

# Dark Matter II

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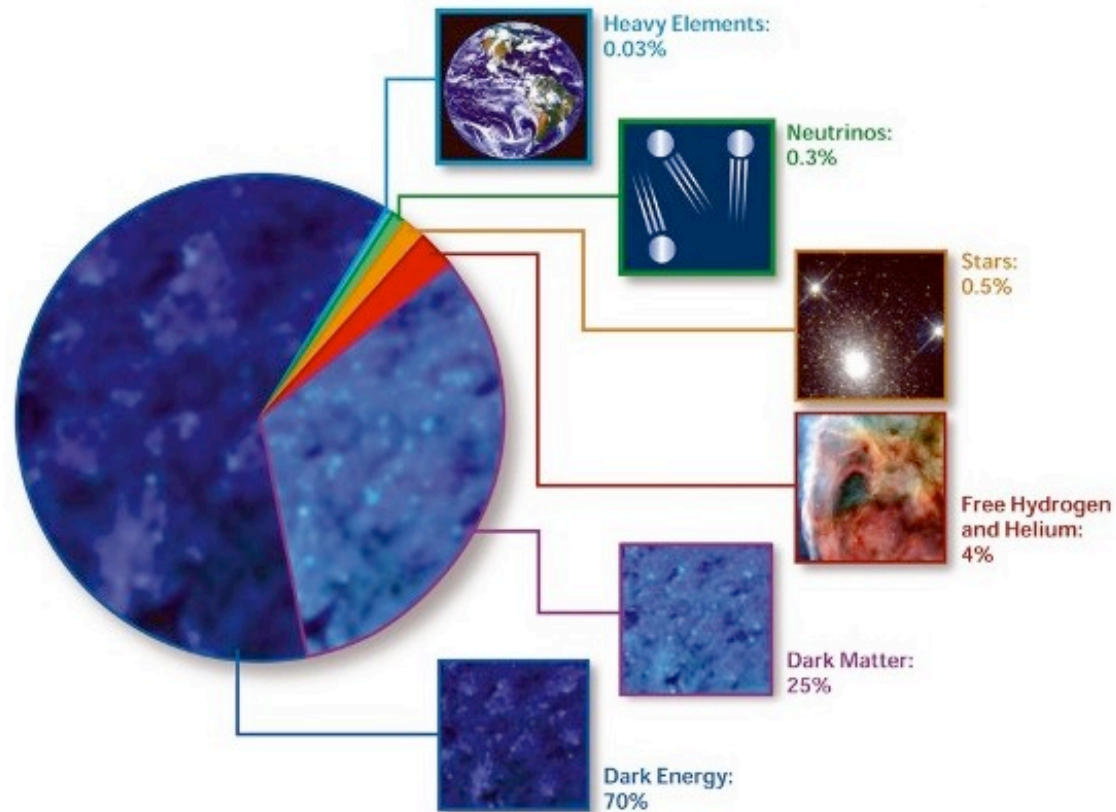
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## Reminder

## Constituents of the present Universe:



# What is Dark Matter?

# Properties of Dark Matter

- Nature of dark matter
  - Is it baryonic?
  - Is it a new type of elementary particle?
- Mass
- Lifetime
  - Must be long lived or stable.
- Kinetic energy expressed in Temperature
- Spin
- Cross section of self-interaction
- Cross section of interaction with ordinary matter

# 'Temperature' of Dark Matter

- Hot

- Moving fast
- Can overcome and escape potential wells
- If all dark matter was hot, structure could not form
  - Contradicts findings of WMAP.

- Warm

- Cold

- Moving slowly
- Could seed structure

# Baryonic Dark Matter

- Cold hydrogen
  - 25% of the hydrogen is in single-atom state.
- Massive Compact Halo Objects (MACHOs)
  - Brown dwarfs/Large planets
    - Not enough around to explain all dark matter
  - Compact stars, such as White Dwarfs
    - Too many of them will need more He than could be produced by Big Bang nucleosynthesis
  - Neutron stars and Black holes
    - Rarer than white dwarfs.
- Constraints from micro-lensing
  - Observed several millions stars for years
    - In Large Magellan Cloud, found 4 candidates
    - Near the galactic centre, found 45 candidates
  - < 20% of the Milky Way halo is MACHOS with

# Non-baryonic Dark Matter

- Hot - Neutrinos
- Warm - Sterile neutrinos, gravitino
- Cold
  - LSP (Lightest Super-symmetric Particle, eg. neutralino, axino)
  - LKP (Lightest Kaluza-Klein particle)
  - Axions, axion clusters
  - Solitons (Q-balls, B-balls)
  - WIMPs (Weakly Interacting Massive Particles), wimpzilla

# Finding The Correct Dark-matter Candidate

- In early Universe, candidate  $x$  is in thermal equilibrium by annihilation and couples to ordinary matter:



