \rightarrow 4\ell$
    - **MSSM-specific:**  $A/H \rightarrow \mu\mu, \tau\tau, tt; H \rightarrow hh, A \rightarrow Zh; H^\pm \rightarrow \tau\nu$
  - if SUSY particles accessible:
    - $H/A \rightarrow \chi^2_0 \chi^2_0 \rightarrow 4\ell + \text{missing Energy}$
    - $h$  produced in cascade decays (e.g.  $\chi^2_0 \rightarrow h \chi^1_0$ )
- **Studies performed in two steps:**
  - i. SUSY particles are heavy: no contribution to Higgs production/decay
  - ii. SUSY particles contribute in production/decays
    - Impact on Higgs decay to SM particles generally small
      - $h \rightarrow \gamma\gamma$  10% smaller
      - $A/H \rightarrow \text{SM}$  at most 40% smaller

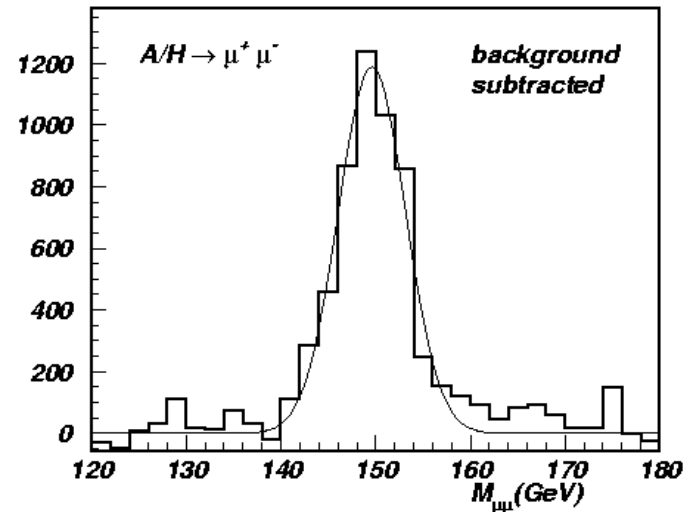
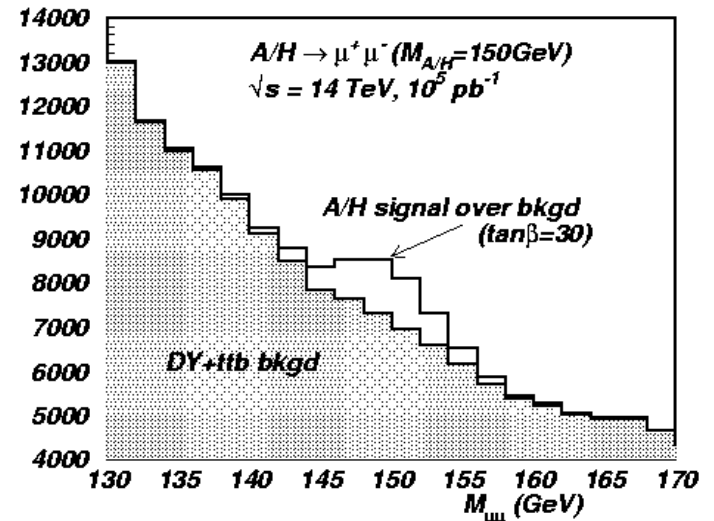
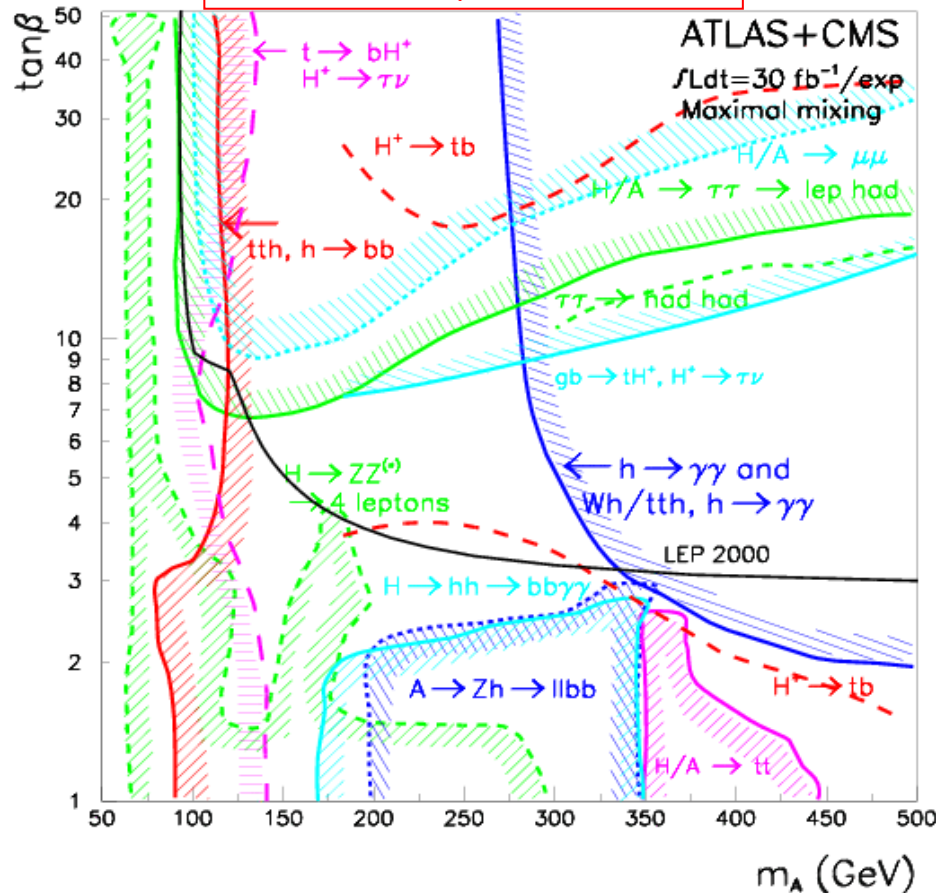


# MSSM Higgs Discovery Potential



- ✓ Plane fully **covered** with **30 fb<sup>-1</sup>**
- ✓ 2 or more Higgses observable in large fraction of plane