Number density follows

the Boltzmann distribution:

$$n_x^{\text{eq}} \propto (mT)^{3/2} e^{-m/kT}$$

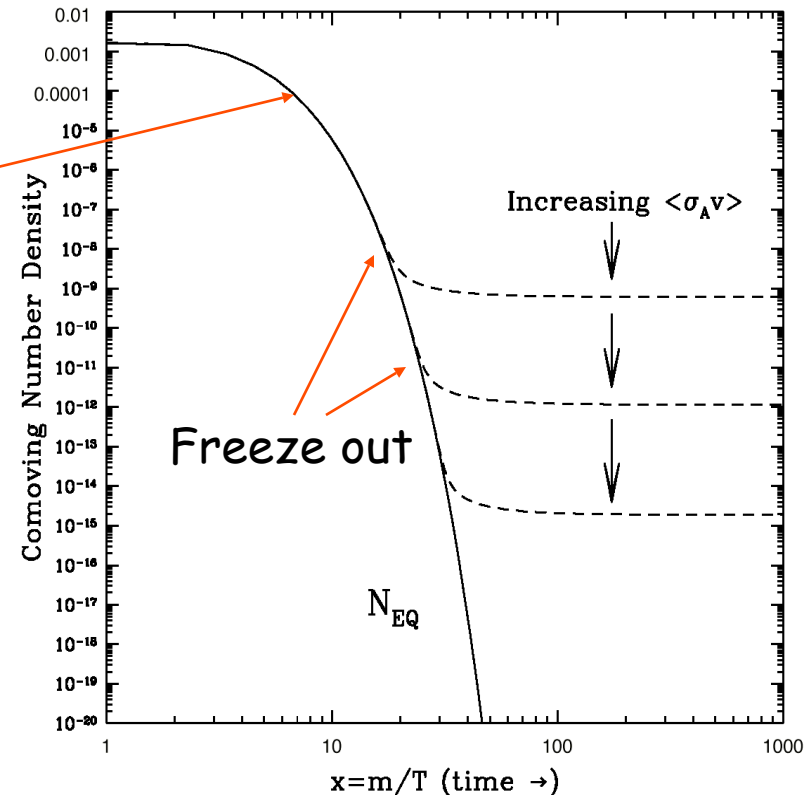
When  $H(t) > \Gamma_{\text{ann}} \sim n_x \langle \sigma_a v \rangle$ ,

$n_x$  freezes.

After freeze out:

$$\Omega_x = \frac{n_x(\text{now})}{\rho_c(\text{now})} \approx \frac{3 \times 10^{-27} \text{ cm}^3 \text{ sec}^{-1}}{\langle \sigma_a v \rangle h^2}$$

↑  
Hubble parameter  
 $\approx 0.73$



- How to find the correct dark-matter candidate?
  - choose a particle
  - work out all the possible annihilation channels
  - calculate  $\Omega_x$

# Dark-matter Candidate: Neutrinos

- Promising candidate
  - Neutral
  - Weakly interacting
  - Possess mass
  - Their abundance in the Universe is known
- From relic abund  $\Omega_v \approx \frac{\sum m_v}{50\text{eV}} = 113/\text{cm}^3$  per flavour :
- From WMAP+BAO+ $H_0$  combined result (2010):
$$\sum m_v < 0.58 \text{ eV}$$
$$\Omega_v \ll \Omega_m$$
- Active neutrinos are not abundant and massive enough to account for dark matter.



# Dark-matter Candidate: Sterile Neutrinos

- Mass of sterile neutrinos  $< \sim \text{keV}$  (Tremaine-Gann bound)
- Have no coupling to  $W$  and  $Z$ 
  - Interact with ordinary matter extremely weakly through mixing with the active neutrinos:

$$\text{Strength} \sim \theta_\nu G_F$$

where  $\theta_\nu \ll 1$  is the mixing angle

$$\theta_\nu^2 = \sum_{\alpha=e,\mu,\tau} \frac{\lambda_{\alpha\nu}^2}{M_\nu^2}$$

Majorana neutrino mass,  $M_\nu < 10^{-7} \text{ GeV}$

Yukawa couplings  $\lambda_{ij} \sim 10^{-7} \sqrt{M_\nu / \text{GeV}}$  (proportional to  $\Delta m^2$ )

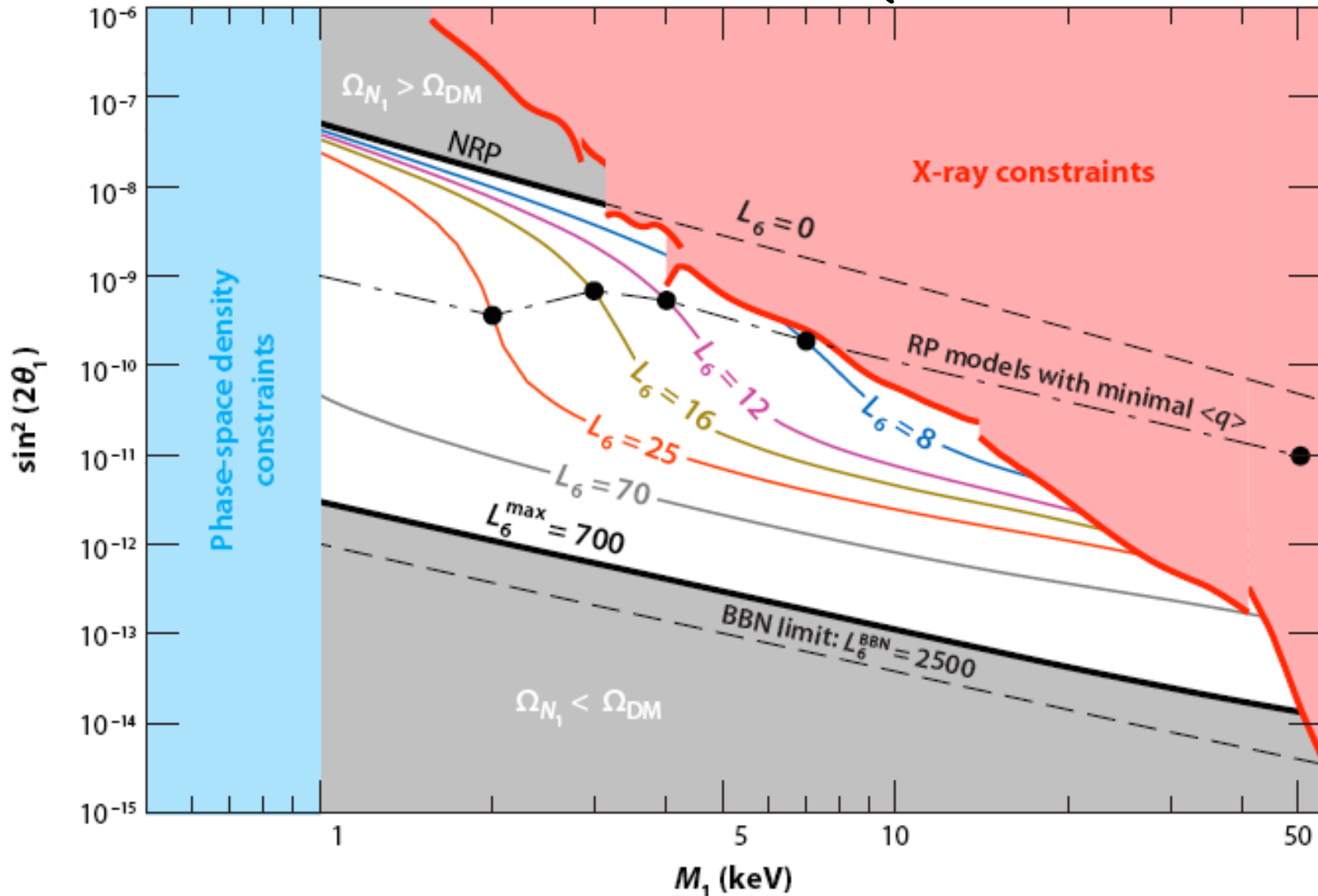
- Sterile neutrinos can decay to active neutrinos

$$\nu_s \rightarrow \nu + \gamma$$

- to be a dark-matter candidate,  $\nu_s$  must have a very long lifetime greater than the age of the Universe, setting a bound on  $\theta_\nu$ .

# Bound on Mass of Sterile Neutrinos

Mass of sterile neutrinos  $< \sim \text{keV}$  (Tremaine-Gann bound)

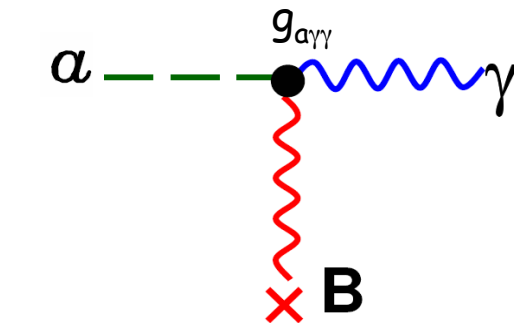


# Dark-matter Candidate: Axion

- Arise from the Peccei-Quinn's idea to solve the strong CP problem.
- Axion got its mass during the QCD phase transition:

$$m_a \approx 0.62 eV \frac{10^7 GeV}{f_a} \quad f_a = \text{energy scale of the transition}$$

- For cold dark matter,  $f_a < 10^{12} GeV$ , implying  $m_a = 10^{-6}-10^{-4} eV$ .
- Axion is extremely weakly interacting.
- It can couple to electromagnetic field:



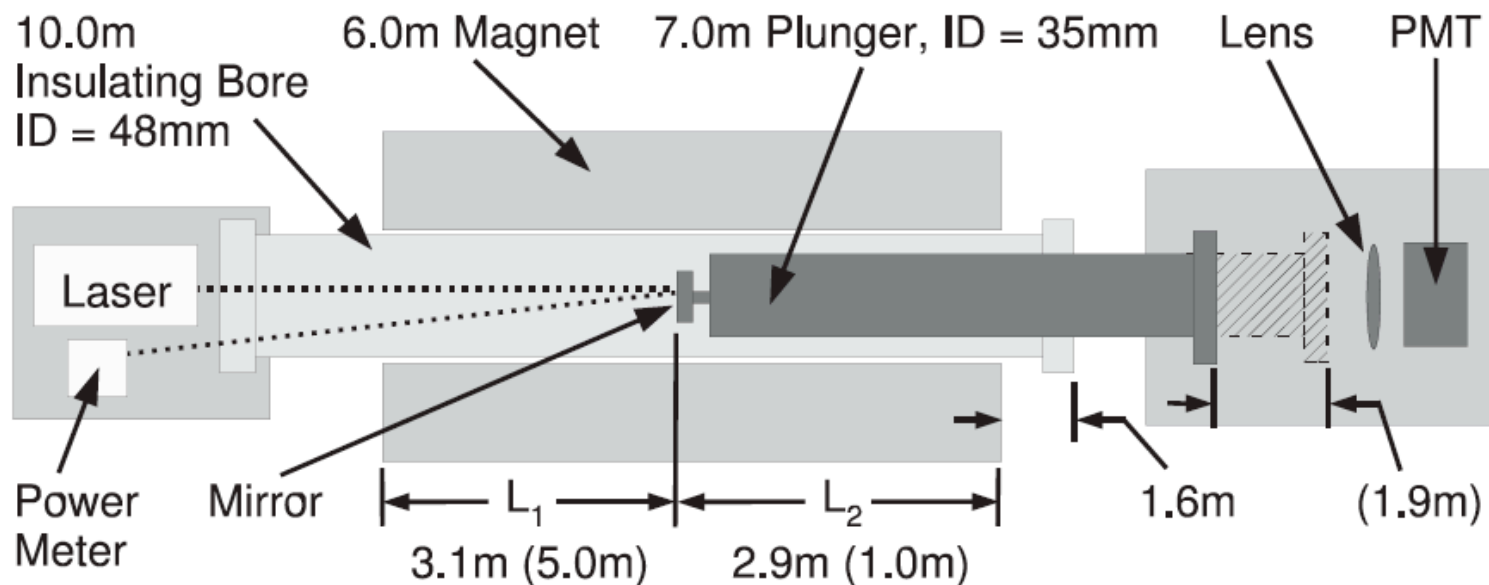
$$g_{\alpha\gamma\gamma} \approx \left( \frac{\alpha}{4\pi} \right) \frac{1}{f_a}$$