**5 $\sigma$  discovery contours**





# Higgs Boson Parameters: Mass

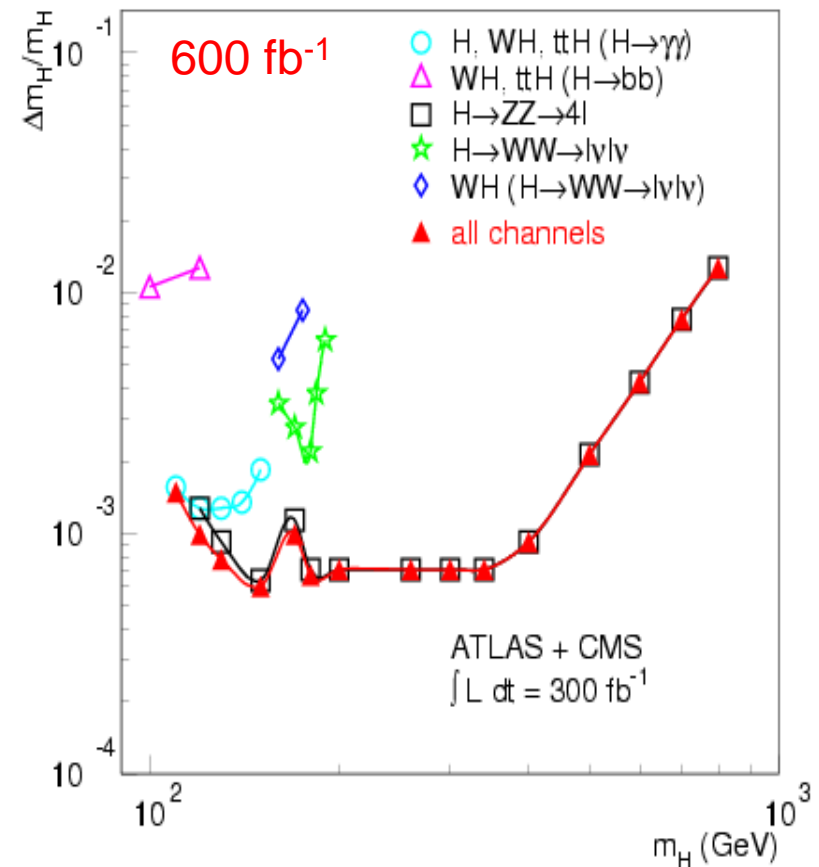


## SM

- ✓ Limited by absolute energy scale
  - 0.1% for  $\ell/\gamma$  ( $Z \rightarrow \ell\ell$  calibration)
  - 1% for jets
- ✓ Resolutions:
  - For  $\gamma\gamma$  &  $4\ell \approx 1.5 \text{ GeV}/c^2$
  - For  $bb \approx 15 \text{ GeV}/c^2$
- ✓ At large masses: decreasing precision due to large  $\Gamma_H$

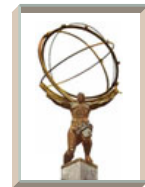
## MSSM

- ✓  $h$  as in SM case
- ✓  $H/A$ : 0.1 - 0.5% in modes  $\gamma\gamma$ ,  $4\ell$ ,  $\mu\mu$
- ✓ 1 - 2% in modes  $bb$ ,  $bb\gamma\gamma$  ( $hh$ ),  $bb\ell\ell$  ( $Zh$ )

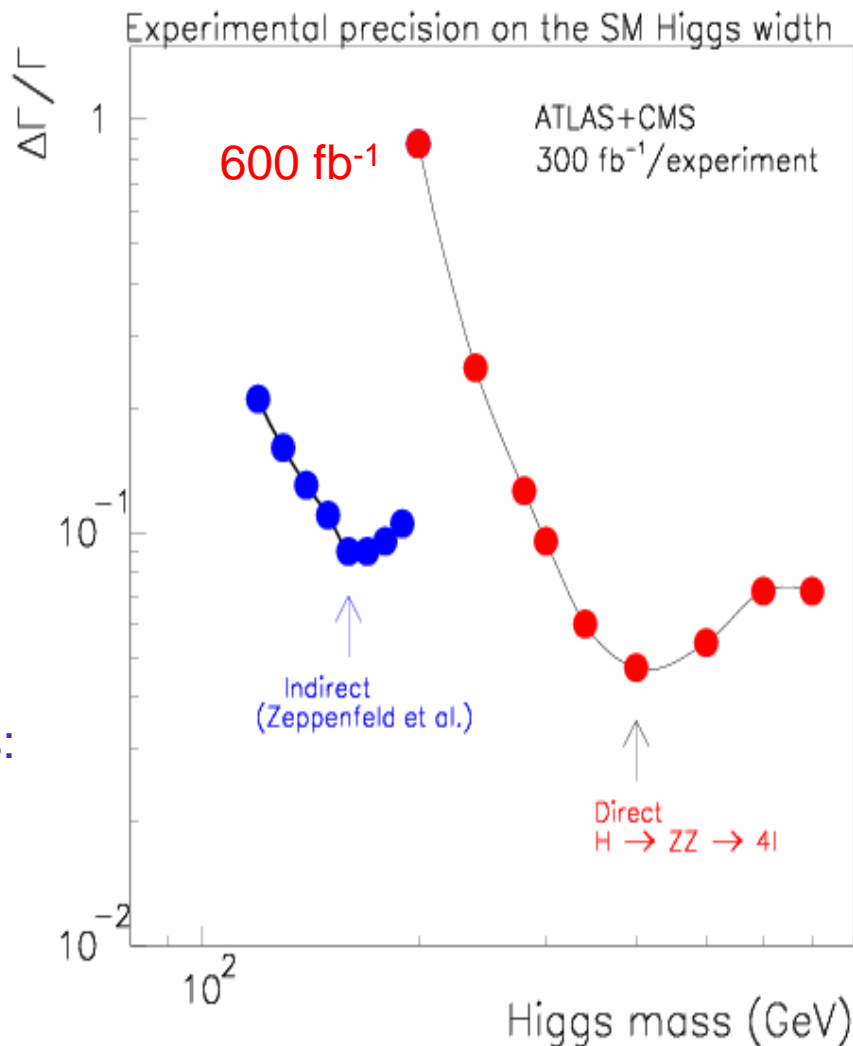




# Higgs Parameters: Width



- Direct width measurement:
  - Mass peak width of  $H \rightarrow ZZ \rightarrow 4\ell$  channel  
for  $M_H > 200$  GeV  
( $\Gamma_H > \Gamma_{\text{exp}}$  in SM)
  - Systematics: radiative decays (1.5%)
- Indirect width measurement:
  - Comparison of rates in channels:
    - $qq \rightarrow qqH, H \rightarrow \gamma\gamma, \tau\tau, WW$
    - $gg \rightarrow H, H \rightarrow \gamma\gamma, ZZ^*, WW^*$
  - Assumption:
    - o  $\text{BR}(H \rightarrow cc, \text{non-standard}) < 10\%$





# Higgs Parameters: Couplings (1)



## Ratio of boson-boson couplings

Direct: ratio between W and Z partial width

$$\frac{\sigma \times BR(H \rightarrow WW^*)}{\sigma \times BR(H \rightarrow ZZ^*)} = \frac{\Gamma_g \Gamma_W}{\Gamma_g \Gamma_Z} = \frac{\Gamma_W}{\Gamma_Z}$$