## The Birth of Axions

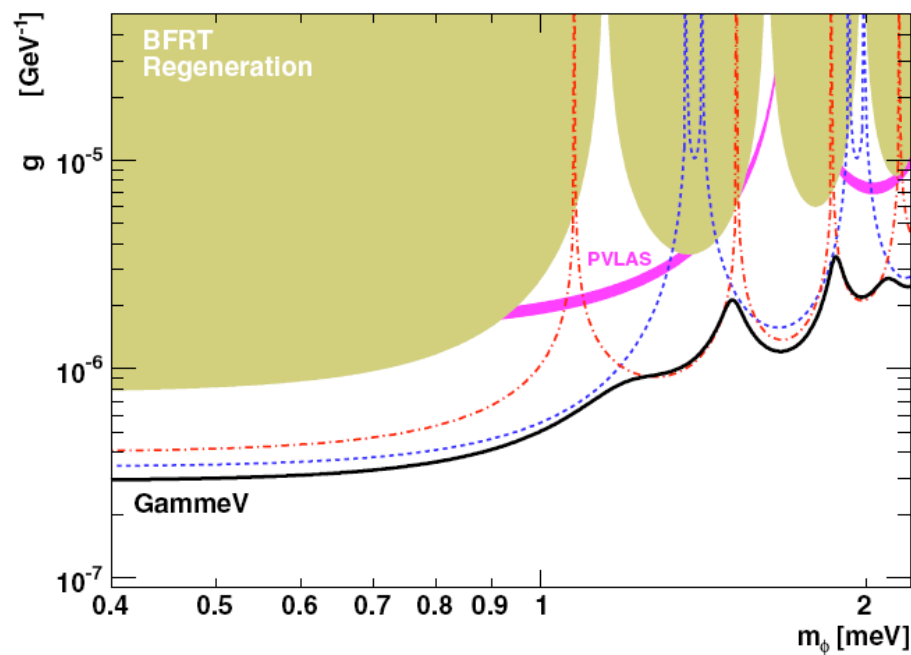
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usual, very light particle. I called this particle the *axion*, after the laundry detergent, because that was a nice catchy name that sounded like a particle and because this particular particle solved a problem involving *axial* currents.

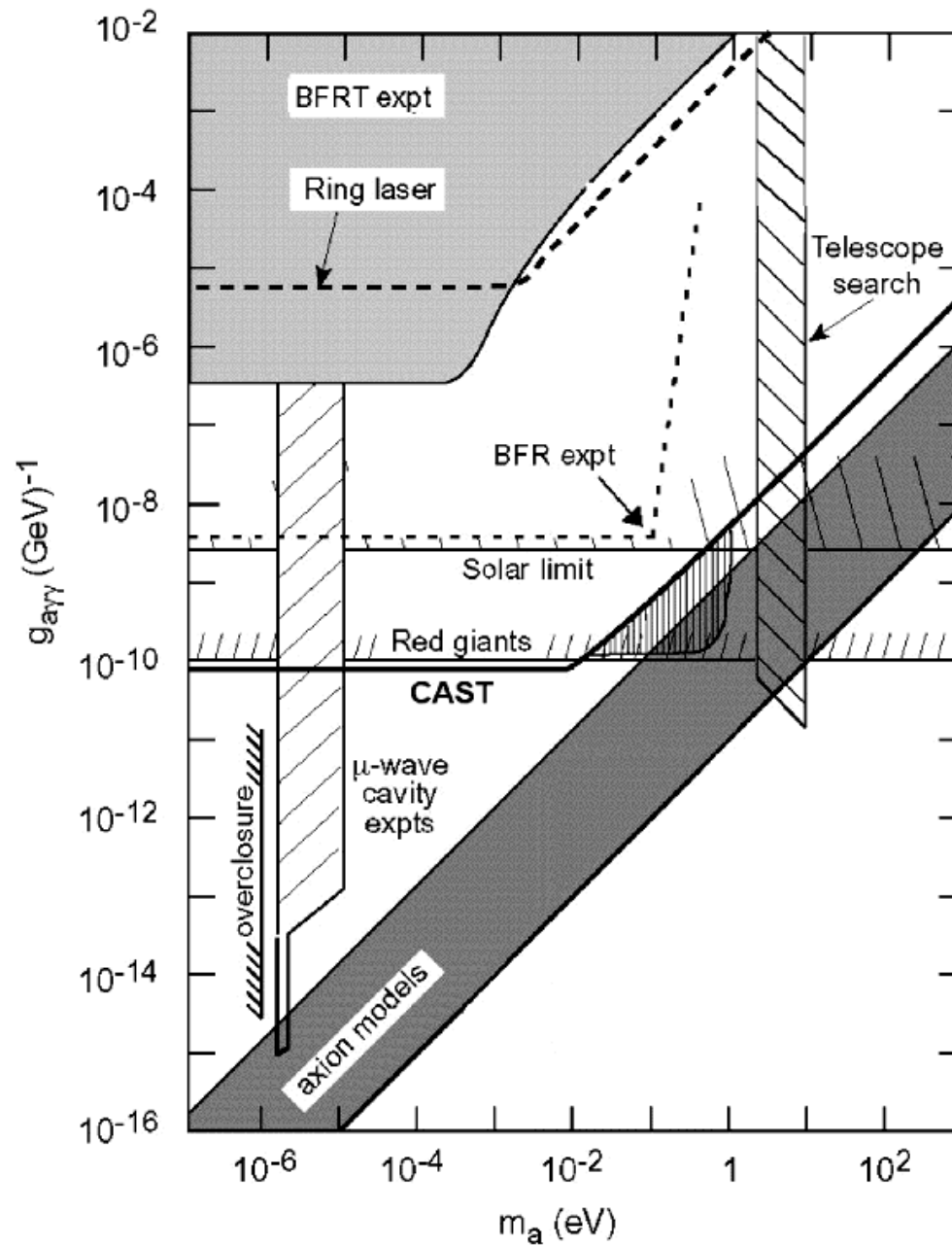
# GammeV: Search For Axion



Bargain, Wester (left) and Chou say their experiment offers high potential payoff at low cost.



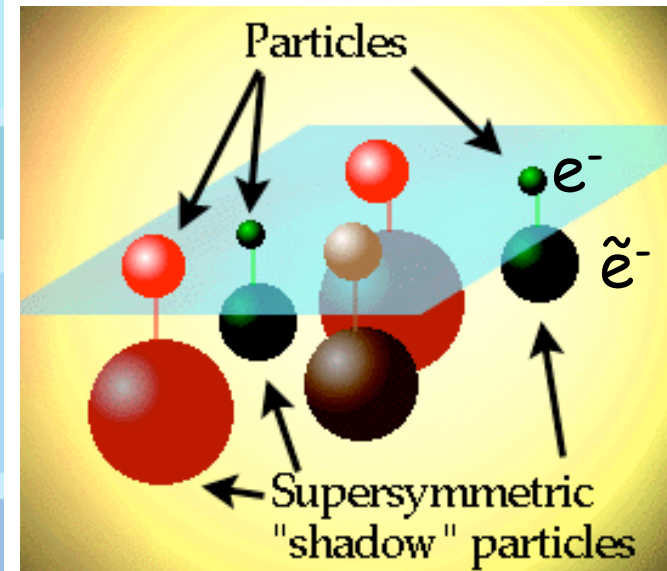
# Limits on Axion



# Supersymmetry (SUSY)

- A new kind of symmetry between fermions and bosons
  - Every particle in SM has a superpartner with same quantum numbers

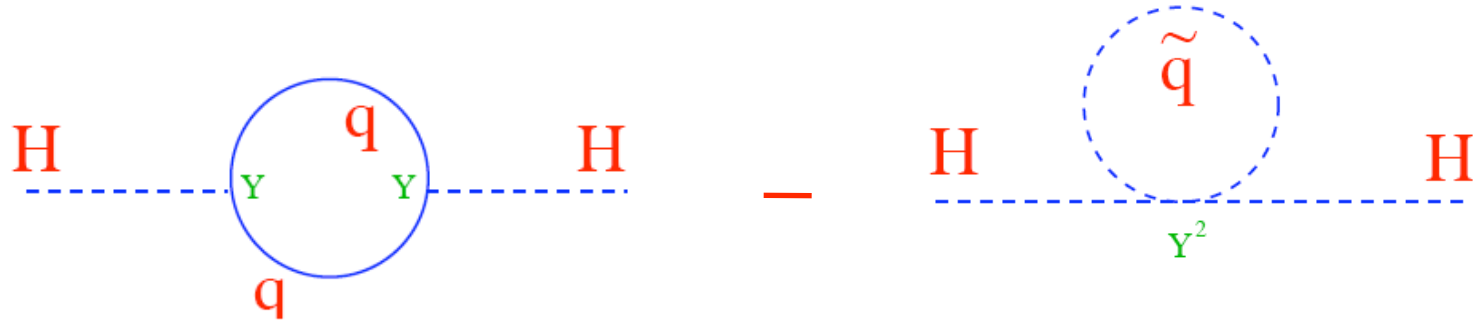
Ordinary Particles		Supersymmetric Partners	
Higgs Boson (spin-0)		Higgsino (spin-1/2)	
Fermions (spin-1/2)		Bosons (spin-0)	
Quarks	Leptons	Squarks	Sleptons
Gauge Bosons (spin-1)		Gauginos (spin-1/2)	
W <sup>±</sup>	Z, B	Winos	Zinos, Binos
	gluons, photons		gluinos, photinos
<i>charged</i>	<i>neutral</i>	<i>charginos</i>	<i>neutralinos</i>
Graviton (spin-2)		Gravitino (spin-3/2)	