QCD correction cancel

VBF: ratio between W and Z partial width

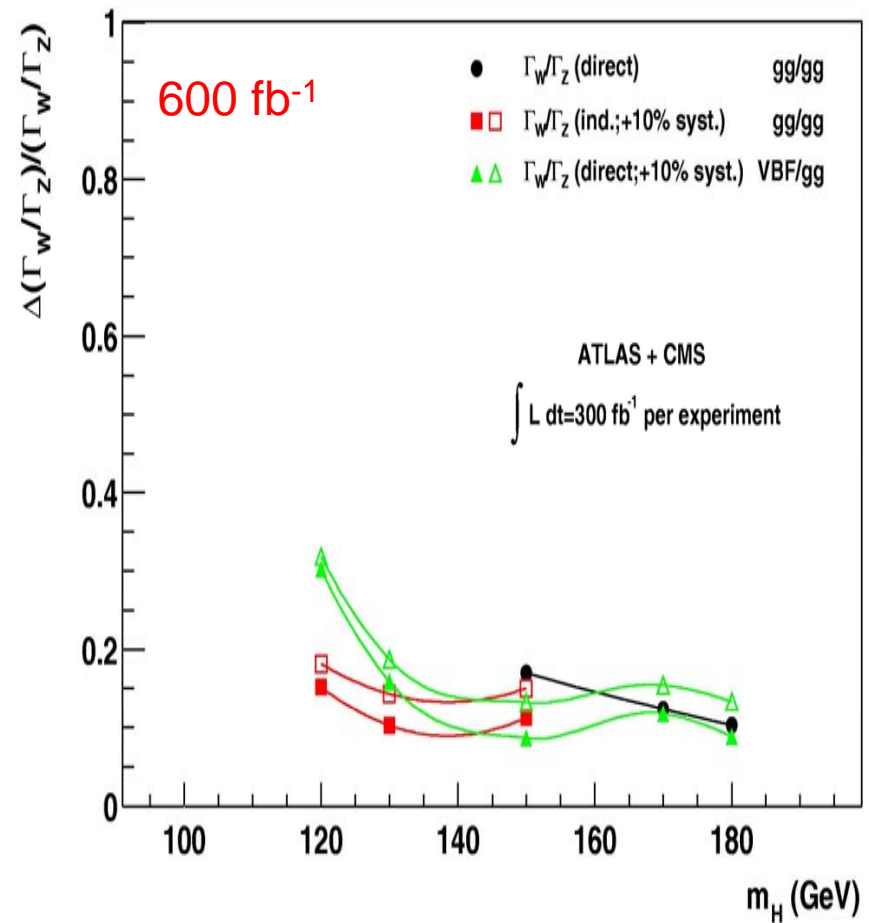
$$\frac{\sigma \times BR(qqH \rightarrow qqWW^*)}{\sigma \times BR(H \rightarrow ZZ^*)} = \frac{\Gamma_W}{\Gamma_Z}$$

different processes, QCD corrections do not cancel, i.e. add. uncertainty

Indirect: ratio between  $\gamma$  and Z partial width

$$\frac{\sigma \times BR(H \rightarrow \gamma\gamma)}{\sigma \times BR(H \rightarrow ZZ^*)} = \frac{\Gamma_g \Gamma_\gamma}{\Gamma_g \Gamma_Z} \approx \frac{\Gamma_W}{\Gamma_Z}$$

Use proportionality between  $G_W$  and  $G_\gamma$ , Needs theoretical input, 10% uncertainty assumed





# Higgs Parameters: Couplings (2)



## Ratio of boson-fermion couplings

VBF: allows a direct measurement of  $G_W/G_t$  in the mass range 120 - 150 GeV

**Direct :**

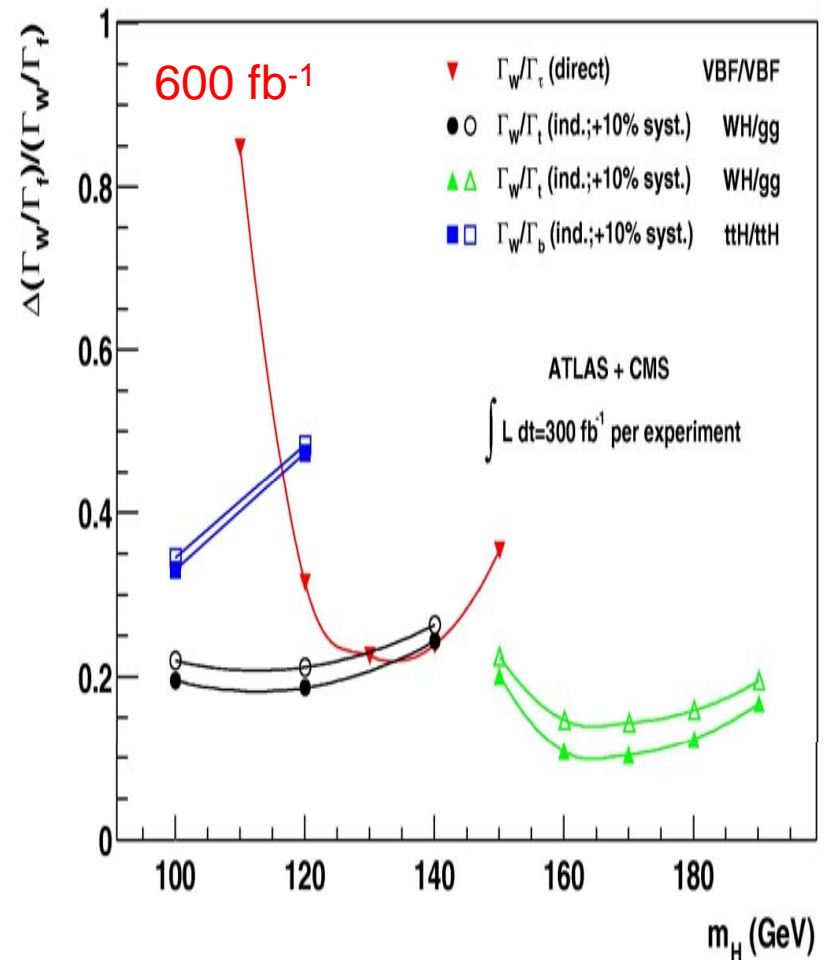
$$\frac{\sigma \times BR(qqH \rightarrow qqWW)}{\sigma \times BR(qqH \rightarrow qq\tau\tau)} = \frac{\Gamma_W \Gamma_W}{\Gamma_W \Gamma_\tau} = \frac{\Gamma_W}{\Gamma_\tau}$$

**Indirect :**

$$\frac{\sigma \times BR(WH \rightarrow W\gamma\gamma)}{\sigma \times BR(H \rightarrow \gamma\gamma)} \sim \frac{\Gamma_W}{\Gamma_t} \cdot C_{QCD}$$

$$\frac{\sigma \times BR(WH \rightarrow WWW)}{\sigma \times BR(H \rightarrow WW)} \sim \frac{\Gamma_W}{\Gamma_t} \cdot C_{QCD}$$

$$\frac{\sigma \times BR(ttH \rightarrow ttbb)}{\sigma \times BR(ttH \rightarrow tt\gamma\gamma)} \sim \frac{\Gamma_b}{\Gamma_W}$$



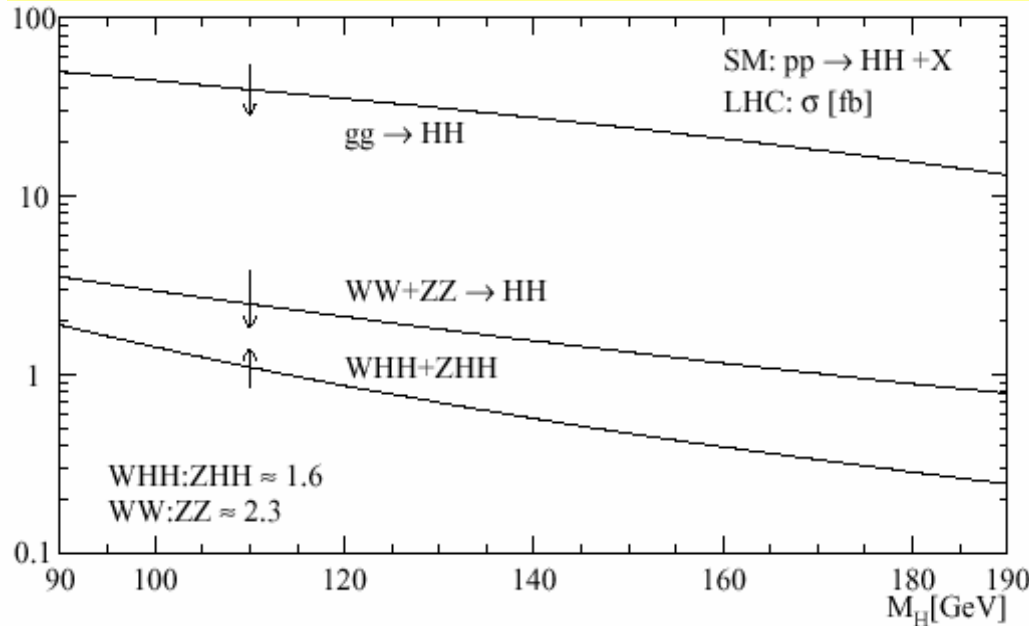


# Higgs Parameters: Self-couplings



## Cross section of the Higgs pair production at LHC

Hep-ph/0204087



Need **Super LHC** to measure Higgs self-interaction couplings:

$$\lambda_{HHH}^{SM} = 3 \frac{m_H^2}{v}, \quad \lambda_{HHHH}^{SM} = 3 \frac{m_H^2}{v^2}$$

Table 8: Expected numbers of signal and background events after all cuts for the  $gg \rightarrow HH \rightarrow 4W \rightarrow \ell^+ \ell'^+ 4j$  final state, for  $\int \mathcal{L} = 6000 \text{ fb}^{-1}$ .

$m_H$	Signal	$t\bar{t}$	$W^\pm Z$	$W^\pm W^+ W^-$	$t\bar{t} W^\pm$	$t\bar{t} t\bar{t}$	$S/\sqrt{B}$
170 GeV	350	90	60	2400	1600	30	5.4
200 GeV	220	90	60	1500	1600	30	3.8

$$gg \rightarrow HH \rightarrow W^+ W^- W^+ W^- \rightarrow \ell^\pm \nu jj \ell^\pm \nu jj$$



# Strongly-Coupled Vector Boson System



**No light Higgs boson?** Study Longitudinal gauge boson scattering in high energy regime (the L-component which provides mass to these bosons).

$$W_L Z_L \longrightarrow W_L Z_L \longrightarrow \ell \nu \ell$$

(Hep-ph/0204087)

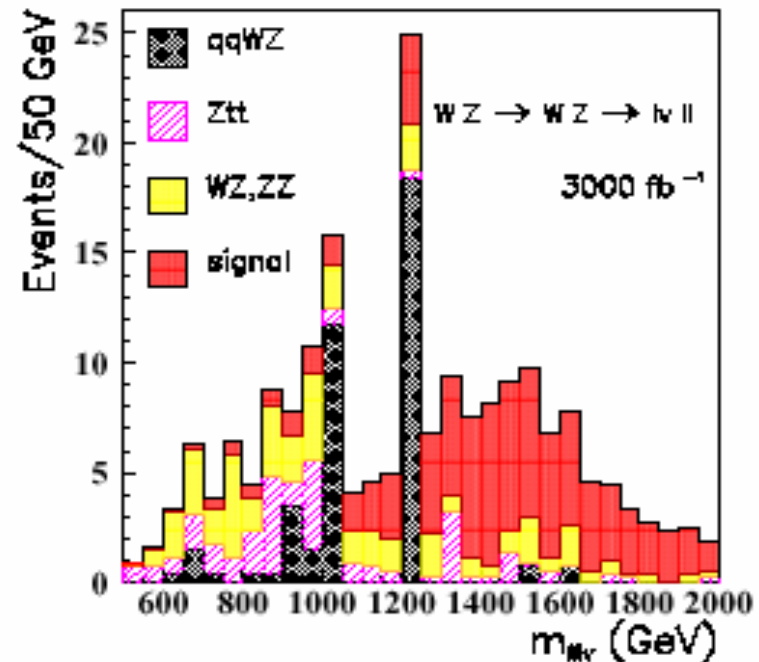
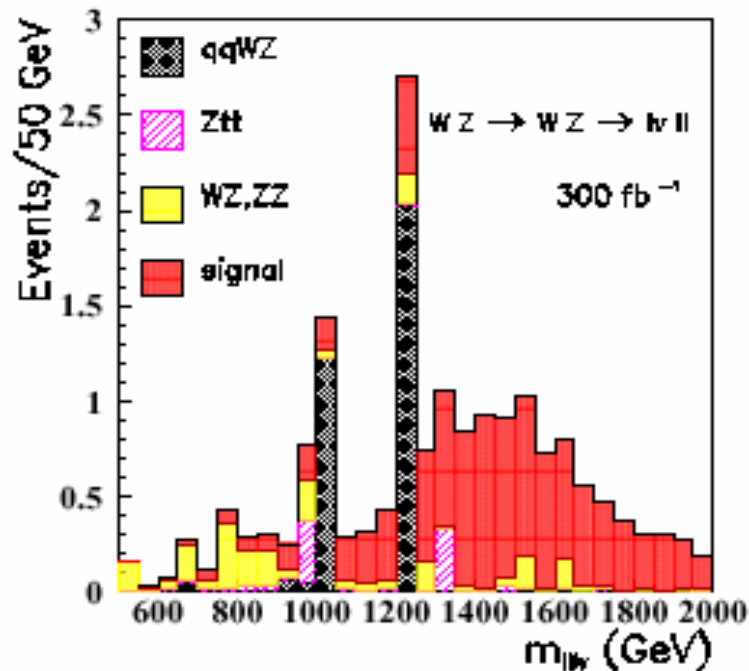
$$p_T(\ell_1) > 150 \text{ GeV}, \quad p_T(\ell_2) > 100 \text{ GeV}, \quad p_T(\ell_3) > 50 \text{ GeV}$$