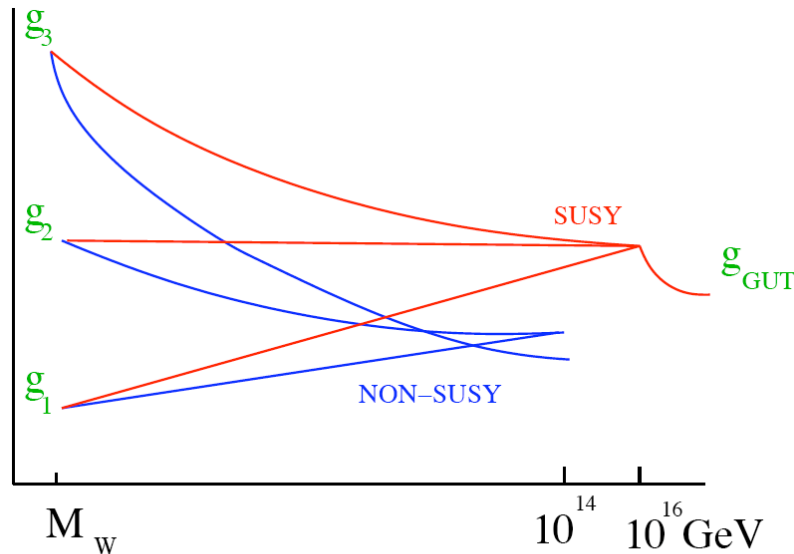
- Interactions of SUSY particles = interactions of quarks and leptons.
- If SUSY was exact,  $m_{\text{sparticle}} = m_{\text{particle}}$  and would have contributed to the Z decay width → SUSY is broken

# Goodies of SUSY

- Remove divergences in calculating  $m_H$ :



- Coupling constants are unified to a single coupling constant at  $M^{\text{SUSY}}$ :



$$M_X^{\text{SUSY}} = 10^{16} \text{ GeV}$$



# MSSM

- MSSM (Minimal Supersymmetric Standard Model) contains the smallest number of new particles and new interactions as well as all supersymmetry-breaking terms.

$$H_d = \begin{pmatrix} H_d^0 \\ H_d^- \end{pmatrix}, \quad H_u = \begin{pmatrix} H_u^+ \\ H_u^0 \end{pmatrix}$$

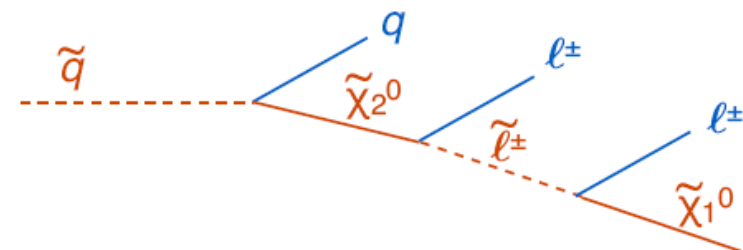
with vacuum expectation values:

$$\langle H_d \rangle_0 = \begin{pmatrix} v_d \\ 0 \end{pmatrix}, \quad \langle H_u \rangle_0 = \begin{pmatrix} 0 \\ v_u \end{pmatrix}$$

such that

$$v_d^2 + v_u^2 = v^2 \quad \tan\beta = \frac{v_u}{v_d} \quad 0 \leq \beta \leq \frac{\pi}{2}$$

- R-parity (to prevent proton decay):
  - $R = (-1)^{3B+L+2s}$        $B = \text{baryon no.}, L = \text{lepton no.}, s = \text{spin}$
  - $R = +1$  (even) for particles;  $R = -1$  (odd) for SUSY particles
  - R-parity is conserved
  - The lightest sparticle (LSP) is absolute stable, e.g. gravitino, neutralino





# The Lightest SUSY Particle

- Gravitinos
  - Could be overproduced in the early Universe and destroy abundance of primordial elements
- Sneutrinos
  - Scattering cross section is much larger than the limits found by direct-detection experiments
- Neutralinos
  - 4 Majorana fermionic eigenstates arise from mixing of bino, wino and higgsinos

$$M_{\chi^0} = \begin{pmatrix} m_1 & 0 & -M_Z c_\beta s_W & M_Z s_\beta s_W \\ 0 & m_2 & M_Z c_\beta c_W & -M_Z s_\beta c_W \\ -M_Z c_\beta s_W & M_Z c_\beta c_W & 0 & -\mu \\ M_Z s_\beta s_W & -M_Z s_\beta c_W & -\mu & 0 \end{pmatrix}$$

$m_1$  = bino mass parameter

$m_2$  = wino mass parameter

$\mu$  = higgsino mass parameter

$c_\beta = \cos\theta_\beta$        $s_\beta = \sin\theta_\beta$

$c_W = \cos\theta_W$        $s_W = \sin\theta_W$

# Minimal Supergravity (mSUGRA)

- Reduce to the ~100 parameters to 4 parameters and 1 sign:

$\tan \beta$

$M_{1/2}$  unified gaugino mass at  $M_{\text{GUT}}$

$m_0$  unified scalar masses at  $M_{\text{GUT}}$  ( $M_\chi \approx 0.4 m_0$ )

$A_0$  trilinear soft-breaking mass

- The lightest neutralino in mSUGRA:  
 $\text{sign}(\mu)$  sign of higgsino mass parameter

$$\chi_1^0 = c_1 \tilde{B} + c_2 \tilde{W} + c_3 \tilde{H}_u^0 + c_4 \tilde{H}_d^0$$