$$|m(\ell_1 \ell_2) - m_Z| < 10 \text{ GeV}$$

$$E_T^{\text{miss}} > 75 \text{ GeV}$$

$$S / B = 6.6/2.2$$

$$S / \sqrt{B} \sim 10$$



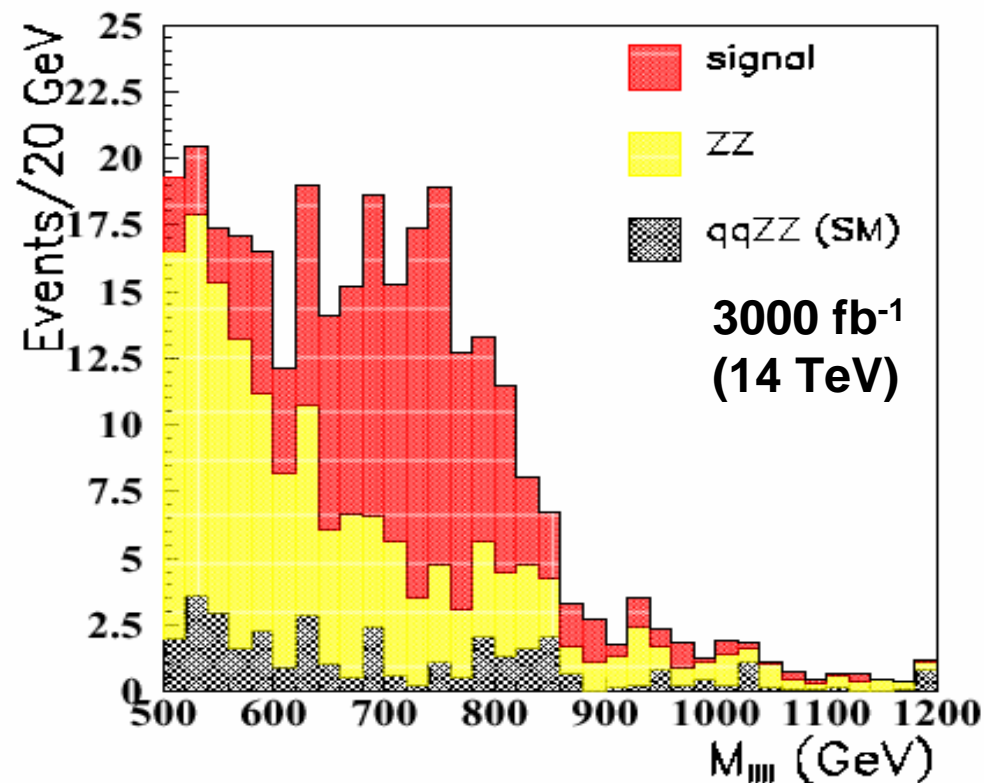


# $Z_L Z_L$ Scalar Resonance



$W_L W_L \rightarrow Z_L Z_L \rightarrow 4 \text{ leptons}$   
 $Z_L Z_L \rightarrow Z_L Z_L \rightarrow 4 \text{ leptons}$

(Hep-ph/0204087)



→ **Search for new resonances**  
(which could regularize vector boson scattering Xsection)



# Extra Dimensions



## Large Extra Dimensions (ADD)

- Gravity in bulk / flat space
- Missing energy / interference / black holes



## Warped Extra Dimensions (RS)

- Gravity in bulk / curved space
- Spin 2 resonances in  $\gg$ TeV range / black holes

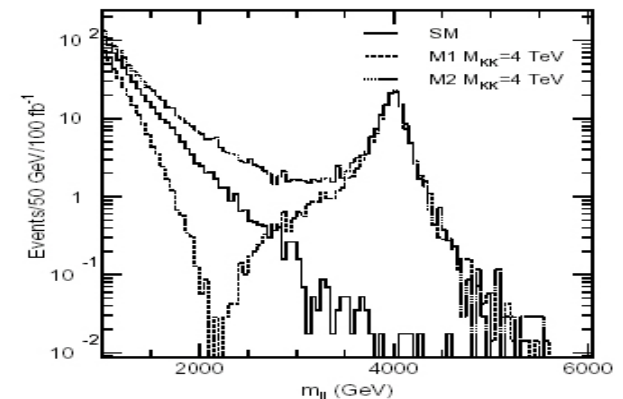
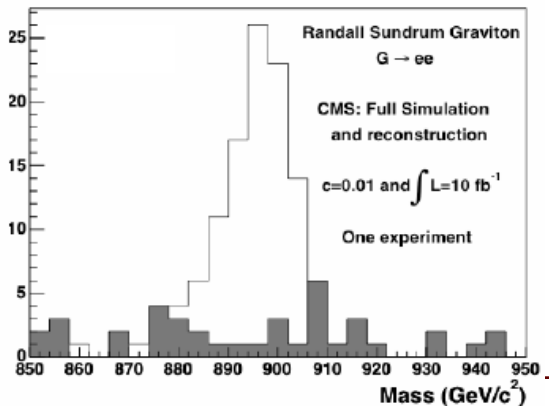
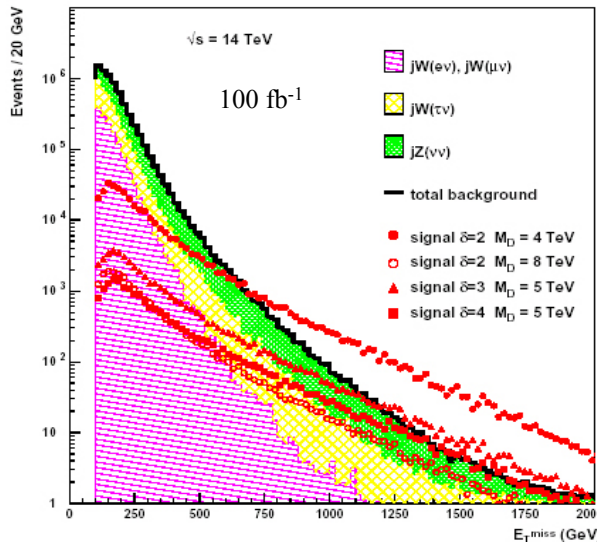
## TeV Scale Extra Dimensions

- Gauge bosons / Higgs in bulk
- Spin 1 resonances in  $\gg$ TeV range
- Interference with Drell-Yan

## Universal Extra Dimensions

- Everybody in the bulk!
- Fake SUSY spectrum of KK states

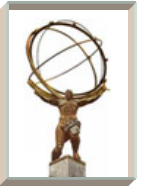
$E_T(\text{jet}) > 1 \text{ TeV}$





# Extra Dimensions:

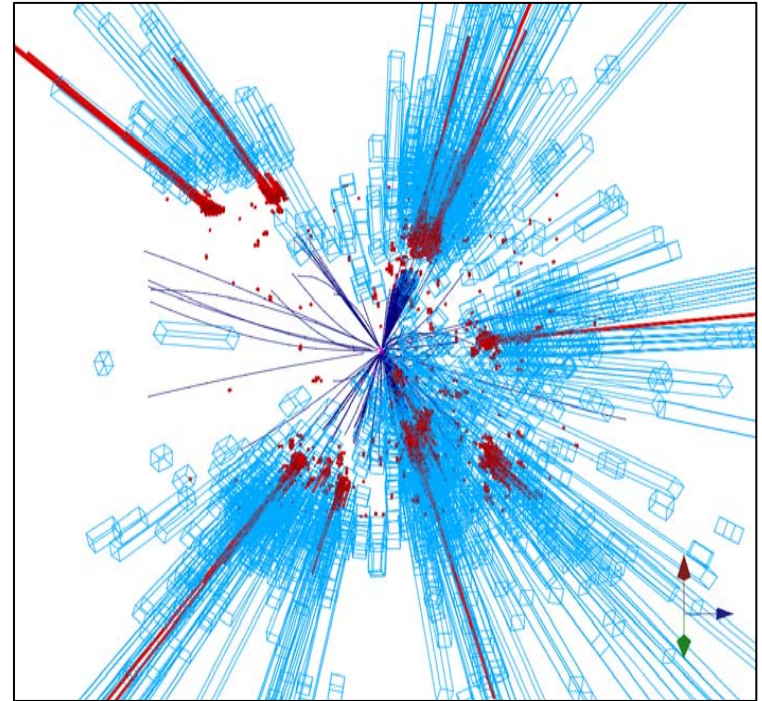
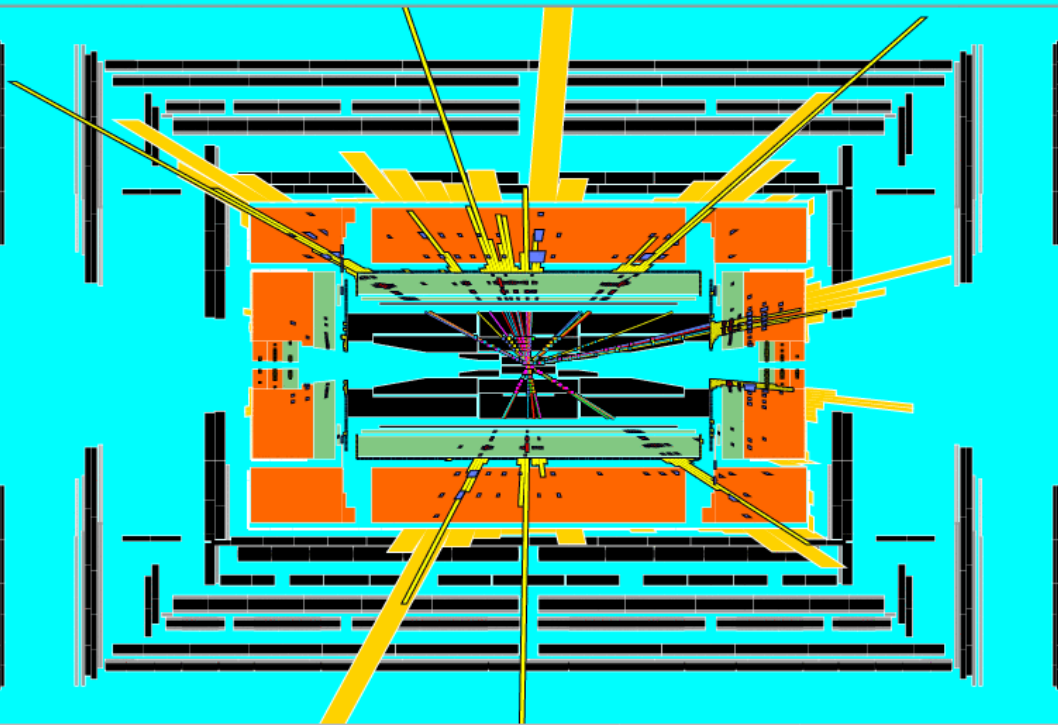
What if Planck Scale  
in TeV Range?



Simulation of a mini black hole  
event with  $M_{BH} \sim 8 \text{ TeV}$  in ATLAS

... and in CMS

ATLAS Atlantis





# Narrow Graviton Resonance: Spin of G



B. C. Allanach, K. Odagiri, A. Parker and B. Webber, *JHEP* 9 (2000) 19

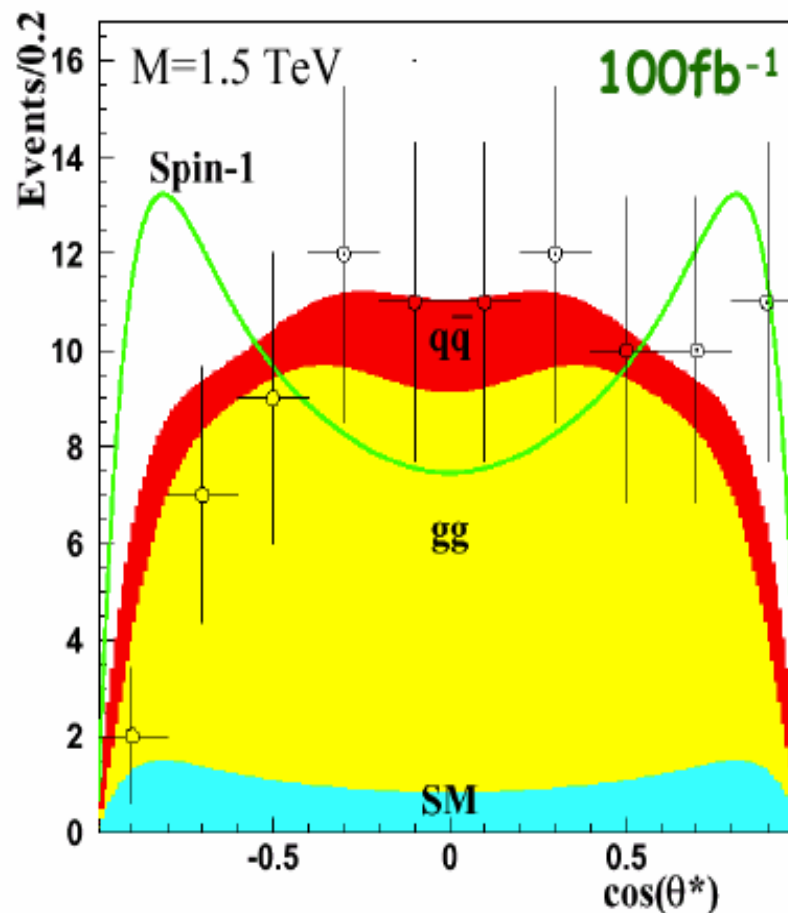
$$gg (q\bar{q}) \rightarrow G \rightarrow e^+ e^-$$

Calculations made using HELAS  
customised for spin-2

Angular distributions

- $qq \rightarrow G \rightarrow ff: 1 - 3\cos^2\theta + 4\cos^4\theta$
- $gg \rightarrow G \rightarrow ff: 1 - \cos^4\theta$
- $qq \rightarrow G \rightarrow VV: 1 - \cos^4\theta$
- $gg \rightarrow G \rightarrow VV: 1 + 6\cos^2\theta + \cos^4\theta$
- DY background:  $1 + \cos^2\theta$