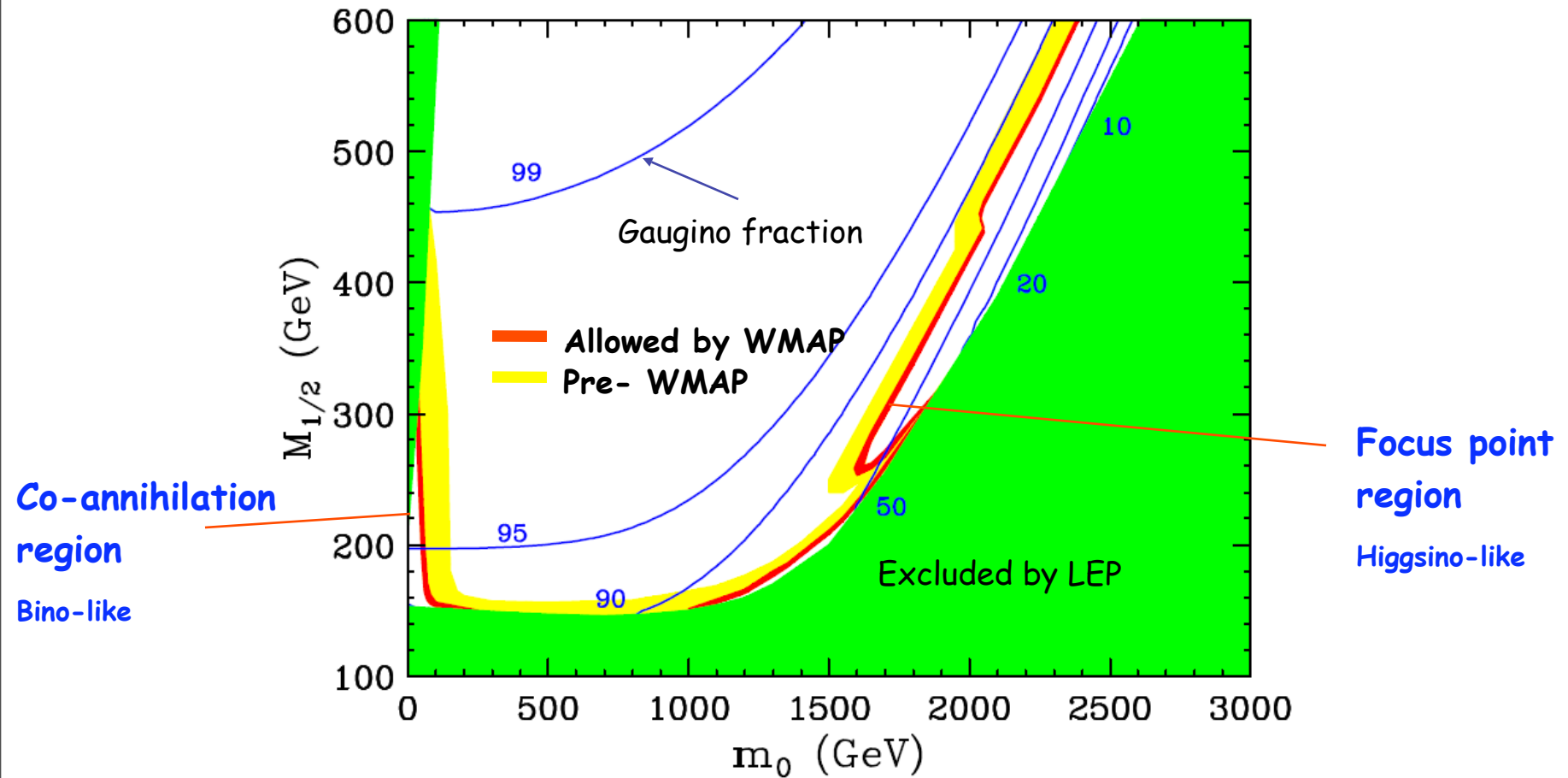
Photino    Wino    Higgsinos

$$\Rightarrow \chi_1^0 = c_U (\tilde{B} + \tilde{W} + \tilde{H}_u^0 + \tilde{H}_d^0)$$

- The mass of  $\chi_1$  : 100 GeV - 1 TeV, weak scale.

# Where To Look For Neutralino ?

mSUGRA with  $\tan\beta = 10$ ;  $A_0 = 0$ ;  $\mu > 0$



K. Matchev hep-ph/0402088

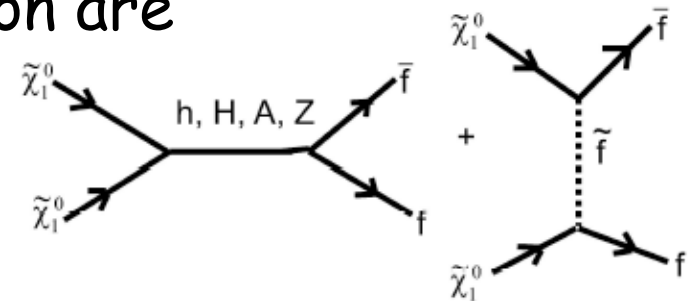
# Neutralino As Dark Matter Candidate

- At present, neutralinos are moving at non-relativistic speed:

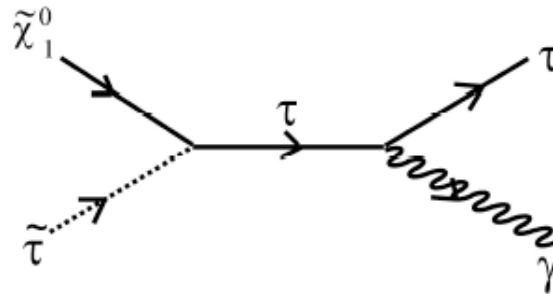
$$\sigma_a v = a + bv^2 + O(v^4) \approx a$$

- The dominant modes of annihilation are

- Fermion-antifermion pairs
- Gauge boson pairs
- Final states containing the Higgs

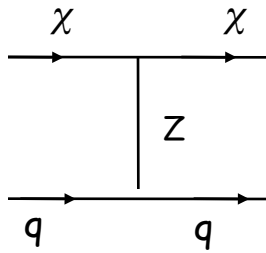


- Other interactions for dark-matter searches are:
  - Co-annihilation, especially when  $M_\gamma \approx M_{\tilde{\tau}}$