Signature : graviton has spin **2**



ATLAS can distinguish spin 2 vs 1 up to 1.72 TeV

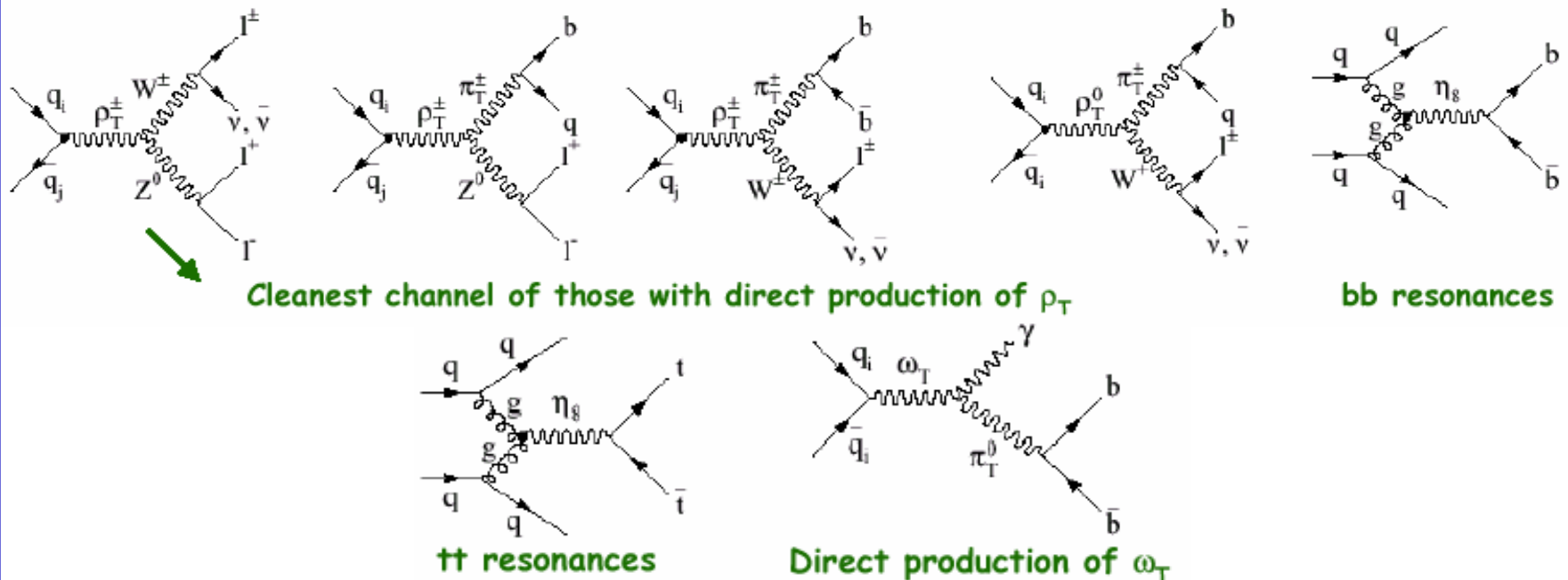


# Strong Symmetry Breaking Technicolor



No fundamental scalar Higgs (it is a new strong force bounded state)  
Technicolor predicts existence of technihadron resonance

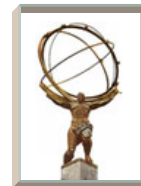
## Technicolor channels investigated



Main background for signals:  $Z$  + jets ( $qq \rightarrow gZ$ ,  $gg \rightarrow qZ$ ,  $qq \rightarrow ZZ$ )  
 $tt$ ,  $WZ$  continuum production

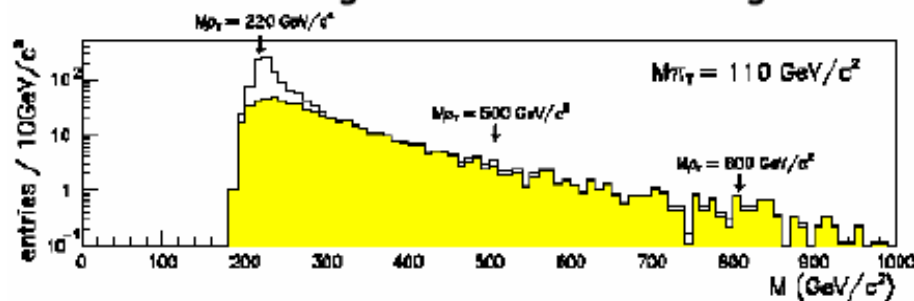


# Signals of the Technicolor

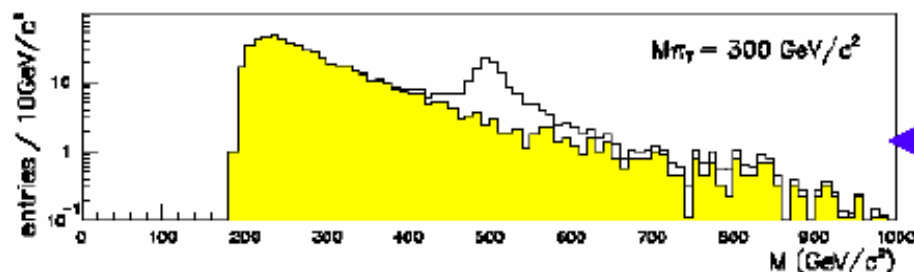


## Examples

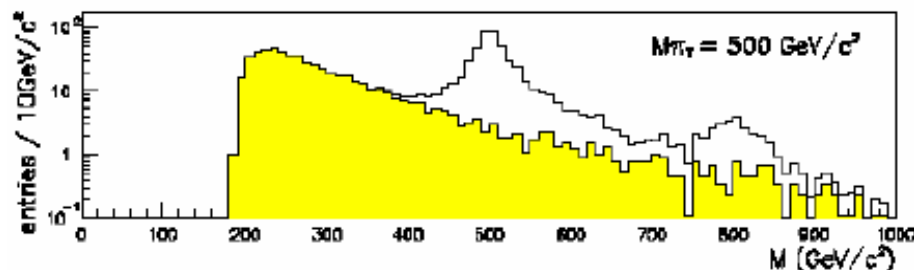
Reconstructed invariant mass for  $\rho_T \rightarrow WZ \rightarrow l\nu ll$  channel.  
Solid line is signal. Filled area is background.



Lower limits required for  $5\sigma$  significance with  $30 \text{ fb}^{-1}$  :  
in some cases,  
signals are below observability,  
but combination of signals  
could provide strong evidence.



$\rho_T \rightarrow WZ \rightarrow l\nu ll$  for  $30 \text{ fb}^{-1}$   
(a)  $\sigma \times \text{BR}_{\text{model}} = 0.16 \text{ fb}$   
 $\sigma \times \text{BR}_{5\sigma \text{ discovery}} = 0.025 \text{ fb}$



$\rho_T \rightarrow W\pi \rightarrow l\nu bb$  for  $30 \text{ fb}^{-1}$   
(c)  $\sigma \times \text{BR}_{\text{model}} = 0.064 \text{ fb}$   
 $\sigma \times \text{BR}_{5\sigma \text{ discovery}} = 0.15 \text{ fb}$



# Summary of New Physics Reach at LHC



Any one of those would change the understanding of our universe !

SM Higgs	100 GeV $\sim$ 1 TeV
MSSM Higgs	covers full ( $m_A, \tan\beta$ )
SUSY (squark, gluino)	2.5 - 3 TeV (300 fb <sup>-1</sup> )
New gauge bosons (Z')	< 4.5 TeV (100 fb <sup>-1</sup> )
Quark substructure ( $\Lambda_C$ )	< 25/40 TeV (30/300 fb <sup>-1</sup> )
q*, l*	< 6.5/3.4 TeV (100 fb <sup>-1</sup> )
Large ED ( $M_D$ for n=2,4)	< 9/5.8 TeV (100 fb <sup>-1</sup> )
Small ED ( $M_C$ )	< 5.8 TeV (100 fb <sup>-1</sup> )
Black holes	< 6 ~ 10 TeV
M(top quark)	$\sigma_M \sim 1$ GeV ( $\sim 0.5$ %)
$M_W$	$\sigma_M \sim 15$ MeV
CP-violation in B-decay	$\sigma(\sin 2\beta) \sim 0.016$ (30 fb <sup>-1</sup> )
Rare B-decay ( $B_s \rightarrow \mu\mu$ )	$\sim 5\sigma$ (130 fb <sup>-1</sup> )

Discovery for sure  
+ some measurements

can say "final word"  
about (low E) SUSY



Both experiments can  
cope with the new physics  
possibilities which were  
not foreseen at the  
beginning of the project.



# Cosmological Connection



- Extremely tempting to assume that EWSB and Dark Matter  
are characterised by the same energy scale
- Likely that new physics contains a stable particle that can be  
copiously produced at the LHC

There are counterexamples, but

if above true  $\Rightarrow$  large cross sections for jets + missing  
energy events at the LHC

$\Rightarrow$  LHC will provide data for astrophysics

$\Rightarrow$  infer DM properties from masses and  
cross sections

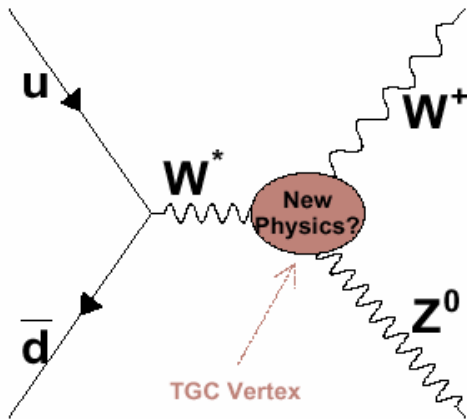
Relic density  $\Omega_\chi h^2 \sim 3 \times 10^{-27} \text{ cm}^3 \text{s}^{-1} / \langle \sigma v \rangle$

requires typical weak interaction annihilation cross sections

How well  $\langle \sigma v \rangle$  can be predicted from LHC depends on model for NP



# Triple Gauge Boson Couplings

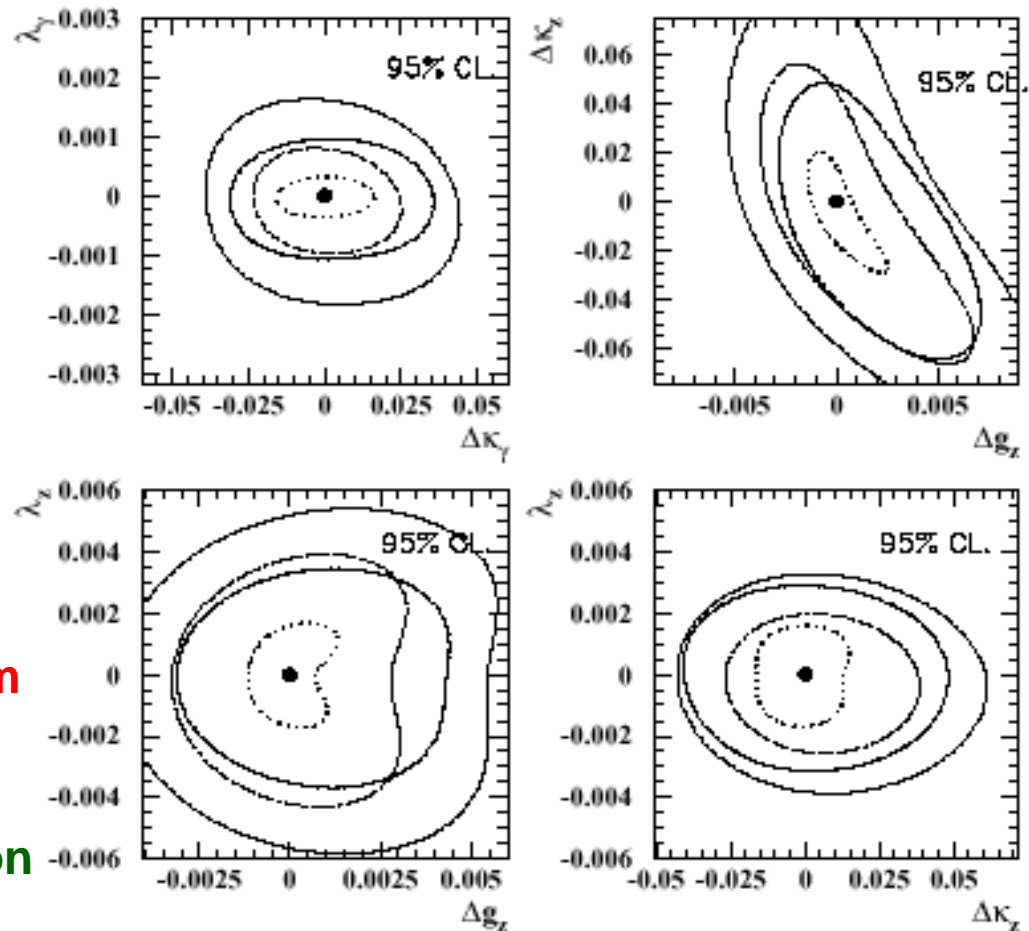


non-abelian  $SU(2)_L \times U(1)_Y$   
gauge group (foundation of SM!)

**Open window to electroweak  
Symmetry breaking mechanism**

**LHC: orders of magnitude  
Improvement over LEP/Tevatron**

*Triple gauge boson couplings*



Expected 95% C.L. constrains contours (outer  $\rightarrow$  inside):

(14TeV, 100fb<sup>-1</sup>), (28TeV, 100fb<sup>-1</sup>), (14TeV, 1000fb<sup>-1</sup>), (28TeV, 1000fb<sup>-1</sup>)