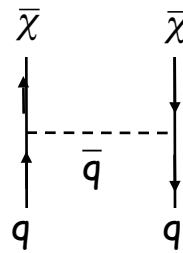


# Elastic Scattering of Neutralinos

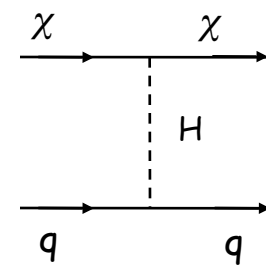
- Processes for detecting neutralinos with ordinary matter:



Spin-dependent



both



Spin-independent

- The cross sections of these processes:

$$\sigma_A = 4G_F^2 \left( \frac{m_\chi M_A}{m_\chi + M_A} \right)^2 C_A F(q^2)$$

$C_A^{SI}$  : for spin-independent coherent interaction  $\propto A^2$

$C_A^{SD}$  : for spin-dependent interaction  $\propto \langle S_{p,n} \rangle^2$

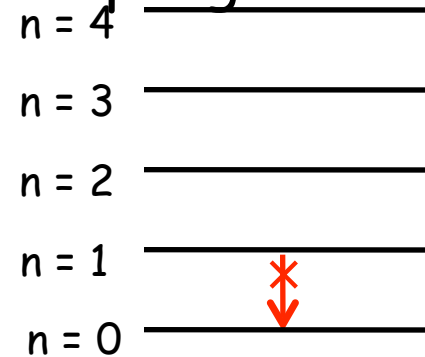
$F(q^2)$  : nuclear form factor, important for large  $q^2$  and large  $A$

# Universal Extra Dimensions (UED)

- All particles can move into the extra compact dimensions with  $R \sim 1 \text{ TeV}^{-1}$ .
- For each particle, there is an infinite tower of states with identical quantum numbers and couplings but with masses given by

$$M_n^2 \propto \frac{n^2}{R^2}$$

SM particles have  $n = 0$ .



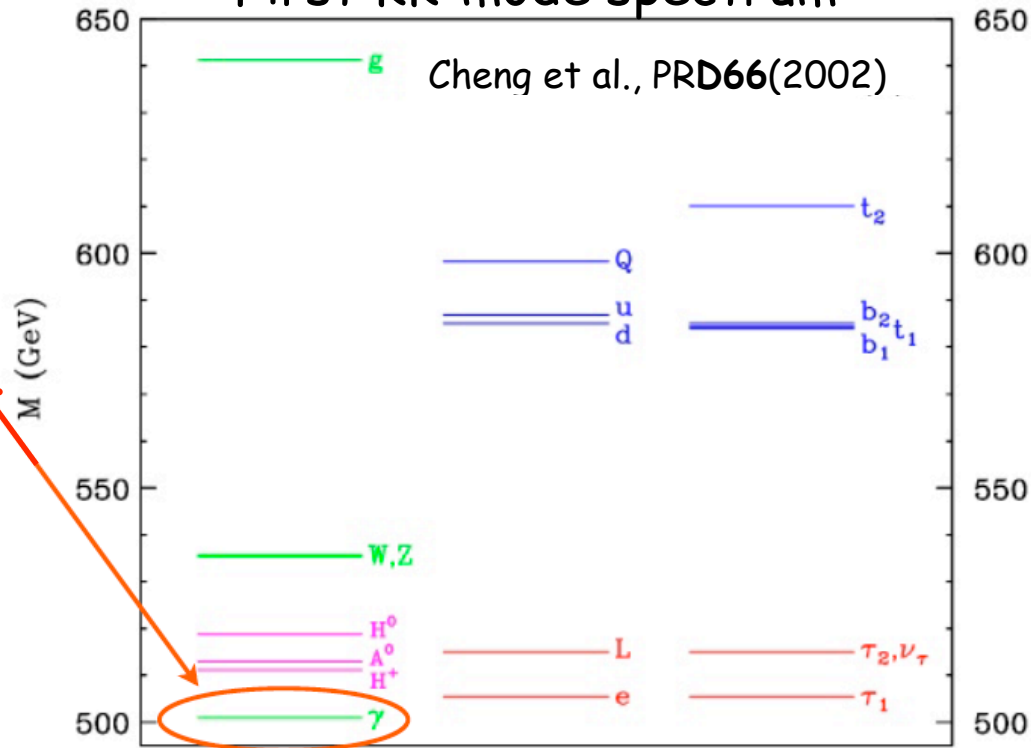
- In the simplest UED model, along the fifth dimension
  - The lightest Kaluza-Klein (KK) state is stable

# Lightest Kaluza-Klein Particle

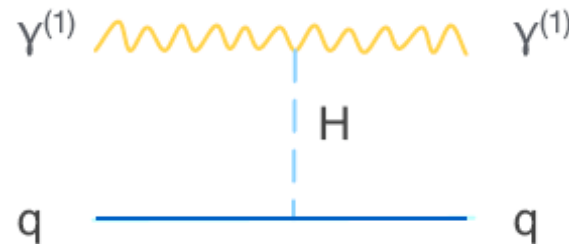
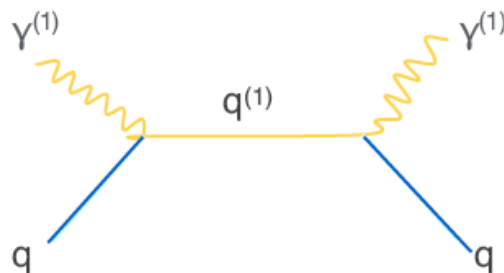
First KK-mode spectrum

Most likely LKP

$\gamma^{(1)}$



- Possible interactions of LKP with ordinary matter:



# Dark-matter Candidates : Summary

